

# Scaling Solutions in Generalized Proca Theory and its Cosmological Implications

Santiago García-Serna

J. Bayron Orjuela-Quintana  
César A. Valenzuela-Toledo

Facultad de Ciencias Naturales y Exactas  
Universidad del Valle



# Motivation

Modified/Extended gravities

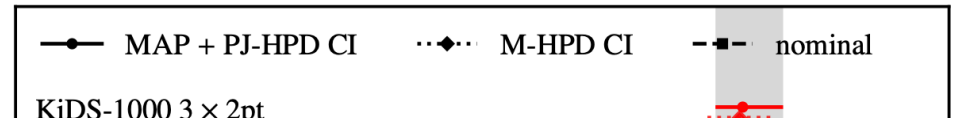
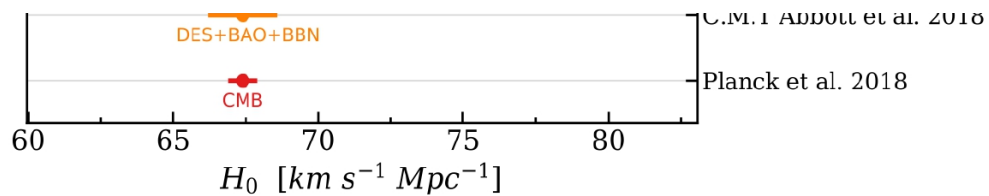
High-order derivatives theories

Extra fields theories

Non-local theories

Higher dimensions theories

Extensions: torsion/emergent theories

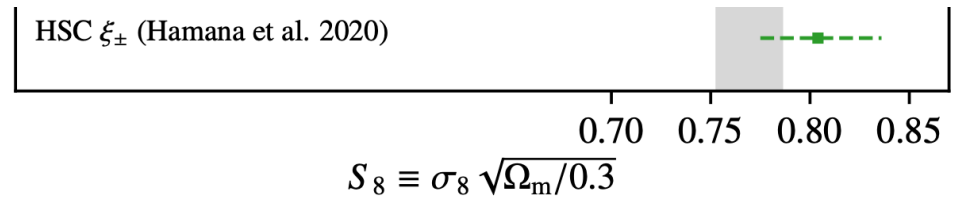


Prescription for dark energy  
dark energy

Standard models:  
non-minimally coupled  
to curvature

Clustering:  
scenarios with fluctuations  
on subhorizon scales

Modified gravity:  
presence of fifth force and  
anisotropic clustering



[arxiv.org/abs/2203.06142](https://arxiv.org/abs/2203.06142)

## **Dynamical Dark Energy**

arXiv:1701.08165  
arXiv:2201.09306

## **Modified Gravity models**

- **Horndeski theory  
(arXiv:1901.07183 )**

## **Interacting Dark Energy**

arXiv:1910.09853  
arXiv:Astro-ph/  
9908023.

What about the Generalized Proca theory coupled?

$$S = \int dx^4 \sqrt{-g} \left( \mathcal{L}_F + \sum_i \mathcal{L}_i + \mathcal{L}_M \right)$$

$$X = -\frac{1}{2} A_\mu A^\mu,$$

$$F = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu},$$

$$\tilde{F} = -\frac{1}{4} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$\mathcal{L}_F = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu}$$

$$\mathcal{L}_2 = G_2(X, F, \tilde{F})$$

$$\mathcal{L}_3 = G_3(X) \nabla_\mu A^\mu$$

$$\mathcal{L}_4 = G_4(X) R + G_{4X} \left[ (\nabla_\mu A^\mu)^2 + c_2 \nabla_\rho A_\sigma \nabla^\rho A^\sigma - (1 + c_2) \nabla_\rho A_\sigma \nabla^\sigma A^\rho \right]$$

$$\begin{aligned} \mathcal{L}_5 = & G_5(X) G_{\mu\nu} \nabla^\mu A^\nu - \frac{1}{6} G_{5X} \left[ (\nabla_\mu A^\mu)^3 \right. \\ & - 3d_2 \nabla_\mu A^\mu \nabla_\rho A_\sigma \nabla^\rho A^\sigma \\ & - 3(1 - d_2) \nabla_\mu A^\mu \nabla_\rho A_\sigma \nabla^\sigma A^\rho \\ & + (2 - 3d_2) \nabla_\rho A_\sigma \nabla^\gamma A^\rho \nabla^\sigma A_\gamma \\ & \left. + 3d_2 \nabla_\rho A_\sigma \nabla^\gamma A^\rho \nabla_\gamma A^\sigma \right]. \end{aligned}$$

# On the background of FLRW

$$ds^2 = -dt^2 + a^2(t)\delta_{ij}dx^i dx^j$$

$$A_\mu = (A_0(t), 0, 0, 0)$$

$$X = \frac{1}{2}A_0^2, \quad F = \tilde{F} = 0$$

Background

$$\mathcal{L}_F = 0$$

$$\mathcal{L}_2 = G_2$$

$$\mathcal{L}_3 = G_3 \dot{A}_0$$

$$\mathcal{L}_4 = \frac{m_{\text{P}}^2}{2} R \equiv \mathcal{L}_{\text{EH}}$$

$$\mathcal{L}_5 = \text{Boundary term}$$

$$ds^2 = a^2(\eta) [-d\eta^2 + \{\delta_{ij} + h_{ij}(x, \eta)\} dx^i dx^j]$$

$$c_T^2 = \frac{2G_4 - A_0^2 A_0' G_{5X}}{2G_4 - 2A_0^2 G_{4X} - H A_0^3 G_{5X}} = 1$$

$$G_4 = \frac{m_{\text{P}}^2}{2}, \quad G_5 = \text{const}^{\text{st}}$$

Tensor perturbations

$$\mathcal{L}_{\text{GP}} = G_2(X) + G_3(X) \nabla_\mu A^\mu$$

# Generalized Proca theory coupled to Cold Dark Matter and Scaling Solution

$$\mathcal{L} = \mathcal{L}_{\text{EH}} + \mathcal{L}_{\text{GP}} + f\mathcal{L}_{\text{CDM}} + \mathcal{L}_b + \mathcal{L}_r \quad \rho_{\text{DE}} = -G_2, \quad w \equiv \frac{p}{\rho}, \quad \Omega \equiv \frac{\rho}{3m_{\text{P}}^2 H^2}, \quad w_{\text{eff}} = \sum_i \Omega_i w_i$$

$$\frac{\rho_{\text{DE}}}{\rho_{\text{CDM}}} = \text{const}$$



$$\frac{d \ln \mathcal{L}_{\text{GP}}}{dX} X' = -3(1 + w_{\text{eff}})$$



$$\mathcal{L}_{\text{GP}} = G_2(X) = V_0 e^{-\gamma \int dX (f_X/f)}$$

Scaling condition

$$\begin{aligned} 3m_{\text{P}}^2 H^2 &= -G_2 + f\rho_{\text{CDM}} + \rho_b + \rho_r \\ -2m_{\text{P}}^2 \dot{H} &= f\rho_{\text{CDM}} + \rho_b + \frac{4}{3}\rho_r \end{aligned}$$

Friedmann equations

$$\begin{aligned} v^2 &= -\frac{G_2}{3m_{\text{P}}^2 H^2}, & \tilde{\Omega}_{\text{CDM}} &= \frac{f\rho_{\text{CDM}}}{3m_{\text{P}}^2 H^2} \\ \Omega_b &= \frac{\rho_b}{3m_{\text{P}}^2 H^2}, & \Omega_r &= \frac{\rho_r}{3m_{\text{P}}^2 H^2} \end{aligned}$$


$$\left( G_{2X} - \frac{f_X}{f} \rho_{\text{CDM}} \right) A_0 = 0$$

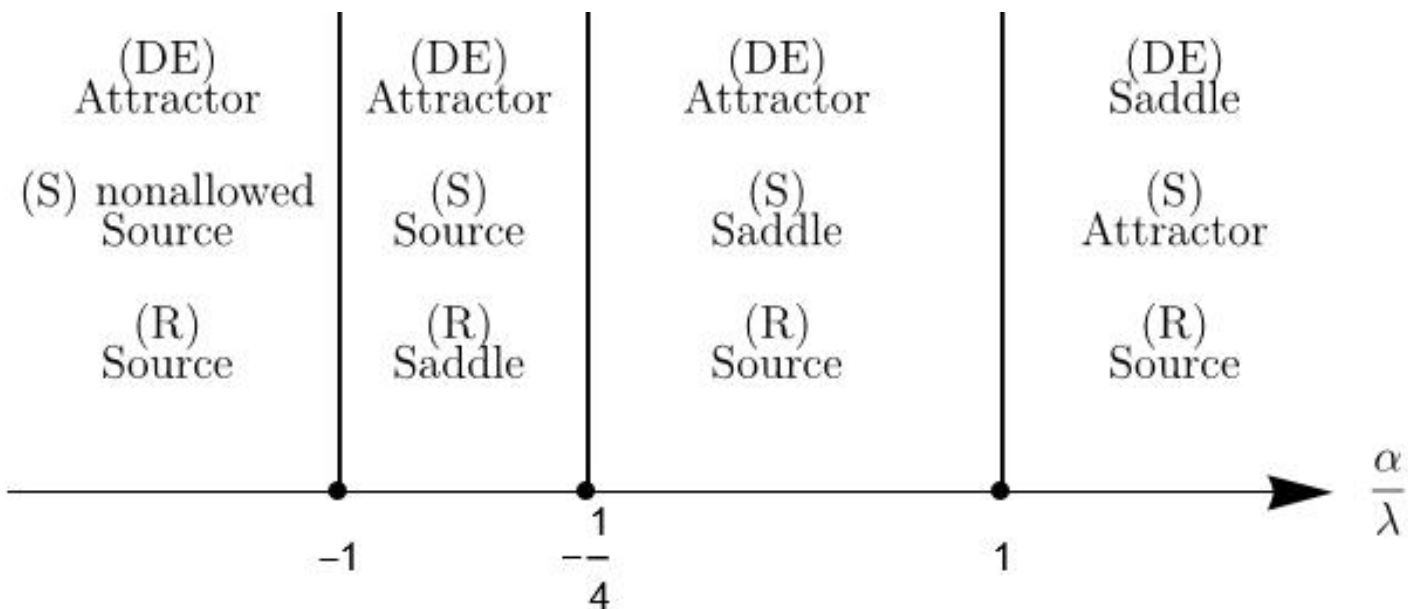
Motion equation of  $A_0$

$$f(X) = f_0 e^{\alpha X/m_{\text{P}}^2}, \quad G_2(X) = V_0 e^{-\lambda X/m_{\text{P}}^2}$$

# Fixed points and stability

Fixed Point	$v$	$\tilde{\Omega}_{\text{CDM}}$	$\Omega_b$	$\Omega_b$	$w_{\text{eff}}$	Stability
(R)	0	0	0	1	1/3	source or saddle
(M)	0	$1 - \Omega_b$	$\Omega_b$	0	0	saddle
(DE)	1	0	0	0	-1	saddle or attractor
(S)	$\sqrt{\frac{a}{\alpha+\lambda}}$	$1 - \frac{a}{\alpha+\lambda}$	0	0	$-\frac{a}{\alpha+\lambda}$	can be anyone

$$\frac{\alpha}{\lambda} > \frac{1}{2} \rightarrow w_{\text{eff}} < -\frac{1}{3}$$


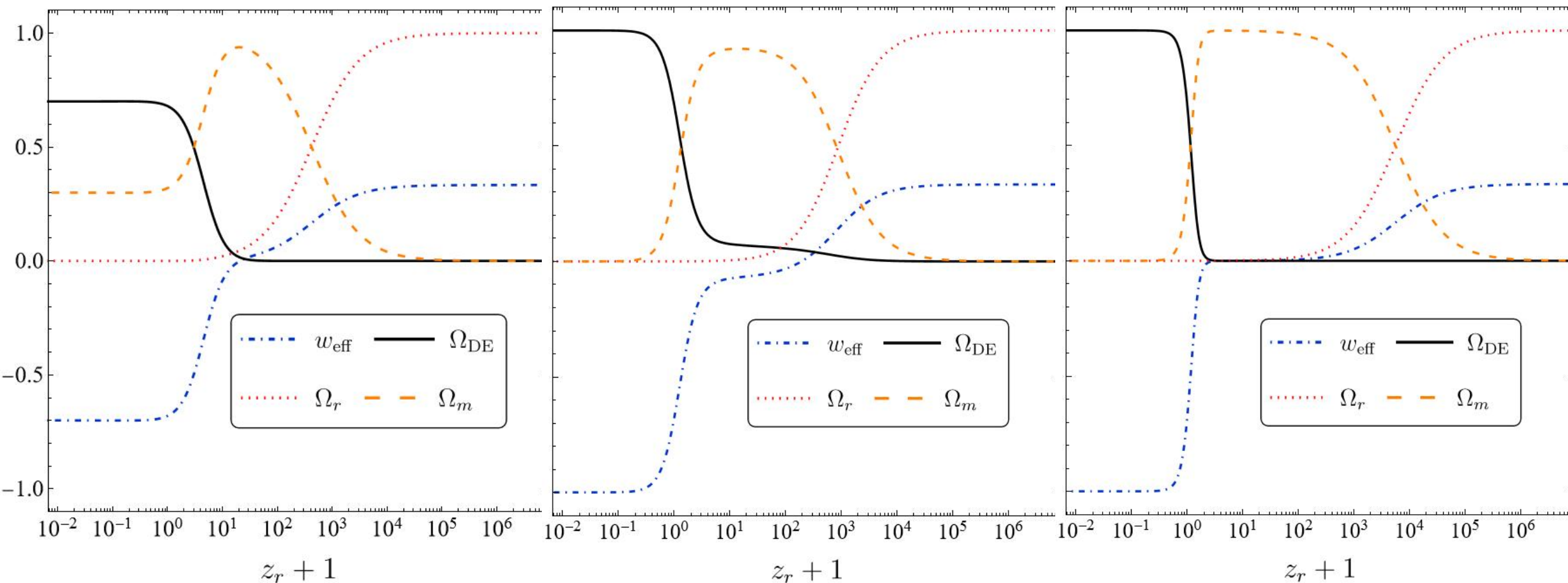


# Evolution of the background

(S)  $\frac{\alpha}{\lambda} = \frac{7}{3}$

(DE) with (S) as saddle  $\frac{\alpha}{\lambda} = 0.1$

(DE) De Sitter  $\frac{\alpha}{\lambda} = -0.7$





$$ds^2 = a(\eta)^2 [-\{1 + 2\Phi(\mathbf{x}, \eta)\}d\eta^2 + \{1 - 2\Psi(\mathbf{x}, \eta)\}d\mathbf{x}^2]$$

$$\tilde{\delta}_c \equiv \frac{\delta(f\rho_c)}{f\rho_c}$$

$$\tilde{\delta}'_c = -V - 3\Phi'$$

$$\Psi \rightarrow -\Phi$$

$$V' = -\mathcal{H}V + \frac{f_X(-A_0A'_0 + A_0^2\mathcal{H})}{a^2f}V + k^2\Psi$$

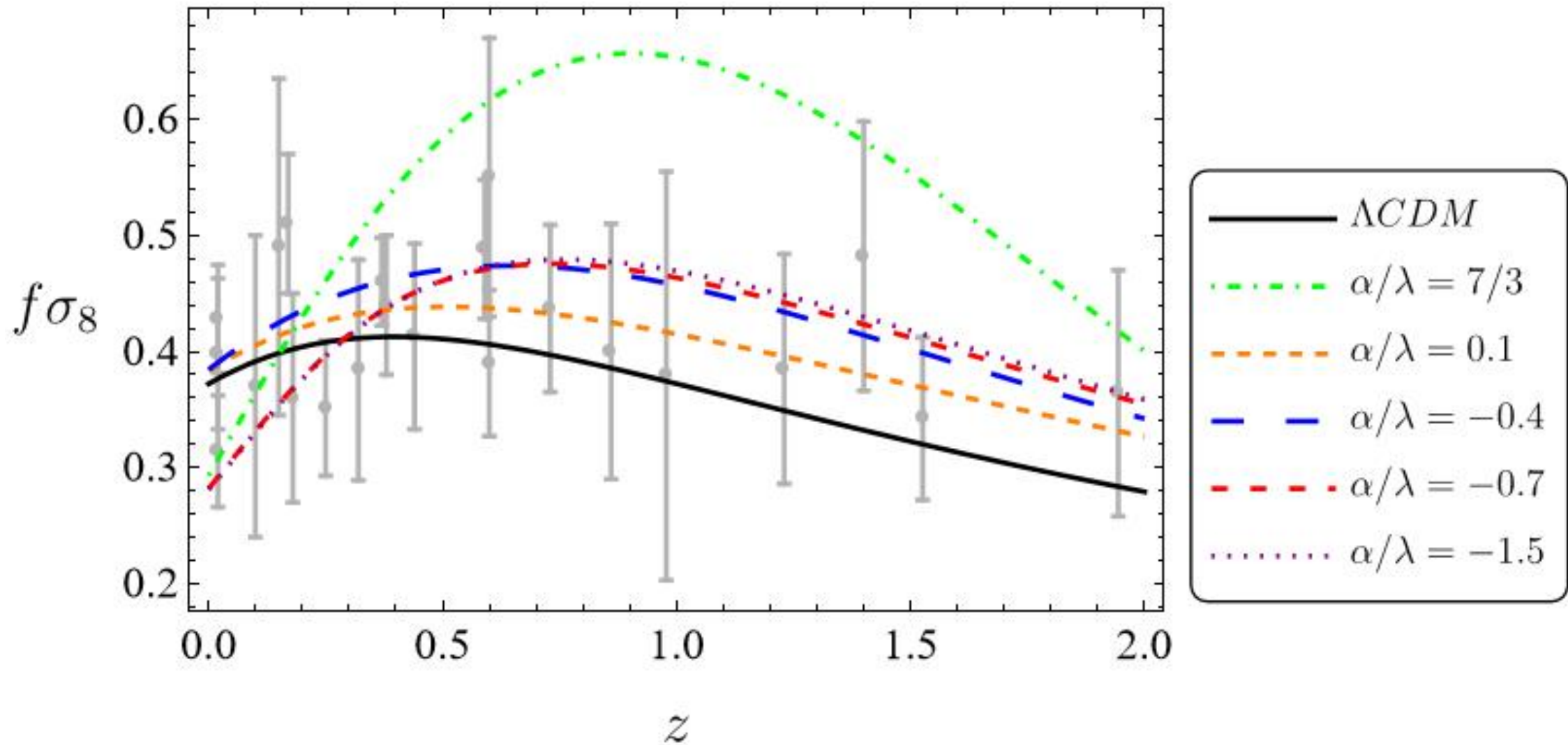
$$\Phi \rightarrow \frac{3a^2}{2k^4m_P^2}f\rho_c + \frac{a^2\tilde{\delta}_cf\rho_c}{2k^2m_P^2} - \frac{A_0^2f_X\tilde{\delta}_cf\rho_c}{2k^2f}$$

$$\Phi'' \rightarrow 0, \quad \Phi' \rightarrow 0$$

QSA-SHA approximation

$$\tilde{\delta}''_c(a) + \left( \frac{d \ln(a^3 H)}{da} + \alpha X'(a) \right) \tilde{\delta}'_c - \frac{3\Omega_{m0}G_{\text{eff}}/G_N}{2a^5 H^2(a)/H_0^2} \tilde{\delta}_c^0 = 0$$

$$G_{\text{eff}} \equiv (1 - 2X(a)\alpha)G_N$$



## Conclusion

- The interaction between DE and CDM allows for modification of the gravity effects during the matter domination
- Despite the interactions is possible to reproduce the evolution at the background level

## Prospects

- MCMC
- Does this model fix  $H_0$  and  $\sigma_8$  or at least one?
- How is this model compared with the others concerning the predicted observables?

