

# Testing Spacetime Symmetries in the Early Universe

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Cosmology from Home 2023

# Symmetries?

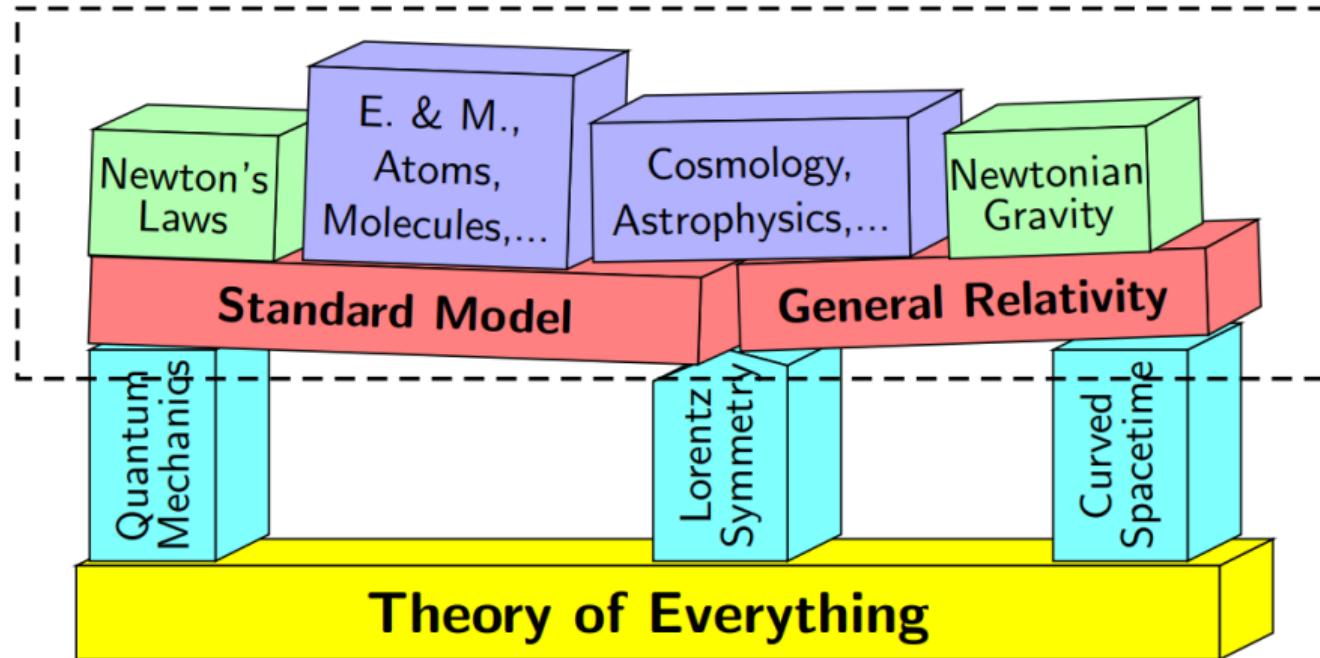


Image credit: Matt Mewes, Cal Poly

# Spacetime symmetries

- General relativity defined in geometrical terms:  $ds^2 = g_{\mu\nu}dx^\mu dx^\nu$ , no torsion

$$R^{\mu\nu} - \frac{1}{2}g^{\mu\nu}R = 8\pi G T^{\mu\nu}$$

• Geometry →  $R^{\mu\nu} - \frac{1}{2}g^{\mu\nu}R$   
• Matter →  $T^{\mu\nu}$

The related conservation laws lead to the Bianchi identities ( $+T_{[\mu\nu]} = 0$ ):

$$\nabla_\mu G^{\mu\nu} = 0$$

- Related symmetries:
  - Local Lorentz symmetry (**6 rotations**):  $B'^\mu = \Lambda_\nu^\mu B^\nu$
  - Diffeomorphism symmetry (**4 translations**):  $B^\mu \rightarrow B^\mu + (\partial_\nu \xi^\mu) B^\nu - \xi^\nu \partial_\nu B^\mu$

# Why test spacetime symmetries?

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## Why are we interested in this?

- **Theoretical motivation**

- Spontaneous Lorentz violation from string field theory (Kostelecký and Samuel, PRD 1989)
- Trickle-down UV Lorentz violation in LQG (Gambini et al, CQG 2011)
- High-order spatial derivatives for renormalisability (Hořava, PRD 2009)

- **What can be measured with precision?**

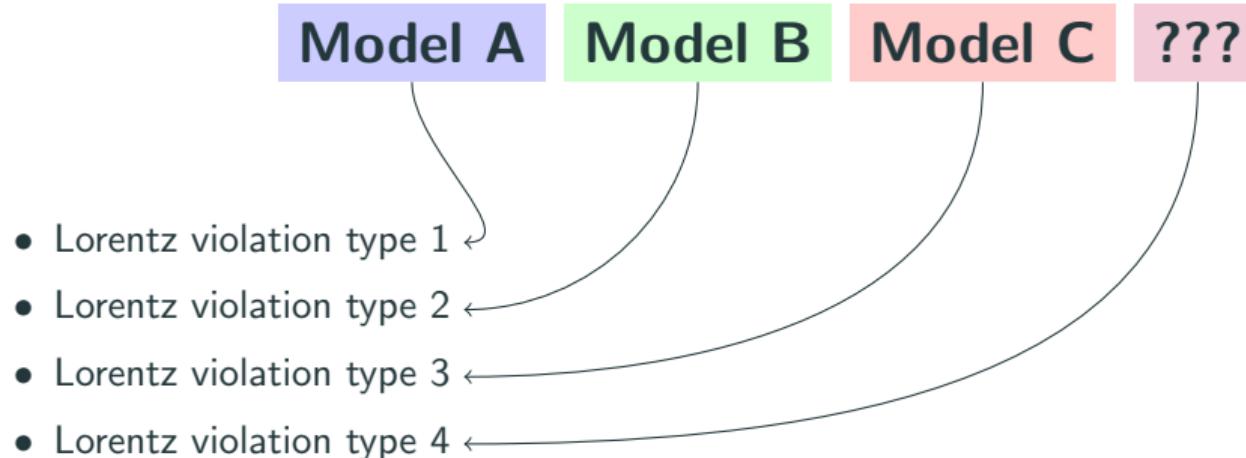
- Symmetries lend themselves easily to high precision tests, easy to describe theoretically

- **Quantum gravity likely alters the structure of spacetime**

- Are there extra spacetime dimensions? (Kaluza-Klein etc.)
- Do spacetime coordinates commute? ( $[x^\mu, x^\nu] = i\theta^{\mu\nu}$ )
- Is spacetime discrete? (LQG etc.)
- Spacetime foam?

[Spacetime symmetries are the basis of ALL of modern physics, we should be testing them!]

# The Standard-Model Extension



Can we find a general framework instead? (Answer: yes, we can!)

# Standard-Model Extension

- The Standard-Model Extension (SME): general framework for studying Lorentz/CPT violation, in all sectors

## Contains:

- Standard Model of particle physics
- General relativity
- ALL operator breaking Lorentz and/or CPT symmetry**

$$\mathcal{L}_{\text{SME}} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{EH}} + \mathcal{L}_{\text{LV}}$$

- operators proportional to *coefficients of Lorentz violation*, e.g.  $(k^{(4)})_{\mu\nu\alpha\beta} R^{\mu\nu\alpha\beta}$
- 20 d.o.f

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Originally developed by Kostelecky, Colladay + collaborators in the 1990's

# Standard-Model Extension

The SME is a **framework**, not a model! Can always map model parameters onto SME coefficients

Model	Link to SME	LV fields	General Test Framework
PPN	Yes	None ( $\alpha_1, \alpha_2, \dots$ )	Yes
SME gravity sector	Yes	Tensors	Yes, EFT
Bumblebee	Yes	Vectors	No
Einstein-Aether	Partial	Vector	No
Hořava gravity	Yes	Vector	No
ATT model	Yes	Antisymm. 2-tensor	No
Cardinal model	Yes	Symm. 2-tensor	No
Massive gravity	Yes	2-tensor	?
CS gravity	?	Scalar	No
GW MDR	Yes	None	Yes, DR
NC gravity	Yes	$\theta^{\mu\nu}$	No

# The pure-gravity sector

Pure gravity: generic Lagrangian up to dimension-6 operators

$$\mathcal{L} = \frac{\sqrt{-g}}{2\kappa} \left[ R + \underbrace{\left( k^{(4)} \right)_{\alpha\beta\gamma\delta} R^{\alpha\beta\gamma\delta}}_{20} + \underbrace{\left( k^{(5)} \right)_{\alpha\beta\gamma\delta\kappa} \nabla^\kappa R^{\alpha\beta\gamma\delta}}_{60} \right. \\ \left. + \underbrace{\frac{1}{2} \left( k^{(6)} \right)_{\alpha\beta\gamma\delta\kappa\lambda} \{ \nabla^\kappa, \nabla^\lambda \} R^{\alpha\beta\gamma\delta}}_{126} + \underbrace{\left( k^{(6)} \right)_{\alpha\beta\gamma\delta\kappa\lambda\mu\nu} R^{\alpha\beta\gamma\delta} R^{\kappa\lambda\mu\nu}}_{210} \right] + \mathcal{L}'$$

- $\nabla_\mu k \neq 0$ , thanks to property of Riemannian geometry
- Dynamical coefficients iff spontaneous breaking, contained in  $\mathcal{L}'$
- Any spontaneous breaking will involve a choice

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Altschul+Bailey PRD 2010 consider generic potentials for antisymmetric tensors. For the pure-gravity sector, see e.g. Kostelecky PRD 2004, Bailey+Kostelecky PRD 2006, O'Neal-Ault+Bailey+NAN PRD 2021

# Strong-field regime

- Bumblebee: cosmological anisotropies, superradiance in Kerr, noncommutative spacetime, monopoles (Jiang et al 2021, Maluf et al 2021, Kumar Jha et al JCAP 2020,...)<sup>a</sup>
- Cosmology: Inflation, deceleration, modified matter evolution in FLRW, Hubble tension, early Universe (Bonder PRD 2017, 2020, NAN 2018, Ault-O'Neal, Bailey, NAN PRD 2021, Reyes et al PRD 2021, 2022)<sup>a</sup>

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<sup>a</sup>and many more...

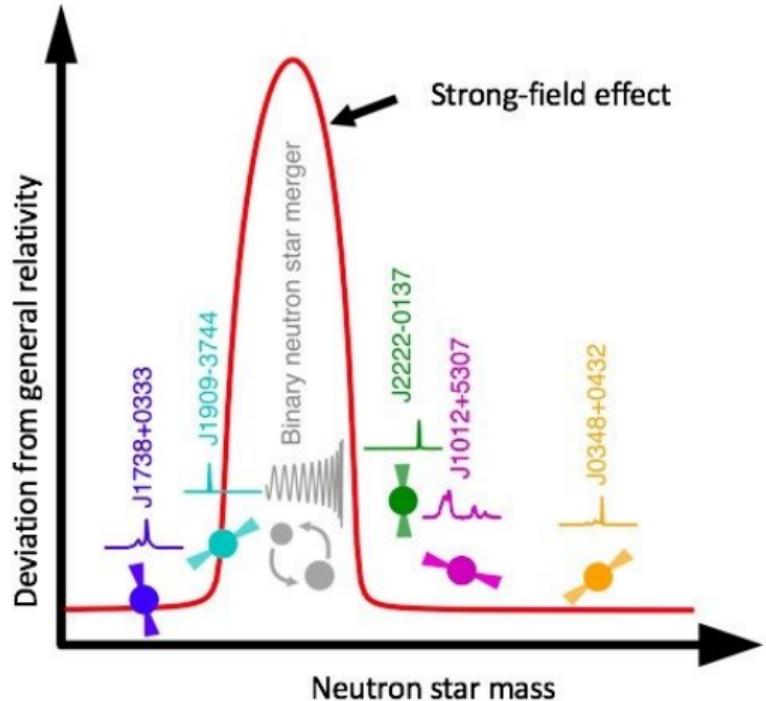


Figure 1: Image credit: L. Shao et al

## Cosmological solutions - Global coefficients only

- First cosmological solution within the SME
- Curvature term
- Energy density

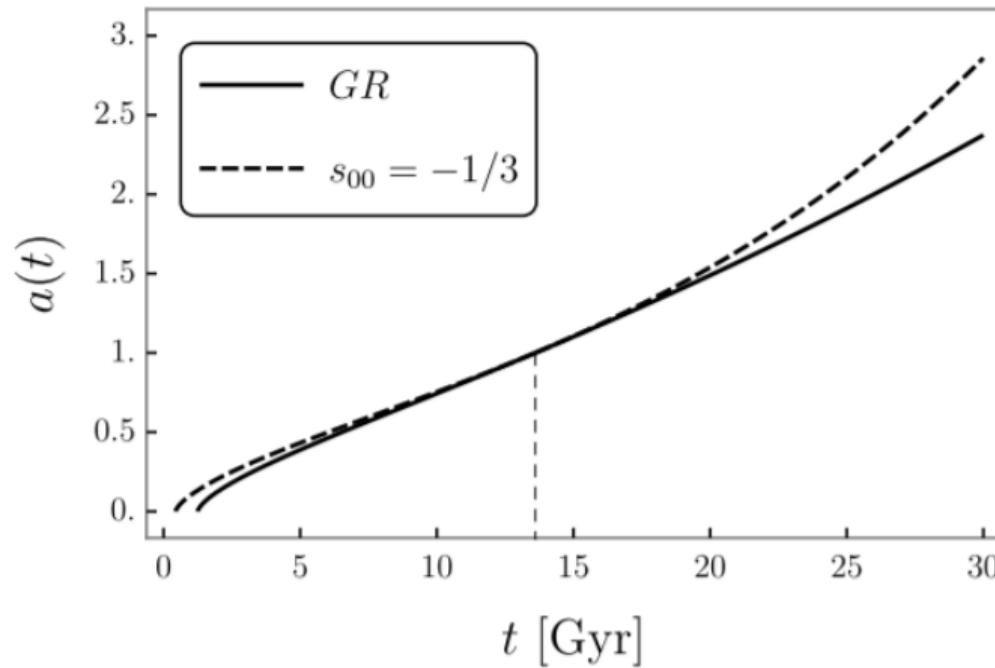
$$\left(\frac{\dot{a}}{a}\right)^2 (1 - s_{00}) = \frac{\kappa \rho}{3} - \frac{k}{a^2} - s_{00} \frac{\ddot{a}}{a} + \frac{\dot{a}}{a} \frac{\dot{s}_{00}}{2},$$

$$\left[\frac{\ddot{a}}{a} + \frac{1}{2} \left(\frac{\dot{a}}{a}\right)^2\right] (1 - s_{00}) = -\frac{\kappa p}{2} - \frac{k}{2a^2} + \frac{\dot{a}}{a} \dot{s}_{00} + \frac{1}{4} \ddot{s}_{00},$$

- Bianchi identities:  $\nabla_\mu (T_s)^{\mu\nu} = -\kappa \nabla_\mu (T_M)^{\mu\nu}$ .

$$\boxed{\nabla_\mu (T_s)^{\mu\nu} = \nabla_\mu (T_M)^{\mu\nu} = 0 \quad \text{OR} \quad \nabla_\mu [(T_s)^{\mu\nu} + \kappa (T_M)^{\mu\nu}] = 0}$$

$$\dot{\rho} + 3\frac{\dot{a}}{a}f(w, s_{00})\rho = 0, \quad f(w, s_{00}) = \frac{2(1+w-s_{00})}{2+s_{00}(3w-2)}$$



# Existing limits

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What are the existing limits on  $s^{\mu\nu}$ ?

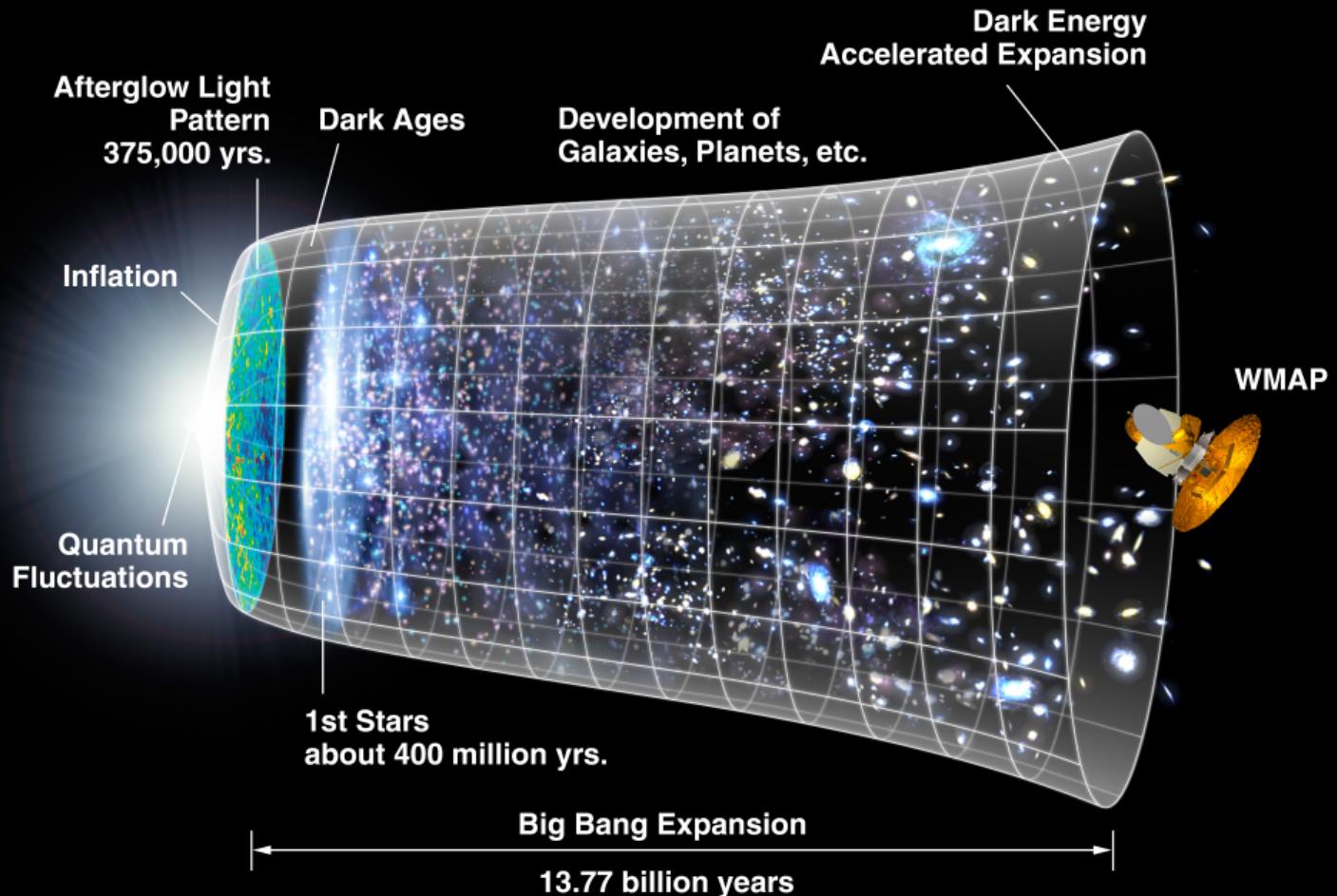
- Choose limits on  $s^{00}$  only, as relevant to this work.
- No limits in the exact sector exist! Only linearised gravity, i.e.  $\bar{s}^{00}$

Gravitational waves:

- Speed of GW's:  $\bar{s}^{00} = -0.2 - 0.07$  (Liu, He PRD 2020)
- GW170817 and GRB 170817A:  $\bar{s}^{00} = (-20 - 2.5) \cdot 10^{-15}$  (LIGO ApJL 2017)
- Gravitational Cherenkov radiation (cosmic rays):  $\bar{s}^{00} > -3 \cdot 10^{-15}$  (Kostelecky & Tasson PLB 2015)

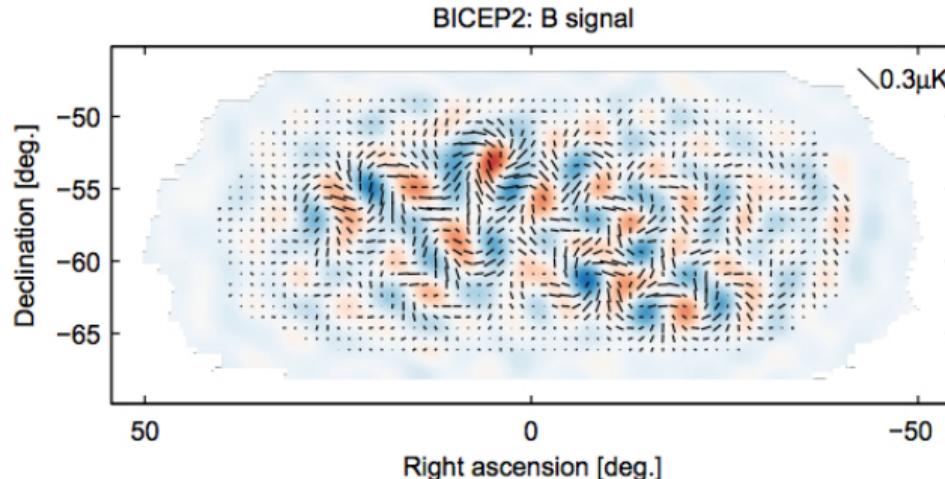
Can we do better?

Answer: Yes, no, and maybe...



# Spacetime-symmetry breaking

- The primordial Universe is an excellent probe of fundamental physics
- What happens to primordial gravitational waves if symmetries are broken?
- Primordial gravitational waves produced by quantum fluctuations during inflation; One possible signal, CMB B modes (**not yet observed**)



# Generating PGW's with spacetime-symmetry breaking

- Introduce tensor perturbations around a flat FLRW background

$$ds^2 = -dt^2 + a(t)(\delta_{ij} + h_{ij}(t, \mathbf{x}))dx^i dx^j$$

- Transverse-traceless perturbation, these are the gravitational waves

$$\delta^2 S = \frac{1}{4} \int dt d^3x a \left[ a^2 (1 - s_{00}) \dot{h}_{ij} \dot{h}^{ij} - \partial_k h_{ij} \partial^k h^{ij} \right]$$

- Equations of motion for the perturbations

$$\ddot{h}_{ij} + \left( 3H - \frac{\dot{s}_{00}}{1 - s_{00}} \right) \dot{h}_{ij} - \frac{1}{a^2(1 - s_{00})} \nabla^2 h_{ij} = 0$$

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See NAN PRD 2022 for all details

See also Bonder PRD 2017 for a discussion about inflation and spontaneous breaking with a Weyl-symmetric 4-tensor.

## Generating PGW's with spacetime-symmetry breaking

- Move to Fourier space using

$$h_{ij}(t, \mathbf{x}) = \sum_{\lambda} \int \frac{d^3 k}{(2\pi)^{3/2}} \epsilon_{ij}^{\lambda}(\mathbf{k}) k_{\mathbf{k}}(t) e^{i\mathbf{k}\cdot\mathbf{x}} \implies \ddot{h}_{\mathbf{k}} + \left( 3H - \frac{\dot{s}_{00}}{1-s_{00}} \right) \dot{h}_{\mathbf{k}} + \frac{k^2}{a^2(1-s_{00})} h_{\mathbf{k}} = 0$$

- From here we can read off that

$$c_T^2 = \frac{1}{1-s_{00}}, \quad m_g = 0, \quad \frac{d \ln M_{Pl}^2}{dt} = -\frac{\dot{s}_{00}}{1-s_{00}}.$$

- Propagation of tensor modes ("speed of gravity") constrained with GW170817+GRB170817A to  $-3 \cdot 10^{-15} < c_T - 1 < +7 \cdot 10^{-16}$ , leading to

$$-6 \cdot 10^{-15} < s_{00} < +1.4 \cdot 10^{-15}$$

## Background evolution

- We need to plug in the background evolution, which is also modified

$$\frac{H^2}{H_0^2} = \Omega_m^0 a^{-3} + \Omega_r^0 a^{-4x_r} + \Omega_\Lambda^0 a^{-x_\Lambda} + \Omega_k^0 a^{-2}$$

- Evolution of radiation and the cosmological constant (!) different from  $\Lambda$ CDM
- As a result, there is **no pure de Sitter phase**. In conformal time, it reads

$$\mathcal{H}_{\text{dS}} = -\frac{2}{(2-x_\Lambda)\eta}, \quad \mathcal{H}_{\text{RD}} = \frac{1}{(2x_r-1)\eta}, \quad \mathcal{H}_{\text{MD}} = \frac{1}{2\eta}$$

## Quantising the perturbations

- de Sitter inflation: simple model with exponential expansion and **scale-invariant** power spectrum
- To find the power spectrum, we need quantisation.

$$\hat{h}_{\mathbf{k}}(\eta) = v_{\mathbf{k}}(\eta)\hat{a}_{\mathbf{k}} + v_{\mathbf{k}}^*(\eta)\hat{a}_{\mathbf{k}}^\dagger, \quad \tilde{h}_{\mathbf{k}} = ah_{\mathbf{k}}$$

- Mode coefficients  $v_{\mathbf{k}}(\eta)$  satisfy the **Mukhanov-Sasaki equation**

$$v''_{\mathbf{k}} + \left( \frac{k}{1 - s_{00}} - \frac{2}{\eta^2} \gamma \right) v_{\mathbf{k}} = 0, \quad (s_{00} \rightarrow 0 \Rightarrow \gamma \rightarrow 1 \Rightarrow \text{GR})$$

- Not possible to rescale away  $s_{00}$ 's at both background and perturbation level

# The power spectrum

- Now we need the two-point correlation function:

$$\langle 0 | \hat{h}_k^\dagger \hat{h}_k | 0 \rangle \equiv \frac{16\pi G}{a^2} \int_0^\infty \frac{dk}{k} \frac{k^3}{2\pi^2} |\nu_k|^2,$$

- and the **power spectrum** is the logarithmic derivative of the correlator

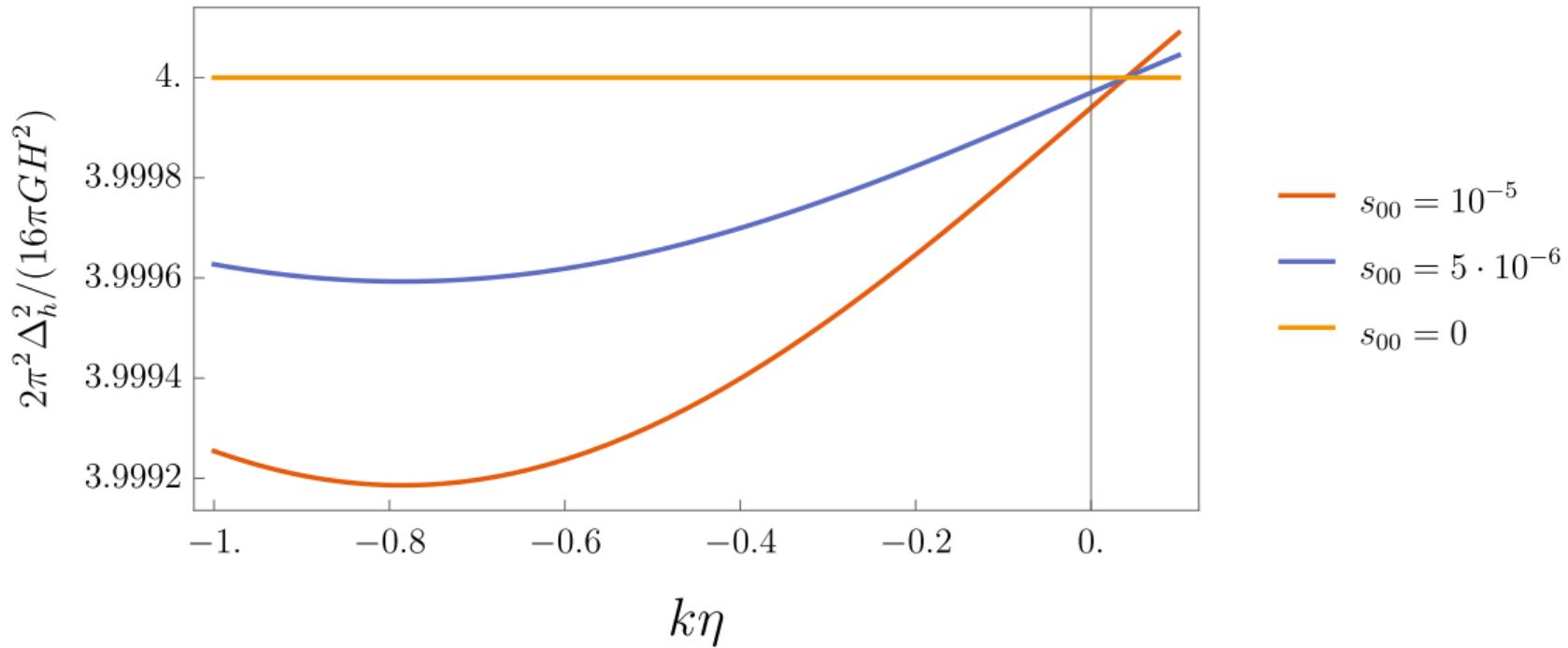
$$\Delta_h^2(\eta, k) \equiv \frac{d \langle 0 | \hat{h}_k^\dagger \hat{h}_k | 0 \rangle}{d \ln k} = \frac{16\pi G}{a^2} \frac{k^3}{2\pi^2} |\nu_k|^2 = \frac{16\pi G}{a^2} \frac{k^3}{2\pi^2} H^2 (2 - x_\Lambda)^2 \gamma^2 |A_k(\eta)|^2,$$

$$A_k(\eta) = e^{\frac{2ik\eta}{1-s_{00}}} \left( 1 + e^{i\pi\sqrt{1+8\gamma}} \right) - ie^{\frac{i\pi}{2}\sqrt{1+8\gamma}}$$

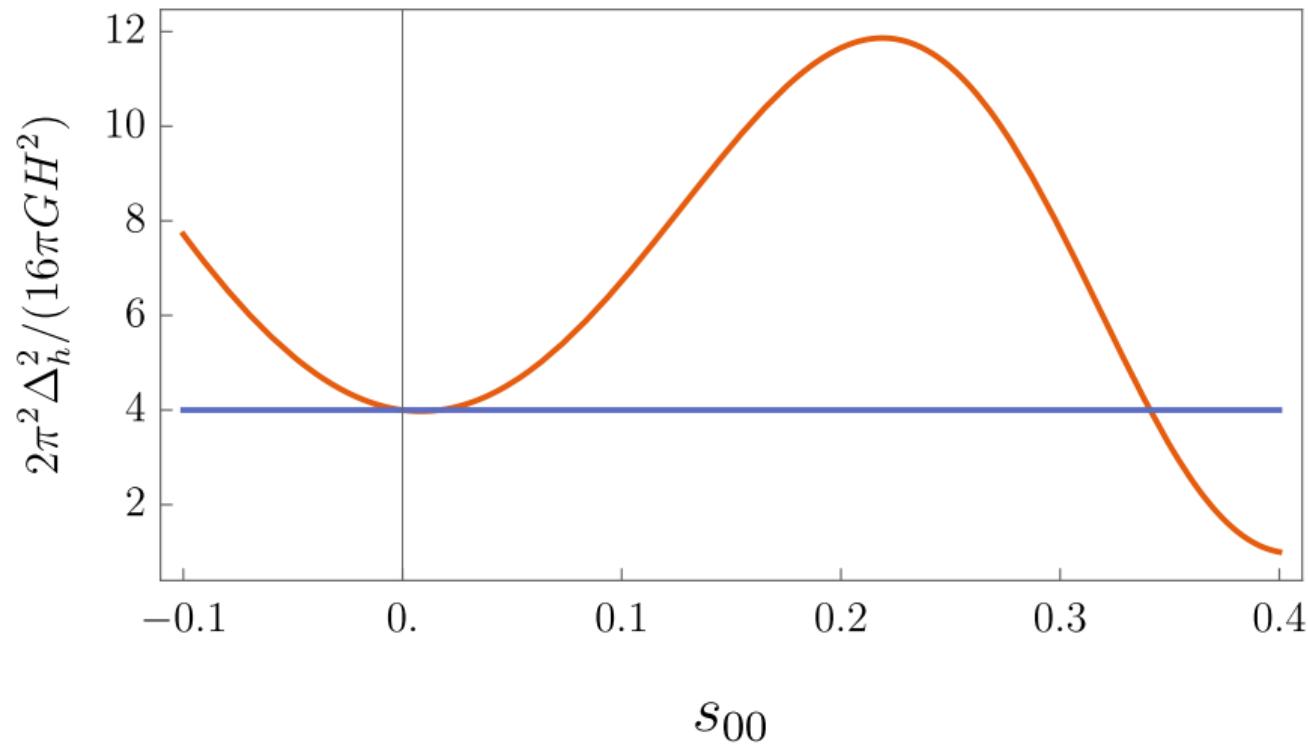
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Wronskian condition holds in the GR limit.

# Solutions

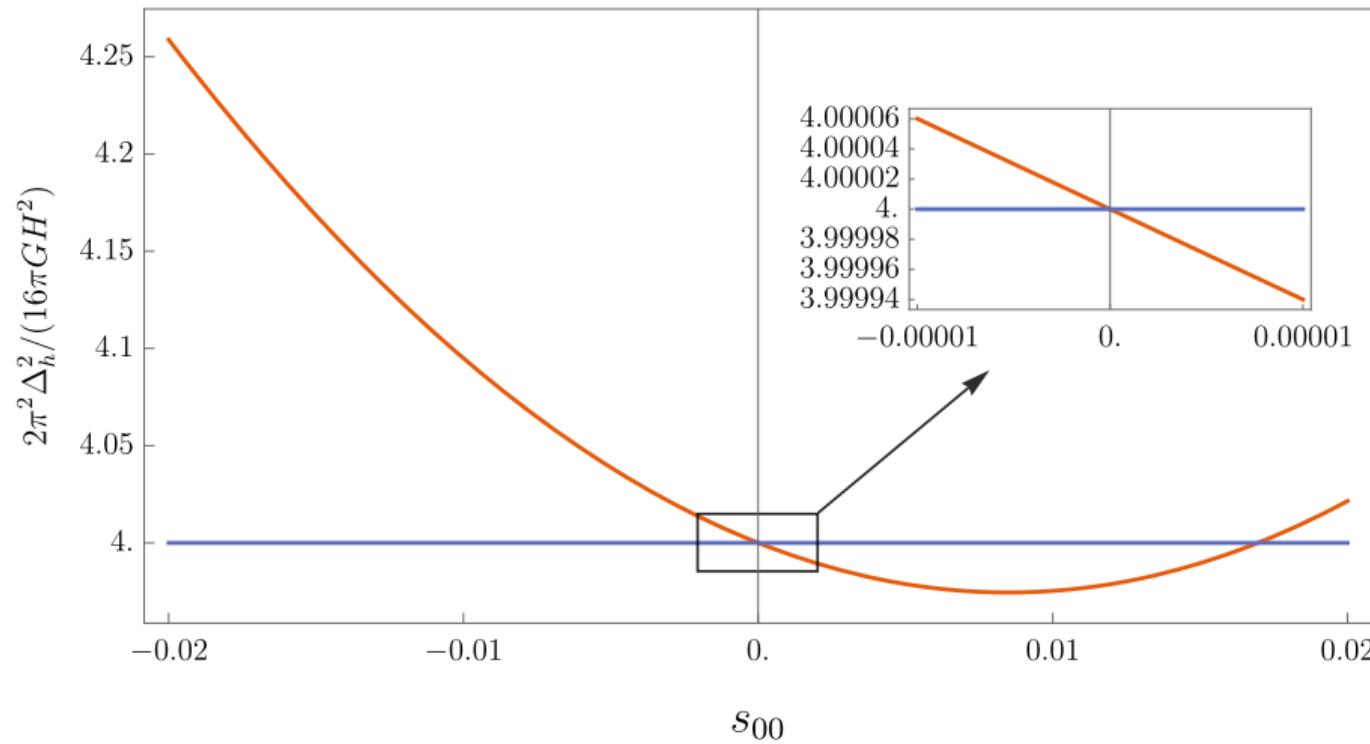


## Solutions



$s_{00}$

# Solutions



# Solutions

