

Mahidol University
Wisdom of the Land

Interacting dark sectors from a dynamical system perspective

Nandan Roy

Centre for Theoretical Physics and Natural Philosophy
Mahidol University, Thailand

In collaboration with Narayan Banerjee and Chonticha Kritpetch

State of the Art

- ❑ Our universe is not only expanding but it is also accelerating!!
- ❑ Λ CDM model has been constrained with unprecedented accuracy.
- ❑ With the improvement in our ability to constrain the cosmological parameters, a few statistically significant tensions has emerged.
- ❑ It seems that the late time cosmological data and early time cosmological data are in tension.
- ❑ We need to extent our imagination beyond standard Λ CDM.

Hubble Tension

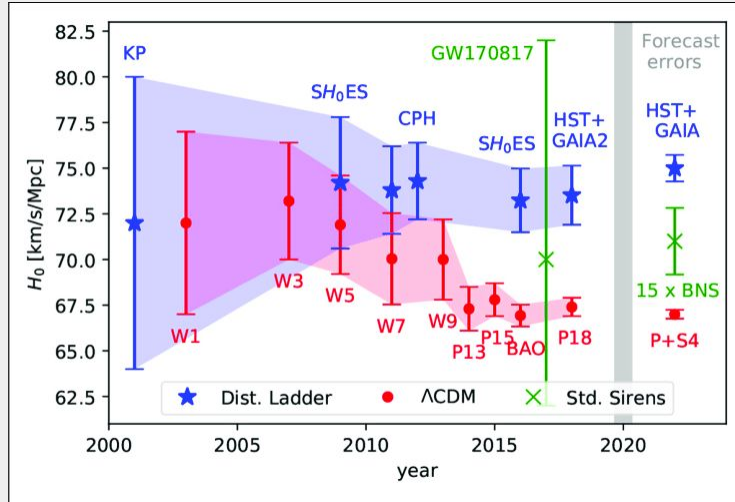


Fig 1(a)

CMB Planck data together with BAO, BBN, and DES have constrained the Hubble parameter to be $H_0 \sim (67.0 - 68.5) \text{ km/s/Mpc}$. On the other hand, cosmic distance ladder and time delay measurement like those reported by SH0ES and HOLiCOW collaborations have reported $H_0 = (74.03 \pm 1.42) \text{ km/s/Mpc}$ and $H_0 = (73.3 +1.7 -1.8) \text{ km/s/Mpc}$ respectively by observing the local Universe.

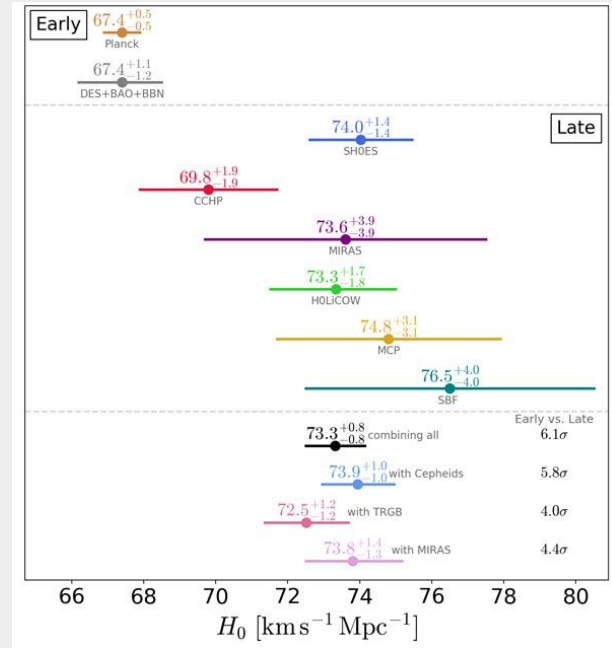
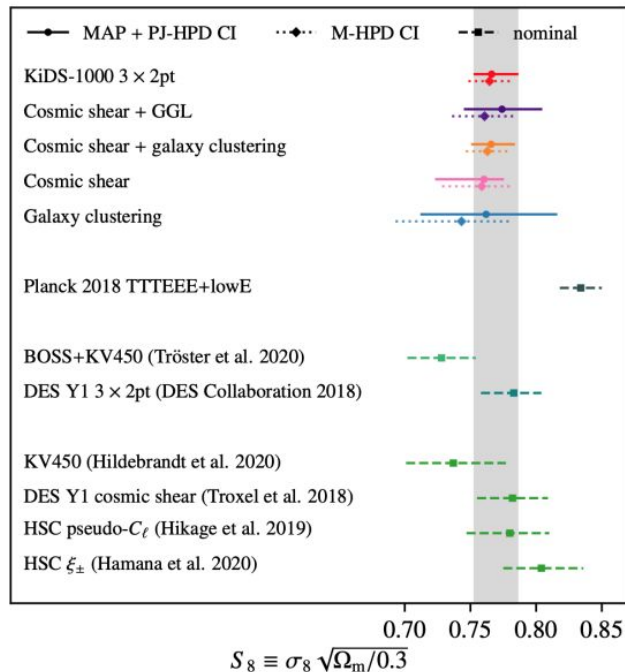
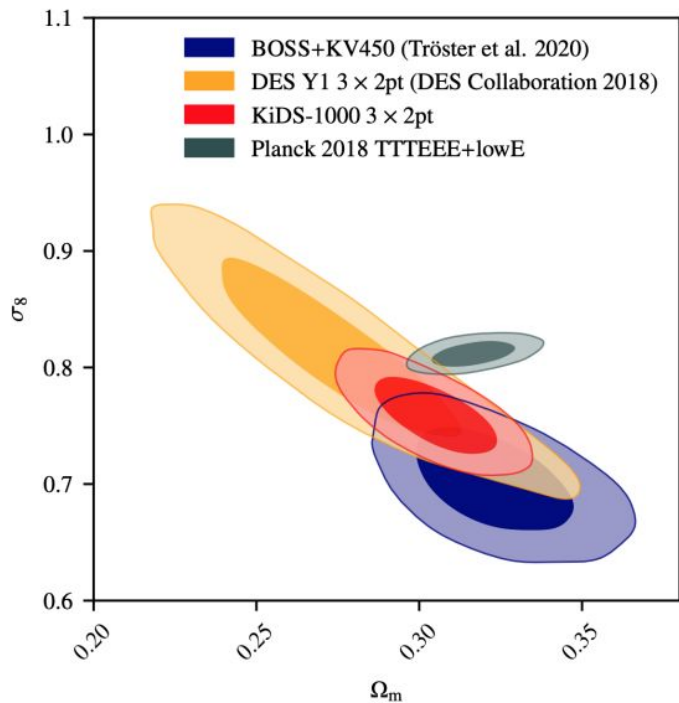


Fig 1(b)

σ_8 Tension



Apart from the Hubble tension, another tension between the Planck data with the weak lensing and the redshift surveys has been reported.

Fig 2

Scalar Field as Dark energy

- The Λ CDM model happens to be most consistent with the observations but it suffers from problems arising from both theoretical and observational aspects.
- From the theoretical side it has to deal with the cosmological constant problem, coincidence problem and the fine tuning problem.
- From the observational side it is unable to explain the tension between the early time (Planck, BAO) and late time observations (SH0ES).
- There could be new physics involved and we should think beyond Λ CDM model.
- Scalar fields models are considered as one of the best alternatives to the cosmological constant.

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{2\kappa^2} R + \mathcal{L}_\phi \right] + S_M$$

$$\mathcal{L}_\phi = -\epsilon \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi)$$

Scalar Field Dynamics

For a spatially flat, homogeneous and isotropic universe filled with matter and nonminimally coupled scalar field components

$$3H^2 = \rho_m + \rho_\phi = \rho_m + \frac{1}{2}\epsilon\dot{\phi}^2 + V(\phi)$$

$$2\dot{H} + 3H^2 = -p_\phi = -\frac{1}{2}\epsilon\dot{\phi}^2 + V(\phi) \quad 1)$$

$$\epsilon\ddot{\phi} + 3\epsilon H\dot{\phi} + \frac{dV}{d\phi} = 0$$

$\epsilon \rightarrow$ Switch parameter

$$\epsilon = \begin{cases} +1, & \text{Quintessence} \\ -1, & \text{Phantom} \end{cases}$$

$\rho_m \rightarrow$ Matter energy density

$\rho_\phi = \frac{1}{2}\epsilon\dot{\phi}^2 + V(\phi) \rightarrow$ Scalar field energy density

$p_\phi = \frac{1}{2}\epsilon\dot{\phi}^2 - V(\phi) \rightarrow$ Pressure component

Interacting Dark Sectors

If we take into account the dark energy scenario featuring a modification to the regular conservation equations, with the addition of an interaction to the dark sectors of the universe.

$$\begin{aligned}\dot{\rho}_m + 3H\rho_m &= -Q \\ \dot{\rho}_\phi + 3H(\rho_\phi + P_\phi) &= Q\end{aligned}$$

Interaction term

Continuity equation of the dark matter and dark energy.

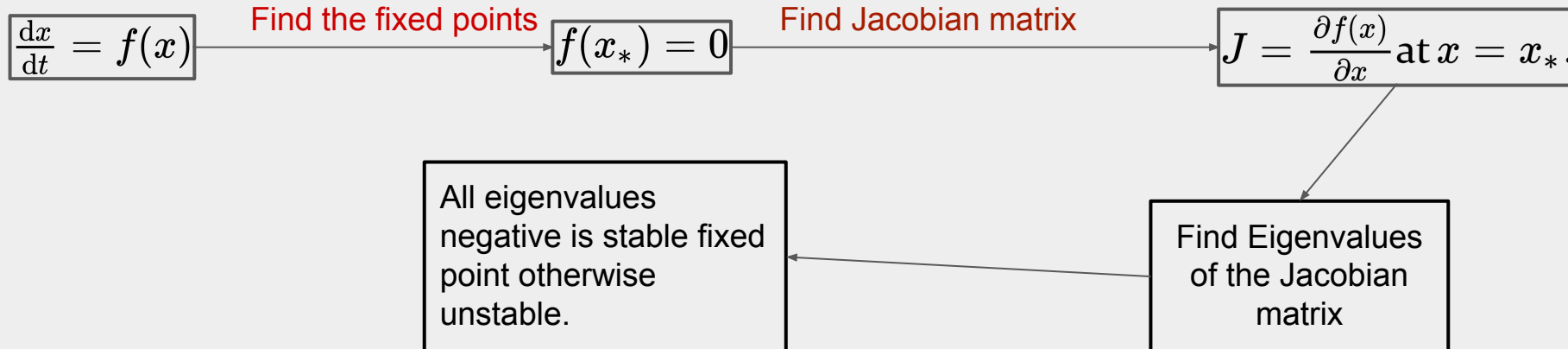
The K.G equation of the scalar field.

$$\ddot{\phi} + 3H\dot{\phi} + \epsilon \frac{dV}{d\phi} = \epsilon \frac{Q}{\dot{\phi}},$$

(3)

Dynamical Systems Analysis

- A dynamical system can be considered to be a model describing the temporal evolution of a system.
- Dynamical system is mathematical tool to understand the qualitative behaviour of the linear or nonlinear systems where finding exact solution to the system is not trivial.
- Dynamical system analysis has been used extensively in many areas of science and engineering.
- Dynamical system analysis has been also widely used in cosmology.



Construction of the autonomous system

$$x = \frac{\kappa\dot{\phi}}{\sqrt{6}H}, y^2 = \frac{\kappa^2 V(\phi)}{3H^2}, \lambda = -\frac{1}{V(\phi)} \frac{dV(\phi)}{d\phi}. \quad (5)$$

$$\begin{aligned} \Omega_\phi &= \epsilon x^2 + y^2, \\ w_\phi &= \frac{\epsilon x^2 - y^2}{\epsilon x^2 + y^2}, \\ q &= -1 + \frac{3}{2}(1 + \epsilon x^2 - y^2). \end{aligned}$$

$$\begin{aligned} x' &= -3x + \sqrt{3/2}\epsilon\lambda y^2 \\ &+ \frac{3}{2}x(1 + \epsilon x^2 - y^2) + \epsilon f(x, y), \\ y' &= -\sqrt{3/2}\lambda xy \\ &+ \frac{3}{2}y(1 + \epsilon x^2 - y^2), \\ \lambda' &= -\sqrt{6}\lambda^2(\Gamma - 1)x, \end{aligned}$$

$$Q = \sqrt{6}\dot{\phi}H^2 f(x, y)$$

Interaction term

$$\Gamma = \frac{V(\phi) \frac{\partial^2 V(\phi)}{\partial \phi^2}}{\left(\frac{\partial V(\phi)}{\partial \phi}\right)^2}$$

(6) Potential

Cosmological parameters in terms of dynamic variables.

Classification of Potential

Depending in the choice of $\Gamma = \frac{V(\phi) \frac{\partial^2 V(\phi)}{\partial \phi^2}}{\left(\frac{\partial V(\phi)}{\partial \phi}\right)^2}$ potential can be classified in two classes.

$$\lambda' = -\sqrt{6}\lambda^2(\Gamma - 1)x,$$

- $\Gamma = 1$ \longrightarrow Exponential Potential
- $\Gamma \neq 1$ \longrightarrow Non-exponential Class of Potentials

Exponential Potential

For the exponential potential $\Gamma = 1$ so that λ is constant.

The system reduces to a 2D system

$$\begin{aligned}x' &= -3x + \sqrt{3/2}\epsilon\lambda y^2 + \frac{3}{2}x(1 + \epsilon x^2 - y^2) + \epsilon f(x, y), \\y' &= -\sqrt{3/2}\lambda xy + \frac{3}{2}y(1 + \epsilon x^2 - y^2),\end{aligned}\tag{7}$$

Fixed Points

List of the fixed points for the exponential potential

Fixed Points	x	y	$f(x, y)$
Q_1	0	0	0
Q_{2i}	$\frac{3}{2}x(1 + \epsilon x^2)$ $-3x + \epsilon f(x, 0) = 0$	0	$f(x, y) = 0$ for $x = 0, x \pm 1$
Q_{3i}	$\frac{1}{2} (\epsilon (2f(x, y) + \sqrt{6}\lambda) + 2\sqrt{6}\lambda x^2 - 2x (\lambda^2\epsilon + 3)) = 0$	$\pm \frac{\sqrt{3x^2\epsilon - \sqrt{6}\lambda x + 3}}{\sqrt{3}}$	$f(x, y)$

Q_1 is completely matter dominated and unstable in nature.

$$\left[\frac{3}{2}, \frac{1}{2} (2\epsilon \partial_x f - 3) \right]$$

Q_2 Depending on the choice of the form of interaction $f(x, y)$ it could be either unstable or saddle.

Q_3 Depends on the choice of the interaction $f(x, y)$.

$$T_{Q_3} = -6 + \epsilon \partial_x f$$

$$D_{Q_3} = -3\epsilon \partial_x f + \sqrt{\frac{3}{2}} \lambda \epsilon \partial_y f + 3\epsilon \lambda^2 + 9$$

For $x = 0, y = \pm 1$.

$$\left[\frac{1}{2} (3x^2\epsilon - \sqrt{6}\lambda x + 3), \frac{1}{2} (2\epsilon \partial_x f + 9x^2\epsilon - 3) \right]$$

Non-exponential class of potentials

For this case the system is 3D

$$\begin{aligned}
 x' &= -3x + \sqrt{3/2}\epsilon\lambda y^2 + \frac{3}{2}x(1 + \epsilon x^2 - y^2) + \epsilon f(x, y), \\
 y' &= -\sqrt{3/2}\lambda xy + \frac{3}{2}y(1 + \epsilon x^2 - y^2), \\
 \lambda' &= -\sqrt{6}\lambda^2(\Gamma - 1)x
 \end{aligned}
 \tag{8}$$

Fixed Points	x	y	λ	$f(x, y)$
P_1	0	0	λ	0
P_{2i}	$\frac{3}{2}x(1 - \epsilon x^2)$ $-\epsilon f(x, 0) = 0$	0	0	$f(x, y) = 0$ for $x = 0, x \pm 1$
$P_{3\pm}$	0	± 1	$\lambda = -\sqrt{\frac{2}{3}}f(x, y)$	$f(x, y)$
$P_{4\pm}$	$3x - \epsilon f(x, y) = 0$	$y^2 = (1 + \epsilon x^2)$	0	$f(x, y)$

Stability

- P_1 is a completely matter dominated unstable fixed point.
- P_2 does not exist for the Phantom field and also unstable in nature. It could be an mixed state of both the matter and dark energy.
- P_3 is completely dark energy dominated fixed point.
- P_4 can have both dark energy and dark matter contribution and its stability depends on the choice of the interaction.

 P_1

$$\left[\frac{3}{2}, 0, \frac{1}{2}(2\epsilon\partial_x f - 3) \right]$$

 P_2

$$\left[0, \frac{3}{2}(x^2\epsilon + 1), \frac{1}{2}(2\epsilon\partial_x f + 9x^2\epsilon - 3) \right]$$

**All are non-hyperbolic
fixed points**

 P_3

$$\left[-3, 0, -3 + \epsilon\partial_x f \right]$$

 P_4

$$\left\{ 0, \frac{1}{6}(-A + 3\epsilon\partial_x f - 18), \frac{1}{6}(A + 3\epsilon\partial_x f - 18) \right\}$$

where

$$A = \sqrt{3}\sqrt{4f^2(\partial_x f - 3\epsilon) + 4\epsilon f\sqrt{9 + \epsilon f^2\partial_y f} + 3(\partial_x f)^2}.$$

Example

Let us consider a form of the $f(x,y)$ and check our general analysis;

$$f(x, y) = \alpha(1 - \epsilon x^2 - y^2)^m x^\gamma \quad (9)$$

Depending on the choice of the it could incorporate a large class of interaction.

Exponential Potential

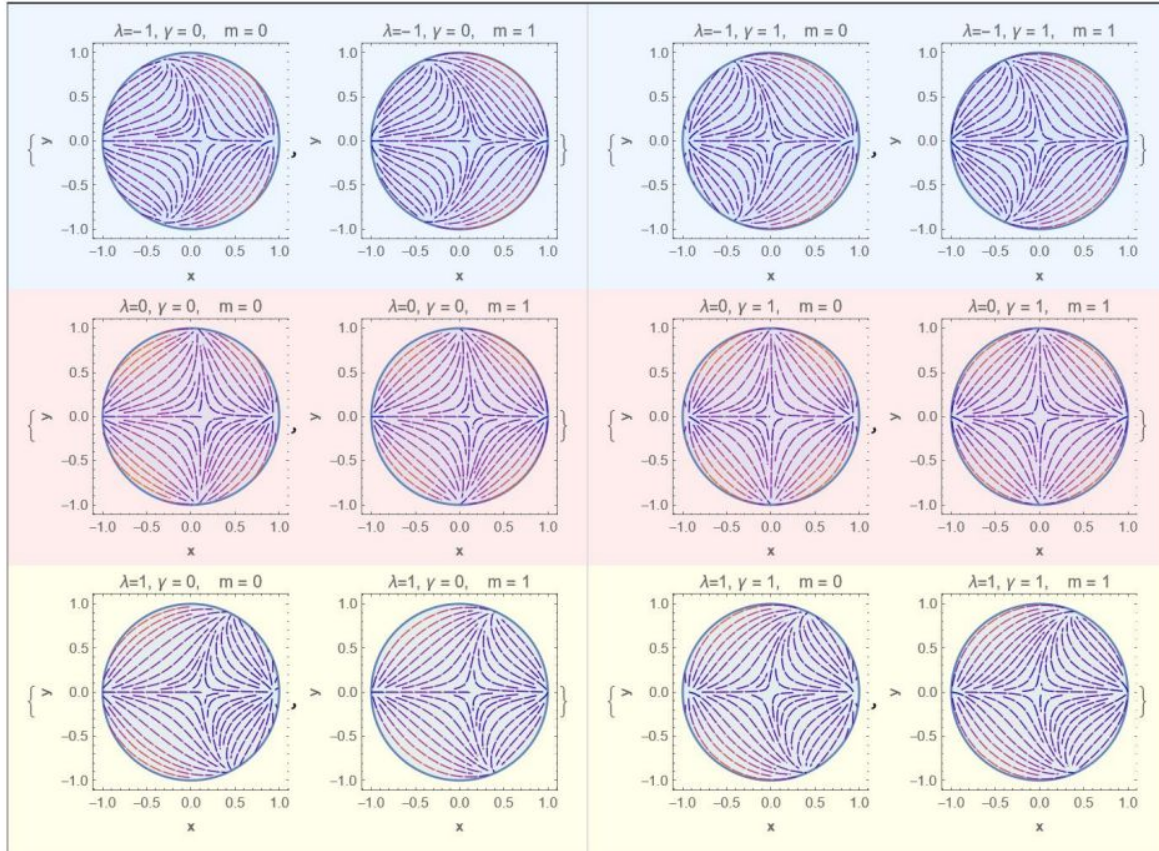


Fig.(3). Phase plot of the quintessence scalar field with exponential potential.

$$0 \leq \Omega_\phi \leq 1;$$

$$0 \leq x^2 + y^2 \leq 1.$$

Exponential Potential

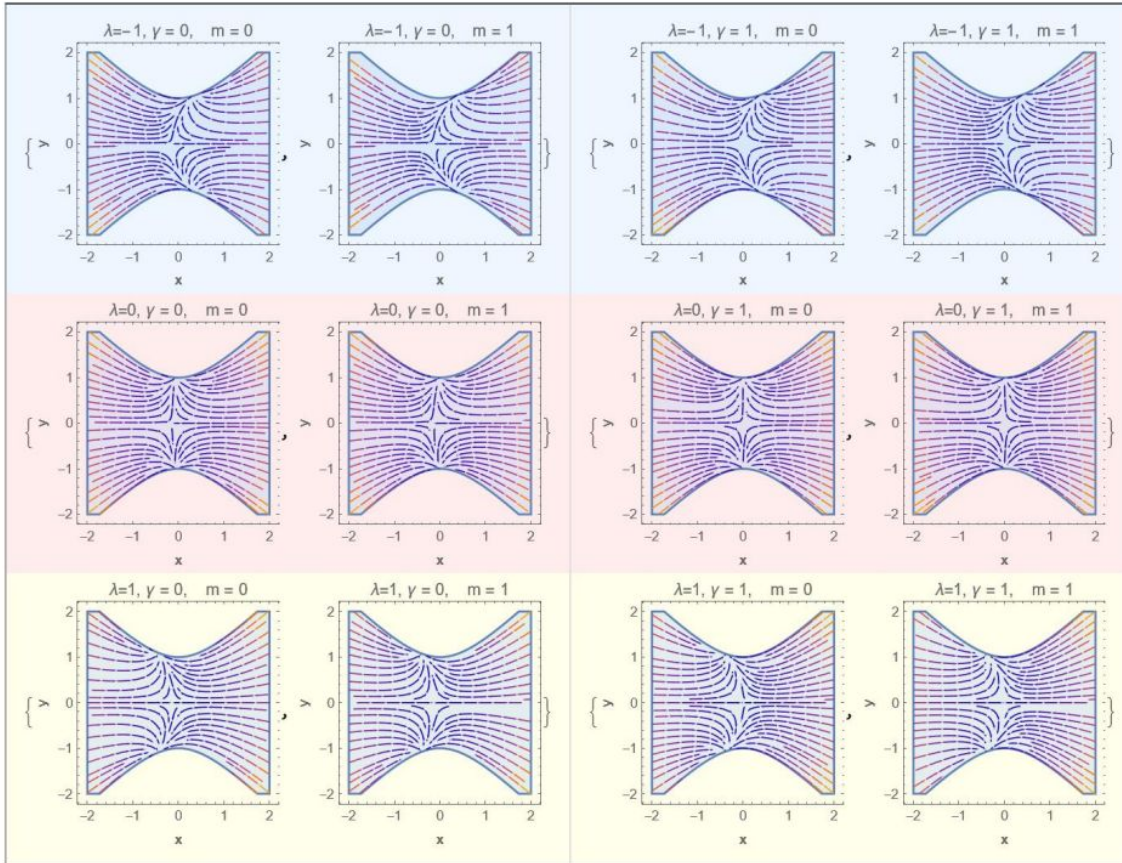


Fig.(4). Phase plot of the phantom scalar field with exponential potential.

$$0 \leq \Omega_\phi \leq 1;$$

$$0 \leq -x^2 + y^2 \leq 1.$$

Exponential Potential

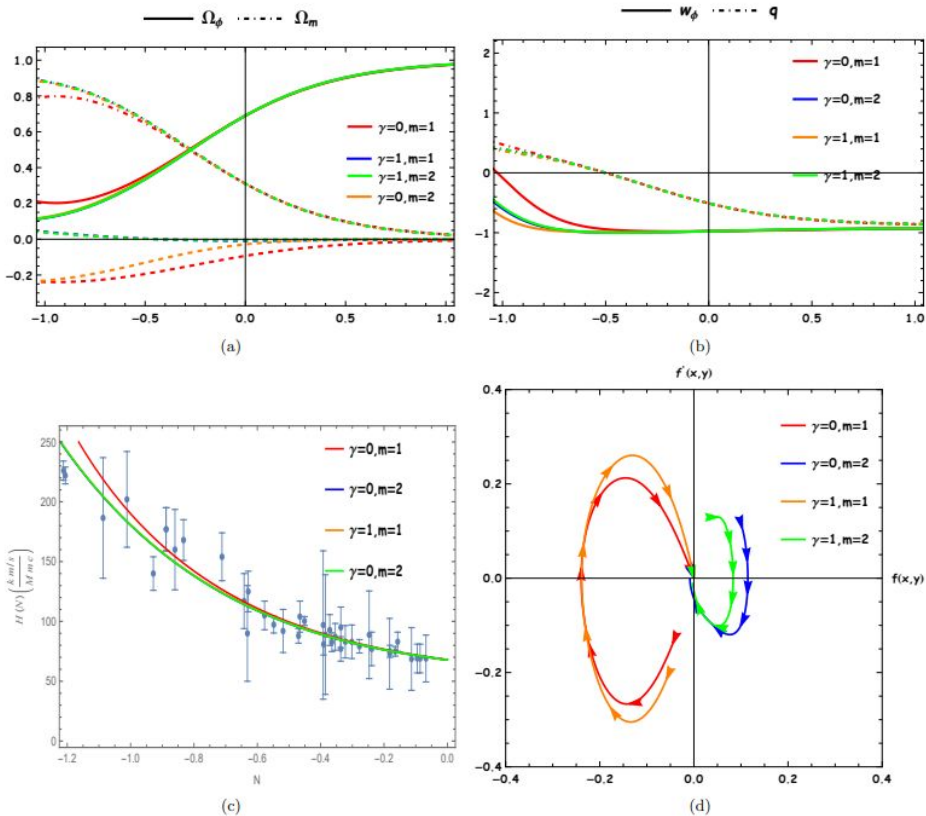


Fig.5, Plots of different cosmological parameters for the quintessence field for the exponential potential.

Exponential Potential

- We have estimated the current values of the x, y from the observation by solving,

$$\Omega_{\phi 0} = \epsilon x_0^2 + y_0^2$$

$$q_0 = -1 + \frac{3}{2}(1 + \epsilon x_0^2 - y_0^2)$$

- For quintessence field we estimate $x_0 = 0.09$, $y_0 = 0.825$ and for the phantom field $x_0 = 0.01$, $y_0 = 0.824$.

$$\Omega_{\phi 0} = 0.68$$

$$q_0 = -0.51$$

1. N. Aghanim et al. Planck 2018 results. VI. Cosmological parameters. 2018
2. Adam G. Riess et al. Astrophys. J. Lett., 934(1):L7, 2022

Exponential Potential

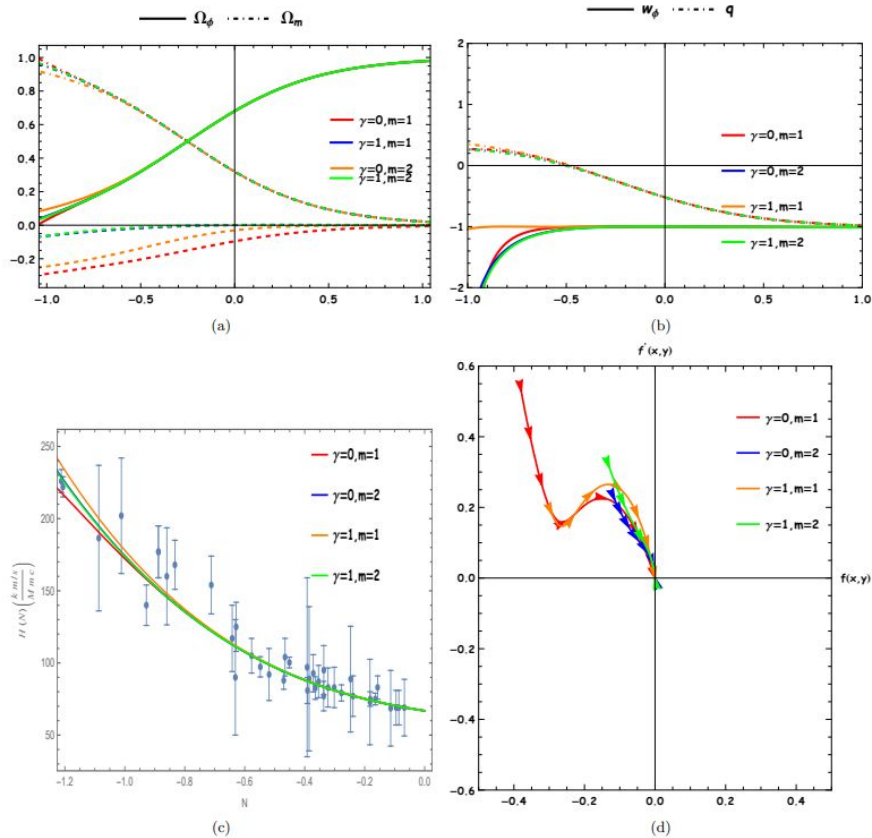


Fig.6, Plots of different cosmological parameters for the phantom field for the exponential potential.

Non-Exponential Potential

Γ could be a function of λ . Here we considered two different choices.

1. $\Gamma = 1/2$ for which the potential becomes to be $V(\phi) = (A + B\phi)^2$.
2. $\Gamma = \frac{1}{2}(\frac{1}{\lambda^2} - 1) + 1$, $V(\phi) = \cosh(\xi\phi) - 1$ where we chose $\xi = 1$.

$\Gamma = 1/2$

Fig.7, quintessence case

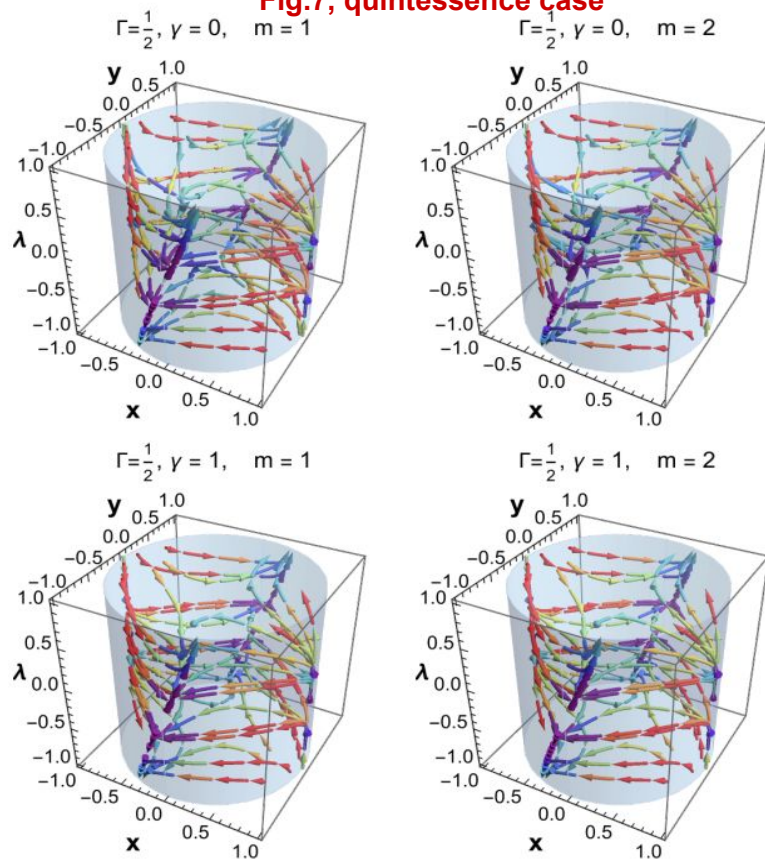
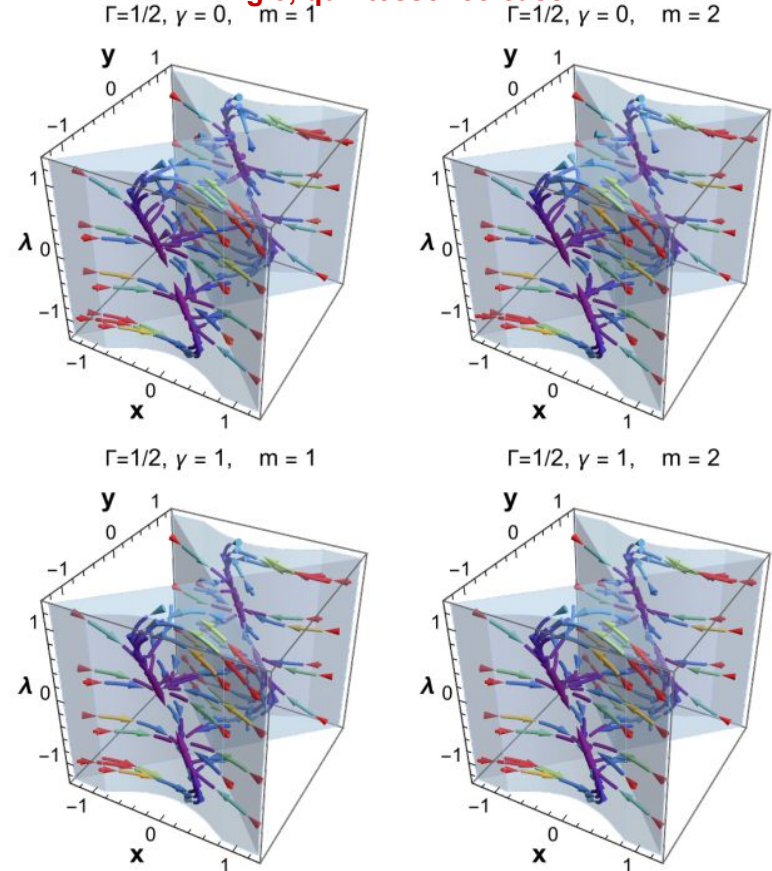


Fig.8, quintessence case



$\Gamma = 1/2$

17

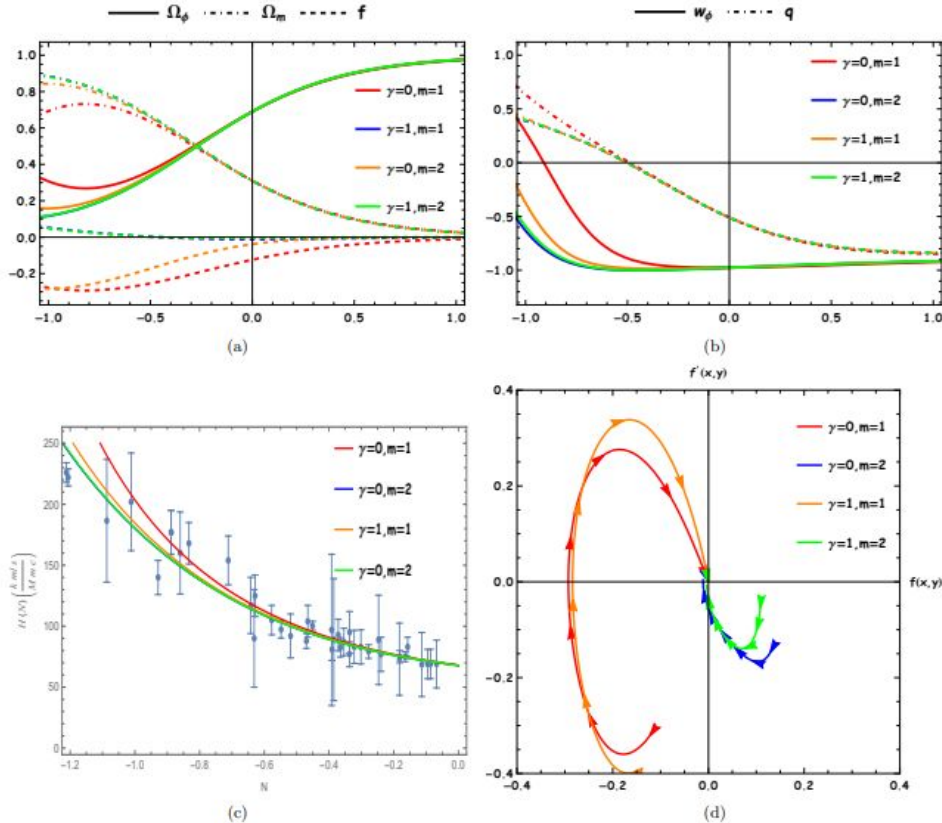


Fig.9, Plots of different cosmological parameters for the quintessence field for the exponential potential.

$$\Gamma = 1/2$$

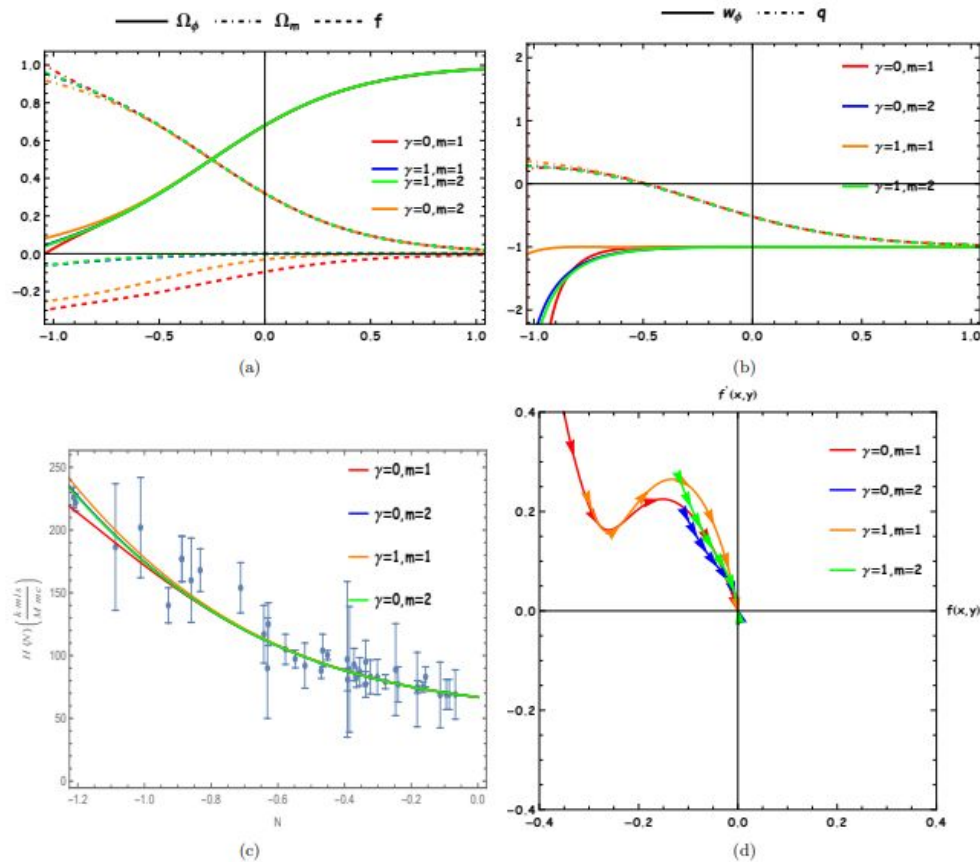
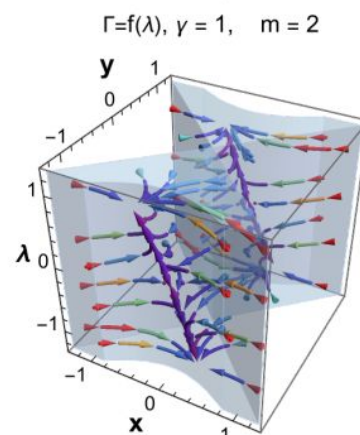
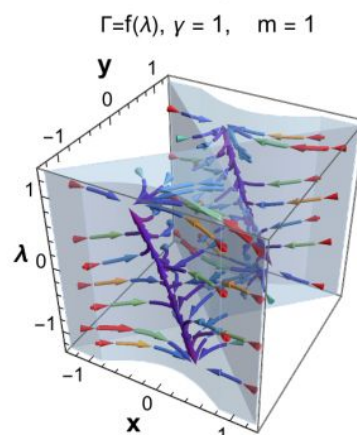
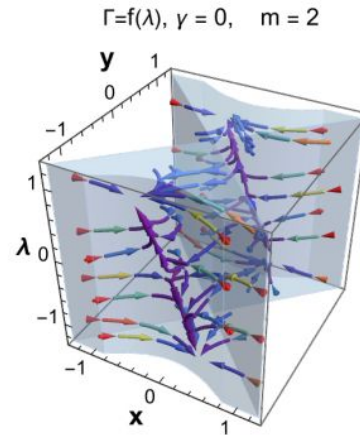
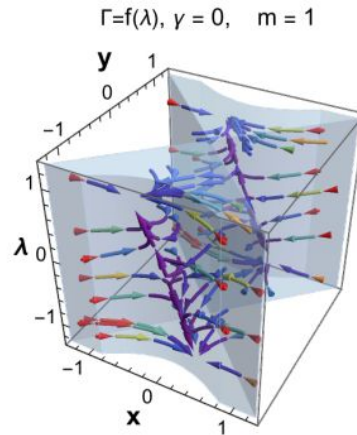
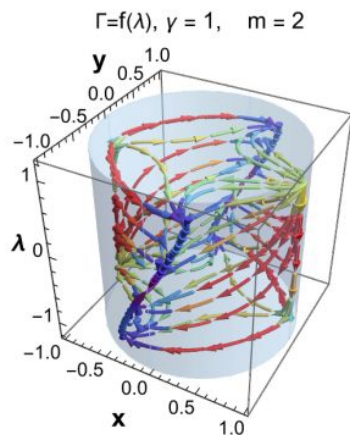
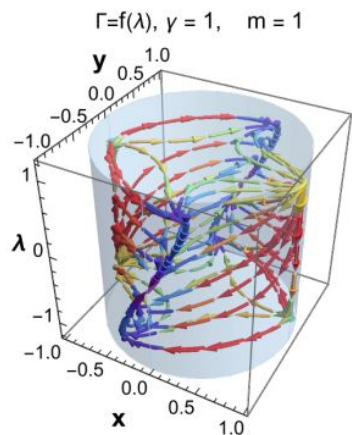
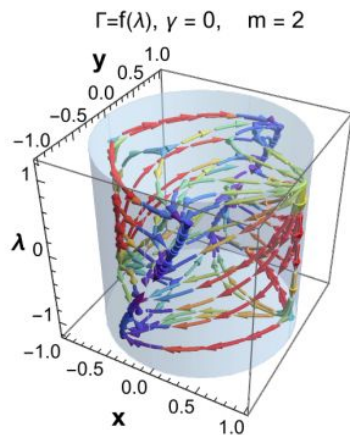
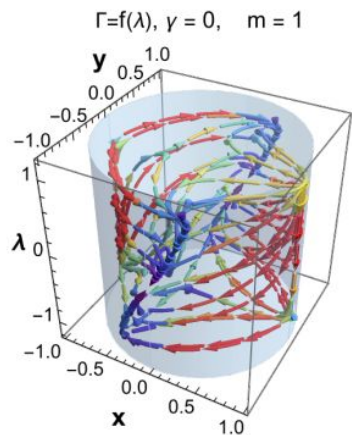


Fig.10, Plots of different cosmological parameters for the phantom field for the exponential potential.

$$\Gamma = \frac{1}{2} \left(\frac{1}{\lambda^2} - 1 \right) + 1$$



$$\Gamma = \frac{1}{2} \left(\frac{1}{\lambda^2} - 1 \right) + 1$$

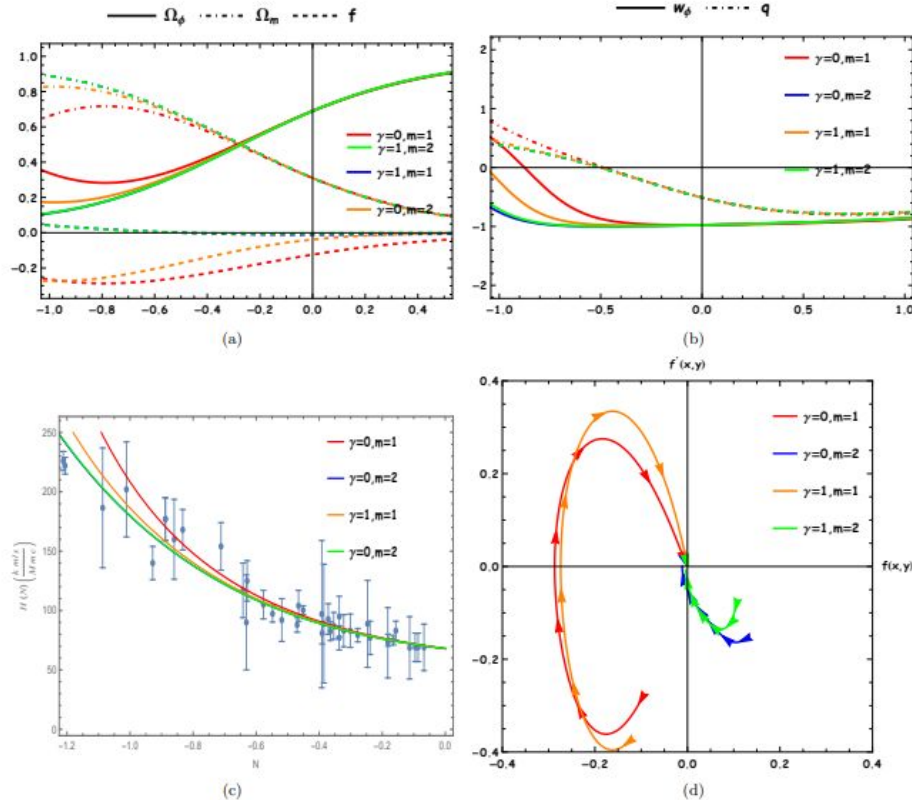


Fig.11, Plots of different cosmological parameters for the quintessence field for the non-exponential potential.

$$\Gamma = \frac{1}{2} \left(\frac{1}{\lambda^2} - 1 \right) + 1$$

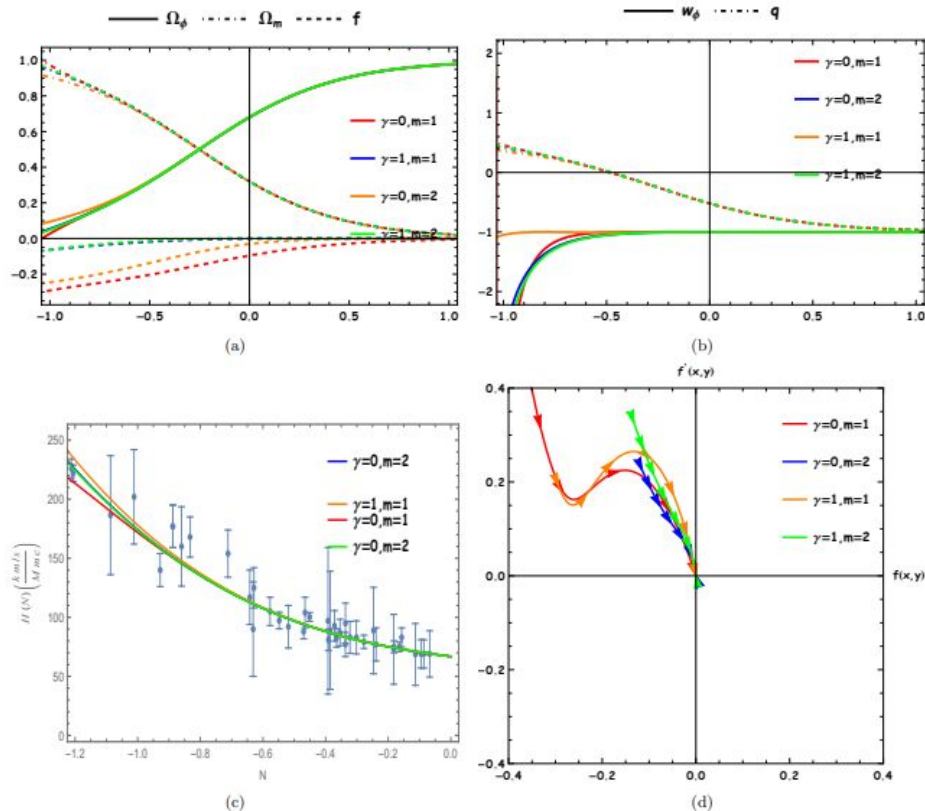


Fig.12, Plots of different cosmological parameters for the phantom field for the non-exponential potential.

Conclusion

- We have studied interacting scalar field model dark energy models using dynamical system analysis.
- In the first part the analysis is done considering a general form of the potential. Our setup includes both the quintessence and the phantom field.
- In the second part both we have considered a particular form of the interaction and a detail phase space behaviour of the system has been studied.
- Our analysis suggests that there could be interaction at present between the dark sectors and in future the universe is completely dominated by the dark energy.

Thank You