

Alternative Methods
for measuring
Primordial Non-Gaussianity

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Outline

- What is primordial non-Gaussianity?
- Traditional methods
- Squeezed limit methods
- Persistent homology

What is primordial non-Gaussianity

Most information about cosmology we have comes from the two-point function.

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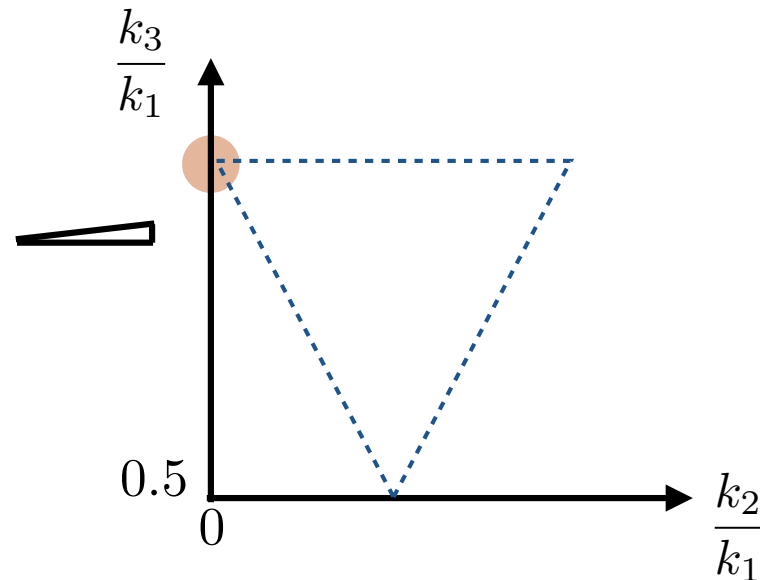
$$\langle \zeta(\vec{k}_1)\zeta(\vec{k}_2)\zeta(\vec{k}_3) \rangle = (2\pi)^3 \delta(\vec{k}_1 + \vec{k}_2 + \vec{k}_3) B(k_1, k_2, k_3)$$

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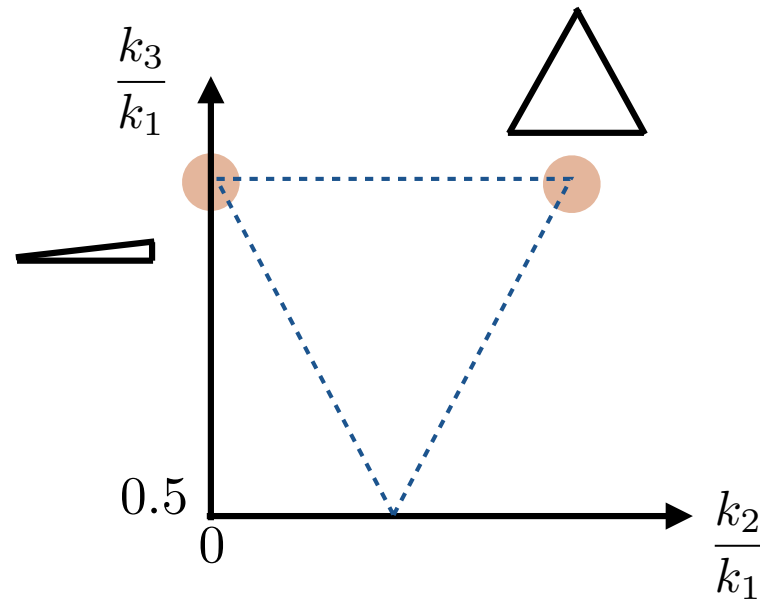


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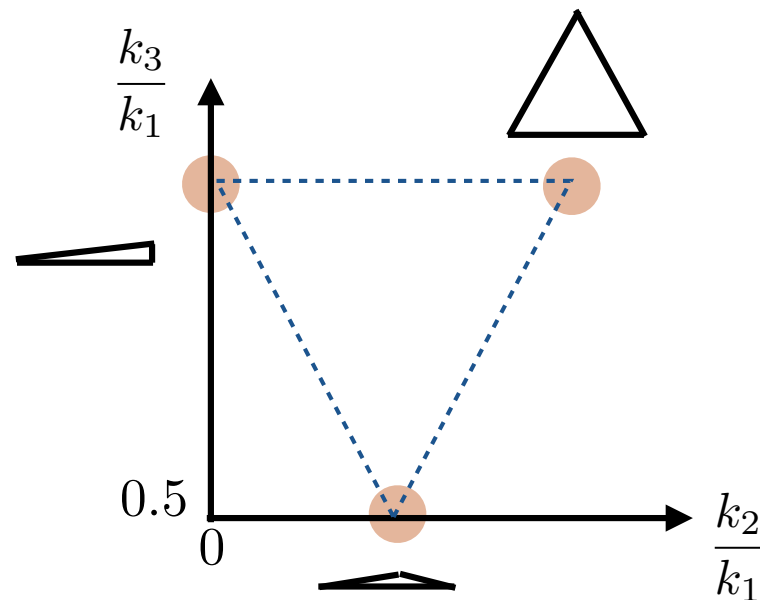


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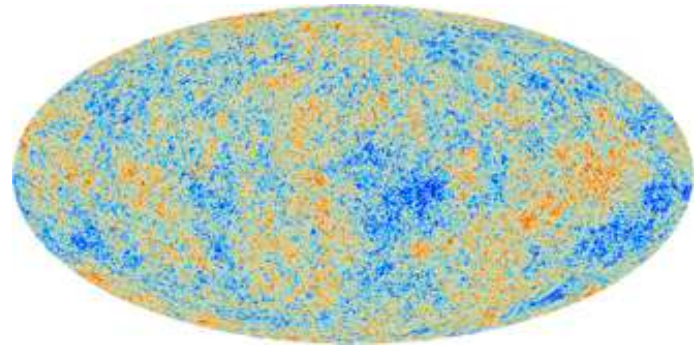
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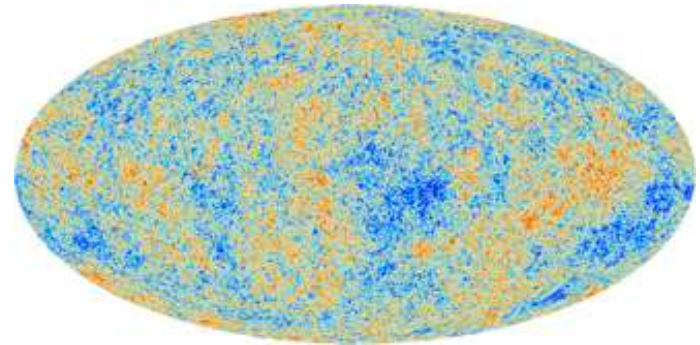
How can it be observed?

$$\left\langle \frac{\delta T}{\bar{T}} \frac{\delta T}{\bar{T}} \frac{\delta T}{\bar{T}} \right\rangle$$



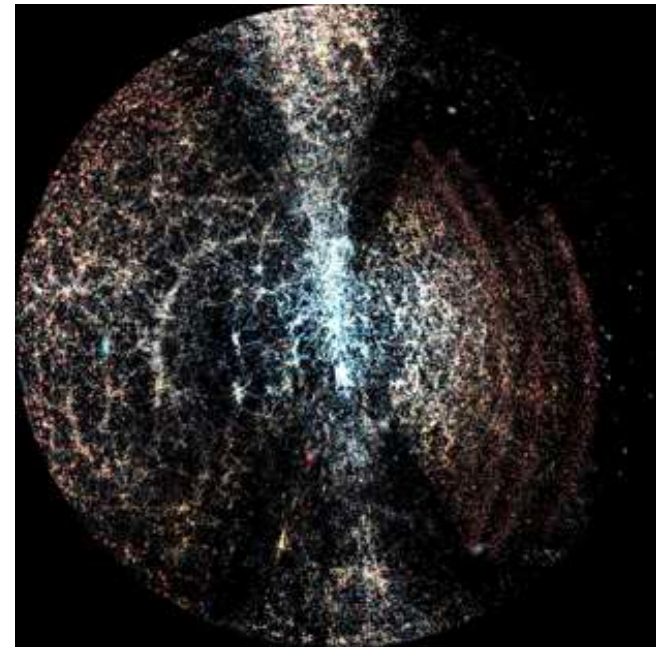
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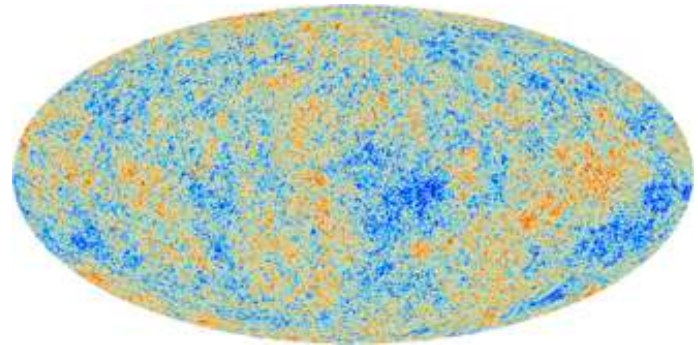
$$\Delta_g = \frac{n_g(z, \hat{n}) - \bar{n}_g(z)}{\bar{n}_g(z)}$$

$$\Delta_g \sim \delta_m$$



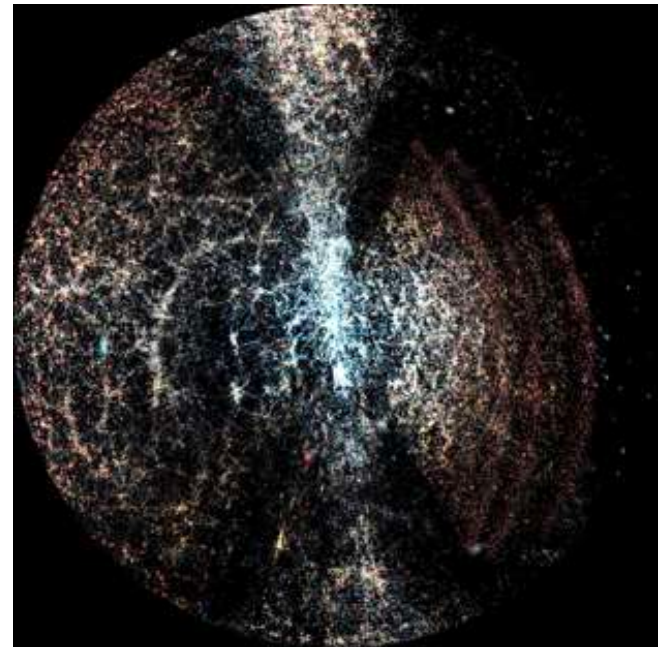
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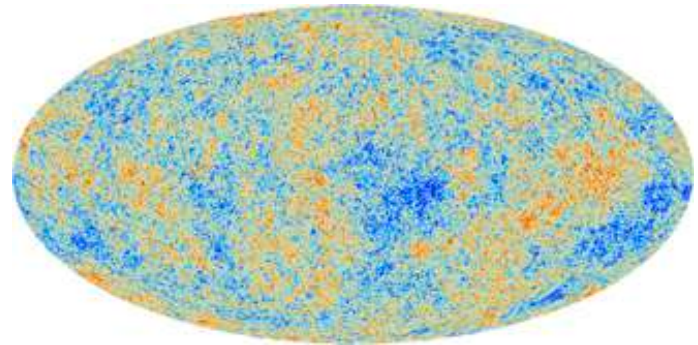
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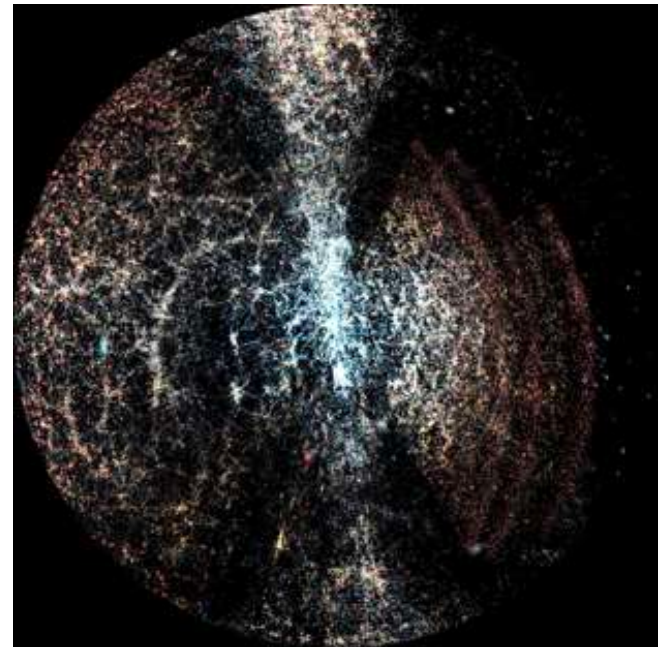


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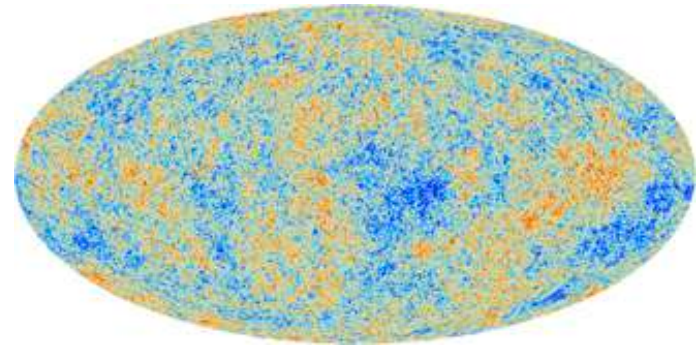


$$\langle \Delta_g \Delta_g \Delta_g \rangle$$



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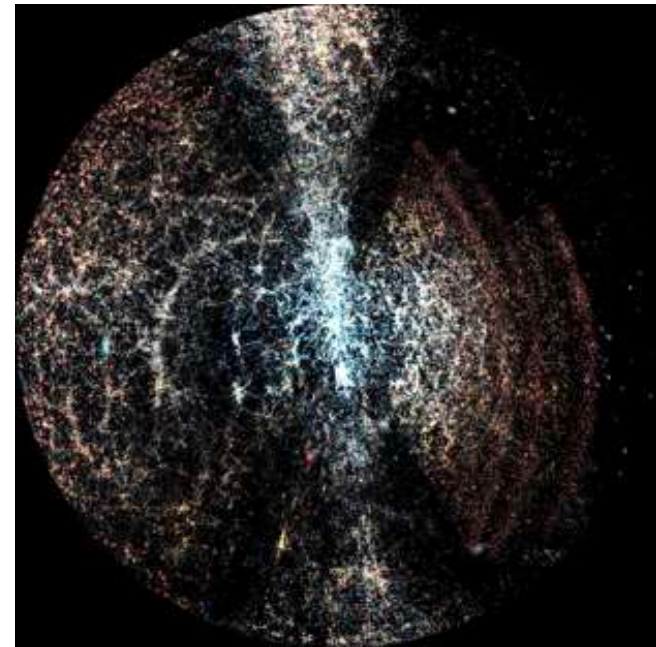


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$$\langle \Delta_g \Delta_g \Delta_g \rangle \text{ and } \langle \Delta_g \Delta_g \rangle$$

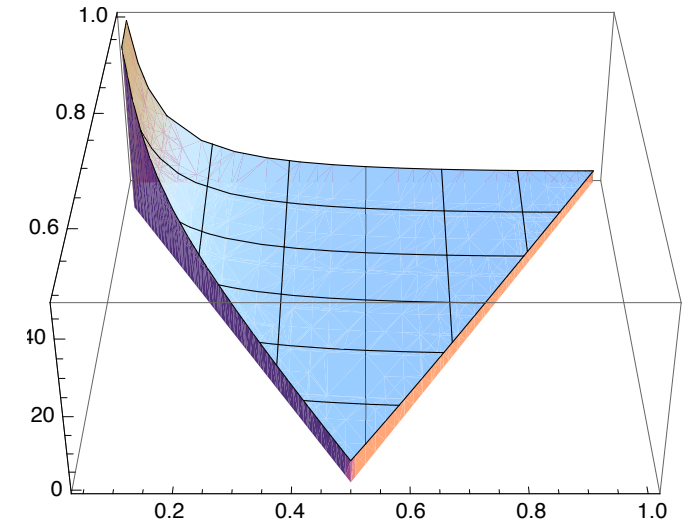


CMB Constraints

Consider a phenomenological model:

$$\zeta = \zeta_g + \frac{3}{5} f_{NL}^{local} \zeta_g^2$$

$$F^{local}(k_1, k_2, k_3) = -2 \frac{3}{5} f_{NL}^{local} A^2 \frac{1}{k_1^3 k_2^3} + 3 \text{ perms.}$$



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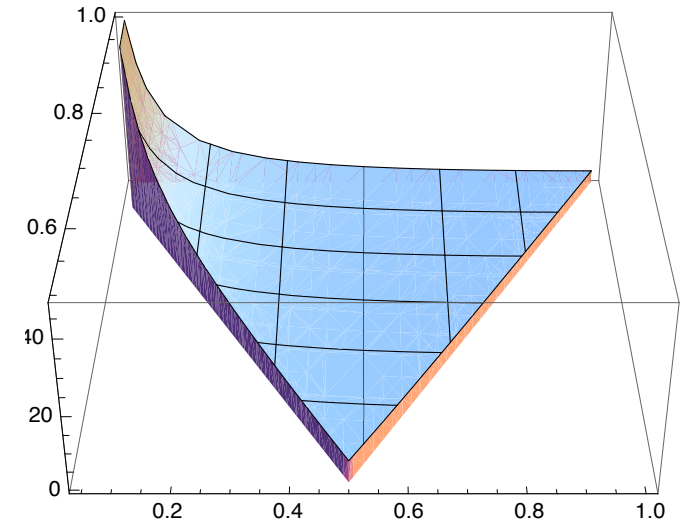
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shape “Overlap”

$$F_1 \cdot F_2 \equiv \sum_{k_1, k_2, k_3} \frac{F_1(k_1, k_2, k_3) F_2(k_1, k_2, k_3)}{\sigma^2(k_1) \sigma^2(k_2) \sigma^2(k_3)}$$

Data is analyzed for simple shapes.



Two shapes are “similar” if they have a cosine of order one.

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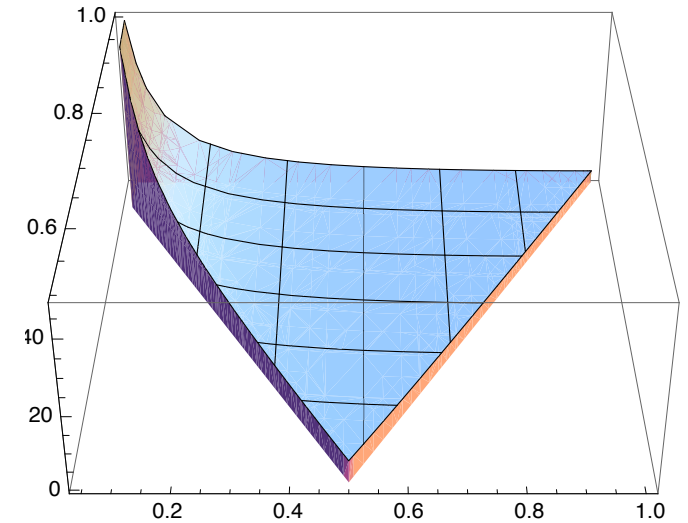
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Data is analyzed for simple shapes.

This shape is the one produced by multi-field models.

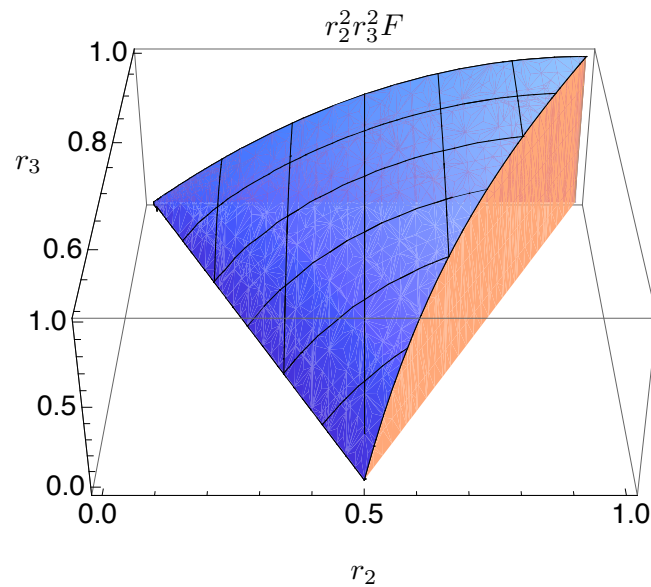


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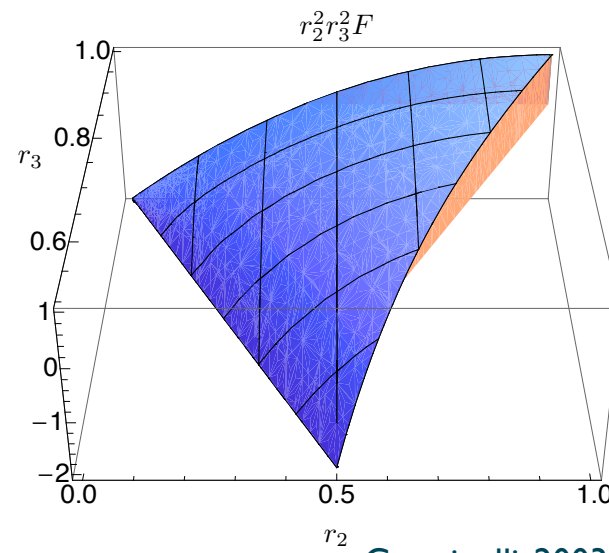
General single field models

Find templates that are like the NG produced by the 2 EFT operators

Equilateral

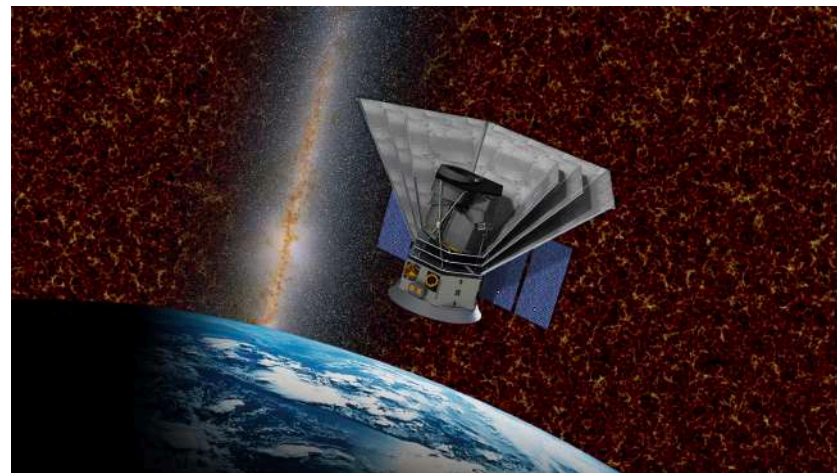
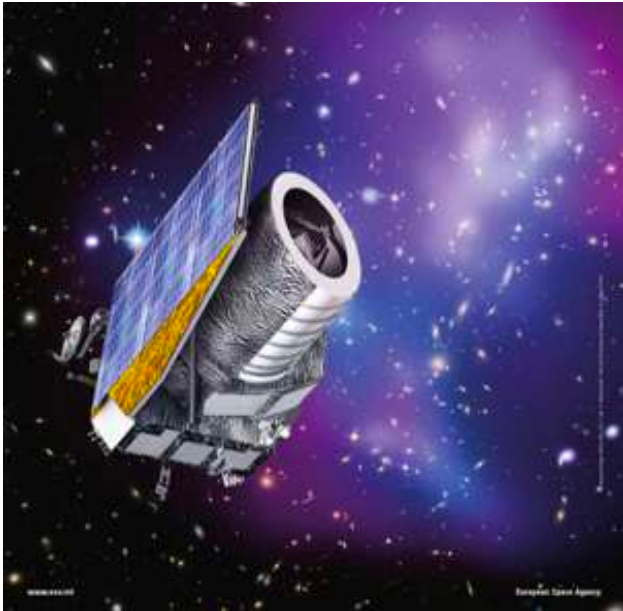


Orthogonal



Creminelli, 2003 [arXiv: astro-ph/0306122]
Cheung et. al., 2008 [arXiv: 0709.0293]
Senatore et. al., 2010 [arXiv: 09053746]

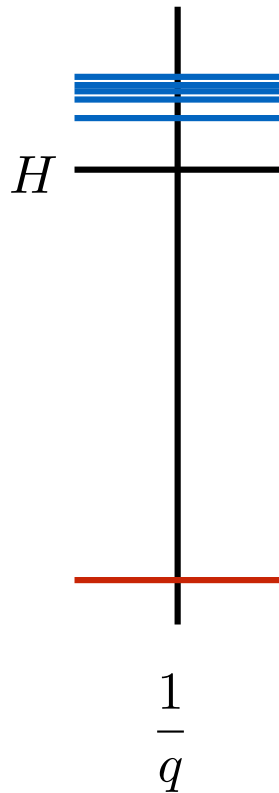
Future: LSS



Squeezed limit information

The squeezed limit contains model independent information about the physics during inflation

Single field



$$B(q, k_1, k_2) \stackrel{q \rightarrow 0}{\sim}$$

J. Maldacena, 2003

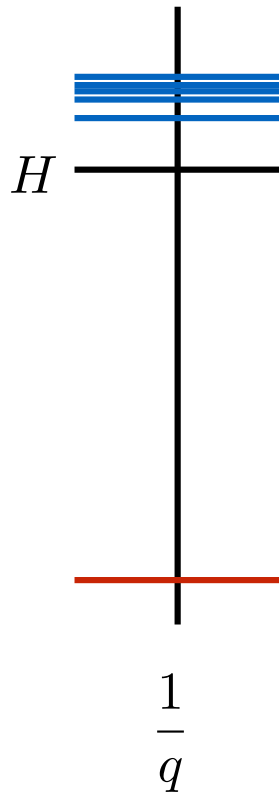
P. Creminelli, M. Zaldarriaga, 2004

P. Creminelli, G. D'Amico, M. Musso, JN, 2011

Squeezed limit information

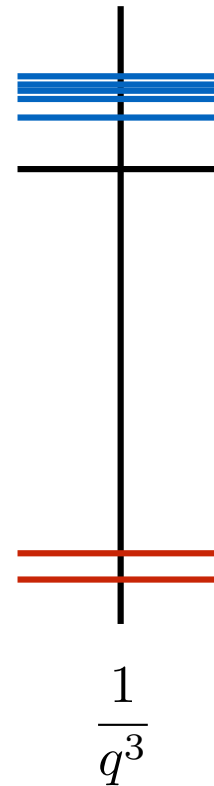
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Multi field

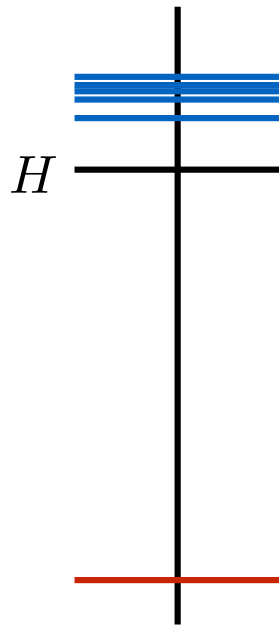


$$\frac{1}{q^3}$$

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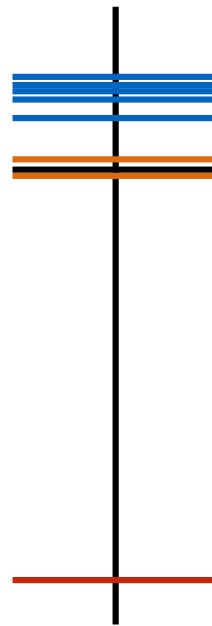
Single field



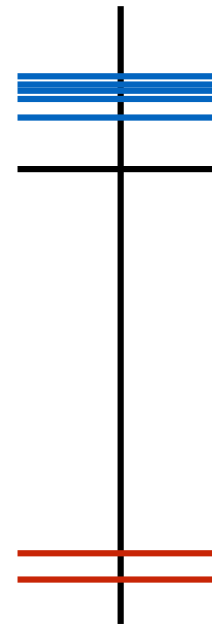
$$\frac{1}{q}$$

$$B(q, k_1, k_2) \stackrel{q \rightarrow 0}{\sim} 0$$

Multi field



$$\frac{1}{q^\alpha} \quad 1 < \alpha < 3$$

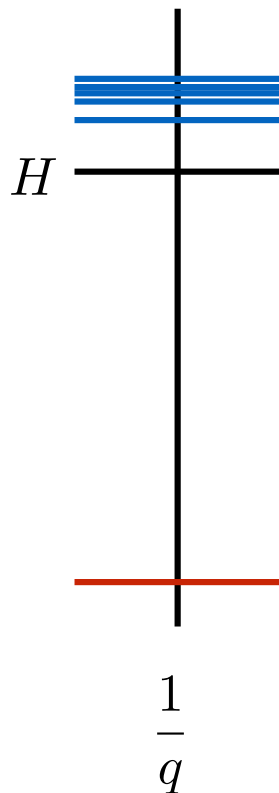


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Single field



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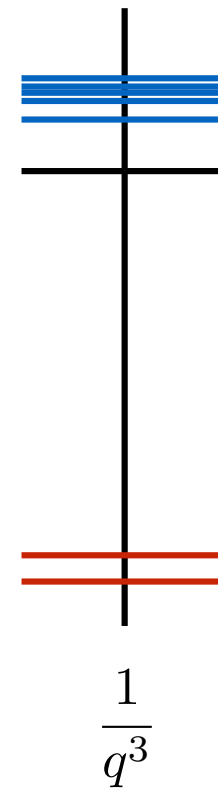
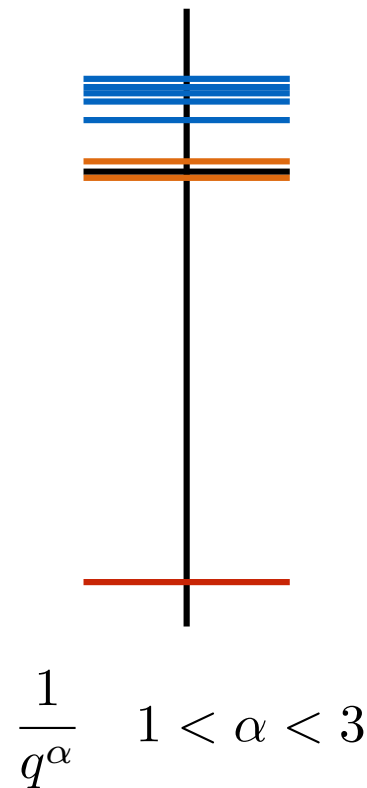
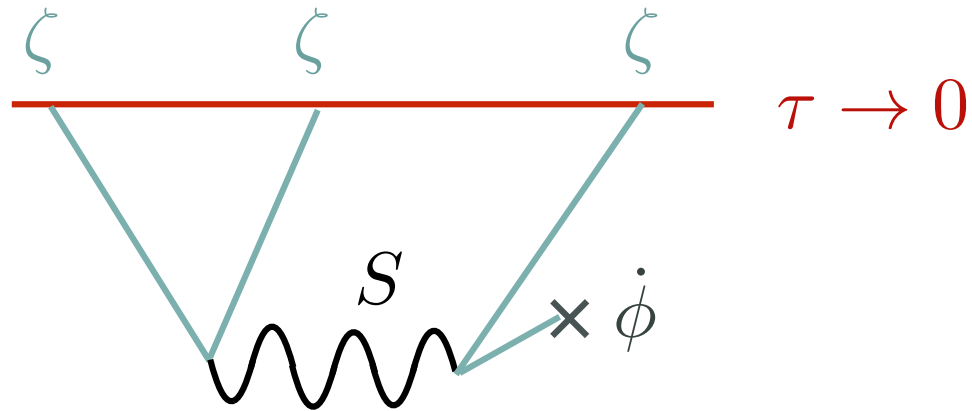


Figure Assassi, Baumann, Green, 2012

Other fields



J. Maldacena, N. Arkani-Hamed, 2015

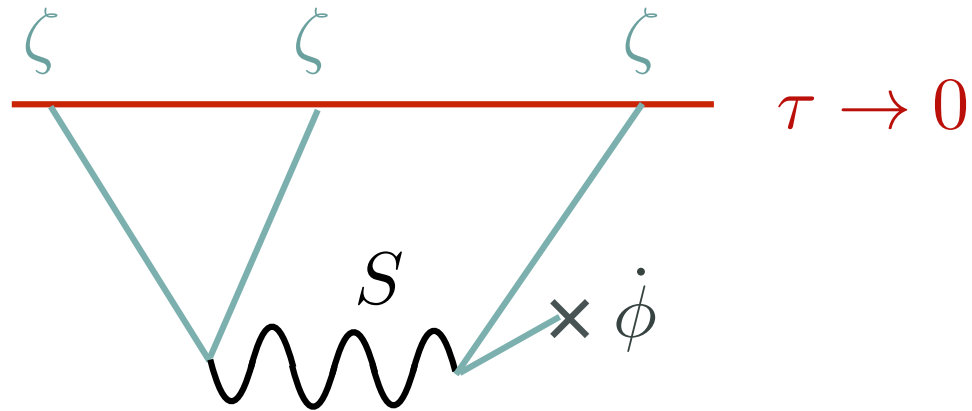
H. Lee, D. Baumann, G. Pimentel, 2016

A. Riotto, A. Kehagias, 2017

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L. Bordin, P. Creminelli, A. Khlemintsky, L. Senatore 2018

Other fields



$$\langle \zeta(q)\zeta(k)\zeta(k) \rangle \sim e^{-\pi\mu} \left[e^{i\delta(\mu)} \left(\frac{q}{k}\right)^{\frac{3}{2}+i\mu} + e^{-i\delta(\mu)} \left(\frac{q}{k}\right)^{\frac{3}{2}-i\mu} \right] P_s(\cos\theta)$$

$$\mu = \sqrt{\frac{m^2}{H^2} - \left(s - \frac{1}{2}\right)}$$

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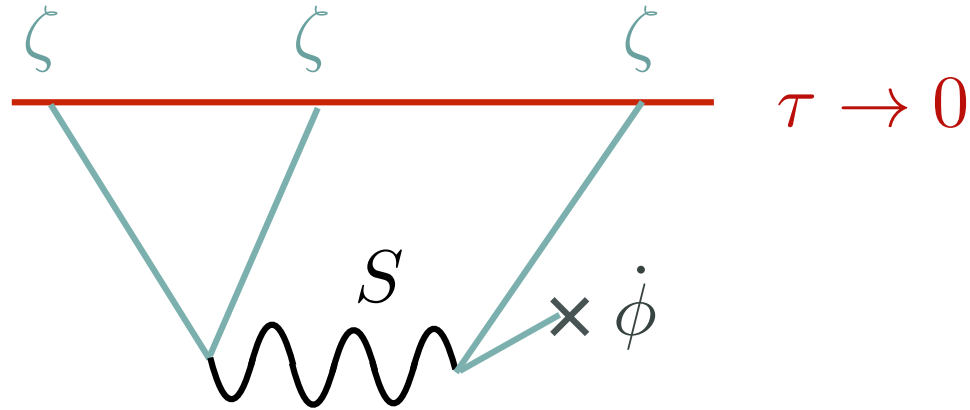
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$\tau \rightarrow 0$

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Characteristic angle dependence

$$\mu = \sqrt{\frac{m^2}{H^2} - \left(s - \frac{1}{2}\right)}$$

J. Maldacena, N. Arkani-Hamed, 2015

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EFT of Inflation

Even within single-field, PNG teaches about the effective Lagrangian

$$\mathcal{L} = -\frac{M_{\text{pl}}^2 \dot{H}}{c_s^2} \left(\dot{\pi}^2 - \frac{c_s^2}{a^2} (\partial_i \pi)^2 \right) + \frac{M_{\text{pl}}^2 \dot{H}}{c_s^2} (1 - c_s^2) \frac{1}{a^2} \dot{\pi} (\partial_i \pi)^2 + \frac{M_{\text{pl}}^2 \dot{H}}{c_s^2} A \dot{\pi}^3$$

Cheung, Creminelli, Fitzpatrick, Kaplan, Senatore, 2007

Constrained through the “equilateral” and “orthogonal” templates.

Senatore, Smith, Zaldarriaga, 2009

Copious particle production during inflation produces large equilateral non-Gaussianity.

The Scale-dependent bias

Bias is the connection of galaxies and matter

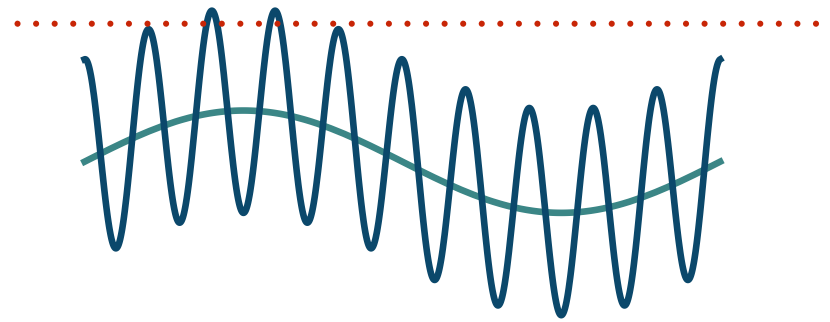
$$\delta_g = b\delta$$

For the local model: $\Phi = \Phi_g + f_{\text{NL}}\Phi_g^2$

Dalal, et. al., 2008

Matarrese, Verde, et. al., 2008

Slosar, et. al., 2008



$$f_{\text{NL}} = 0$$

The Scale-dependent bias

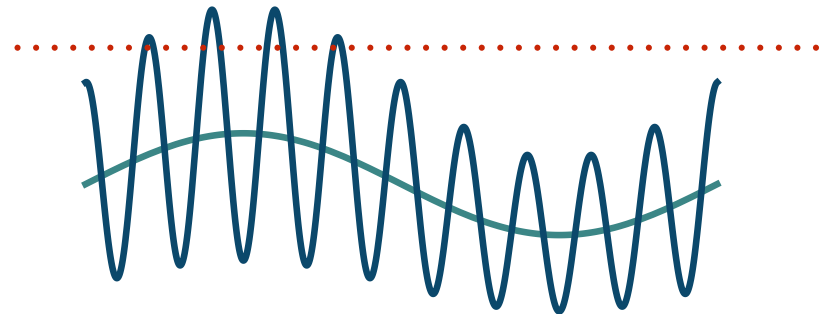
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$$f_{\text{NL}} > 0$$

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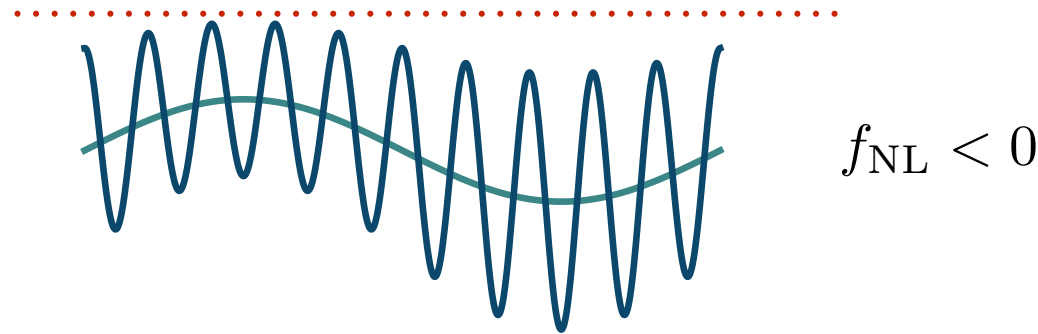
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There is a correlation between Φ and the number of galaxies Δ_g .

$$\langle \Delta_g \Delta_g \rangle \subset \langle \Phi \delta \rangle \sim \frac{1}{k^2} \langle \delta \delta \rangle$$

Sensitive to the squeezed limit!

No coupling with potential

A homogeneous gravitational potential has no physical meaning



No coupling with potential

A homogeneous gravitational potential has no physical meaning

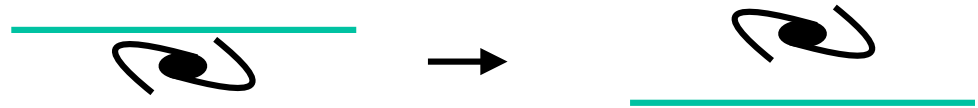
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No coupling with potential

A homogeneous gravitational potential has no physical meaning

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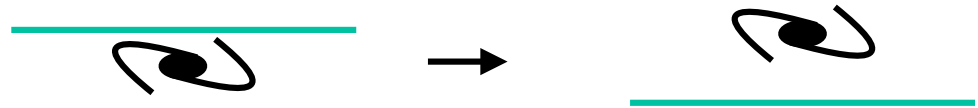
A homogeneous gravitational force can be set to zero by going to a freely falling frame



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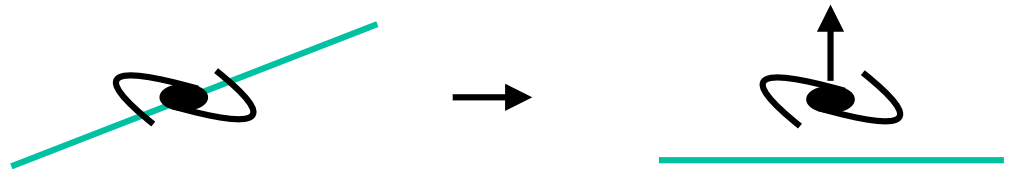
$$\Phi \rightarrow 0$$



A homogeneous gravitational force can be set to zero by going to a freely falling frame

$$\nabla\Phi \rightarrow 0$$

$$\vec{V} \rightarrow \vec{V} - t\nabla\Phi$$



Current constraints

CMB

(68% CL)

Single-field
EFT

Equilateral

$$f_{\text{NL}}^{\text{equi}} = -26 \pm 47$$

Orthogonal

$$f_{\text{NL}}^{\text{orth}} = -38 \pm 24$$

Multi-field

Local

$$f_{\text{NL}}^{\text{loc}} = -0.9 \pm 5.1$$

Planck collaboration 2019

Current constraints

LSS

(68% CL)

Single-field
EFT

Equilateral

$$f_{\text{NL}}^{\text{equi}} = 2 \pm 212$$

Orthogonal

$$f_{\text{NL}}^{\text{orth}} = 126 \pm 72$$

Multi-field

Local

$$f_{\text{NL}}^{\text{loc}} = -33 \pm 28$$

D'Amico, Lewandowski, Senatore, Zhang, 2022

See also

Cabass, Ivanov, Philcox, Simonović, Zaldarriaga, 2022

Forecasts

Karagiannis et. al., 2018

$f_{\text{NL}}^{\text{loc}}$

Doré et. al., 2014

	“Euclid-like”	“LSST inspired”	SPHEREx
P	~ 6	~ 1	~ 1
B	~ 6	~ 0.5	~ 0.2
$P + B$	~ 5	~ 0.5	

Single tracer
Multi tracer

“High-z”

P $\sigma(f_{\text{NL}}^{\text{loc}}) \sim 0.8$

Castorina et. al., 2021

MegaMapper

B $\sigma(f_{\text{NL}}^{\text{equi}}) \sim 17$

$\sigma(f_{\text{NL}}^{\text{orth}}) \sim 8$

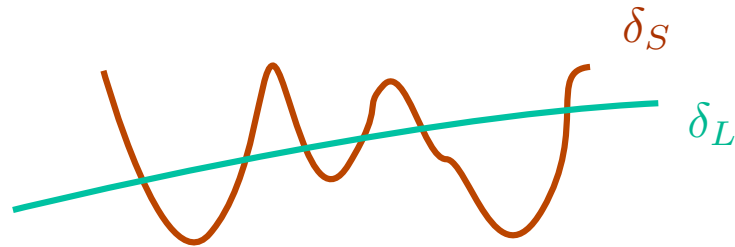
Cabass et. al., 2022

Squeezed limit methods

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$$\langle \delta(\mathbf{q})\delta(\mathbf{k}_1)\delta(\mathbf{k}_2) \rangle = (2\pi)^3 \delta(\mathbf{q} + \mathbf{k}_1 + \mathbf{k}_2) B(q, k_1, k_2)$$

We will be interested in the limit $q \ll k_1, k_2$

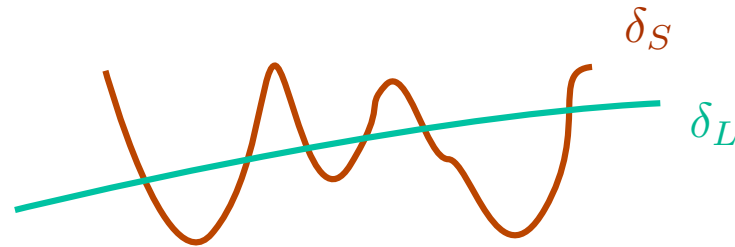


$$\langle \delta(\mathbf{q})\delta(\mathbf{k}_1) \dots \delta(\mathbf{k}_n) \rangle \stackrel{q \rightarrow 0}{\equiv} \langle \delta(\mathbf{q}) \langle \delta(\mathbf{k}_1) \dots \delta(\mathbf{k}_n) \rangle_{\delta_L} \rangle$$

Squeezed limit methods

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$$\langle \delta(\mathbf{q})\delta(\mathbf{k}_1) \dots \delta(\mathbf{k}_n) \rangle \stackrel{q \rightarrow 0}{\approx} \langle \delta(\mathbf{q}) \langle \delta(\mathbf{k}_1) \dots \delta(\mathbf{k}_n) \rangle_{\delta_L} \rangle$$

Correlation between a 2-point function and the potential Φ

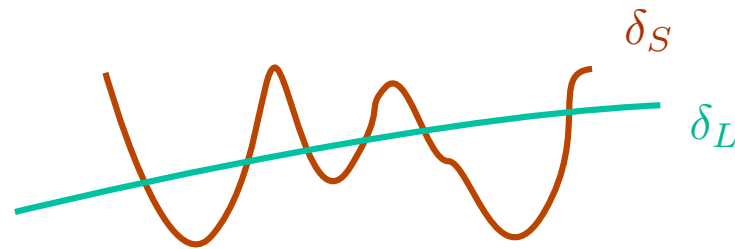
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Local non-Gaussianity.

Squeezed limit methods

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$$\langle \delta(\mathbf{q})\delta(\mathbf{k}_1) \dots \delta(\mathbf{k}_n) \rangle \stackrel{q \rightarrow 0}{\approx} \langle \delta(\mathbf{q}) \langle \delta(\mathbf{k}_1) \dots \delta(\mathbf{k}_n) \rangle_{\delta_L} \rangle$$

Correlation between a 2-point function and the potential Φ

=

Local non-Gaussianity.

Can be pushed to very small scales for the short modes.

Squeezed limit methods

Modulation of the short modes by a long-wavelength mode:

$$\begin{aligned}\delta(\vec{k})|_{\Phi_L} &= \delta(\vec{k})|_{\Phi_L=0} + \frac{(\vec{k} - \vec{q}) \cdot \vec{q}}{\vec{q}^2} \delta_L(\vec{q}) \delta(\vec{k} - \vec{q}) \\ &\quad + \delta_L(\vec{q}) \Delta_1(\vec{k}) + \hat{q}^i \hat{q}^j \delta(\vec{q}) \hat{k}^i \hat{k}^j \Delta_\theta(\vec{k}) + \dots\end{aligned}$$

Squeezed limit methods

Modulation of the short modes by a long-wavelength mode:

Equivalence principle.

↑

$$\delta(\vec{k})|_{\Phi_L} = \delta(\vec{k})|_{\Phi_L=0} + \frac{(\vec{k} - \vec{q}) \cdot \vec{q}}{q^2} \delta_L(\vec{q}) \delta(\vec{k} - \vec{q})$$
$$+ \delta_L(\vec{q}) \Delta_1(\vec{k}) + \hat{q}^i \hat{q}^j \delta(\vec{q}) \hat{k}^i \hat{k}^j \Delta_\theta(\vec{k}) + \dots$$

Squeezed limit methods

Modulation of the short modes by a long-wavelength mode:

$$\delta(\vec{k})|_{\Phi_L} = \delta(\vec{k})|_{\Phi_L=0} + \frac{(\vec{k} - \vec{q}) \cdot \vec{q}}{q^2} \delta_L(\vec{q}) \delta(\vec{k} - \vec{q})$$

↑
Equivalence principle.

$$+ \delta_L(\vec{q}) \Delta_1(\vec{k}) + \hat{q}^i \hat{q}^j \delta(\vec{q}) \hat{k}^i \hat{k}^j \Delta_\theta(\vec{k}) + \dots$$

↓
Dynamical response.

Squeezed limit methods

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Equivalence principle.

$$+ \delta_L(\vec{q}) \Delta_1(\vec{k}) + \hat{q}^i \hat{q}^j \delta(\vec{q}) \hat{k}^i \hat{k}^j \Delta_\theta(\vec{k}) + \dots$$

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Squeezed limit methods

Modulation of the short modes by a long-wavelength mode:

$$\begin{aligned}
 \delta(\vec{k})|_{\Phi_L} &= \delta(\vec{k})|_{\Phi_L=0} + \underbrace{\frac{(\vec{k} - \vec{q}) \cdot \vec{q}}{q^2} \delta_L(\vec{q}) \delta(\vec{k} - \vec{q})}_{\text{Equivalence principle.}} + \underbrace{\Phi_L(\vec{q}) \Delta_{NG}(\vec{k})}_{\text{Local NG}} \\
 &\quad + \underbrace{\delta_L(\vec{q}) \Delta_1(\vec{k}) + \hat{q}^i \hat{q}^j \delta(\vec{q}) \hat{k}^i \hat{k}^j \Delta_\theta(\vec{k}) + \dots}_{\text{Dynamical response.}}
 \end{aligned}$$

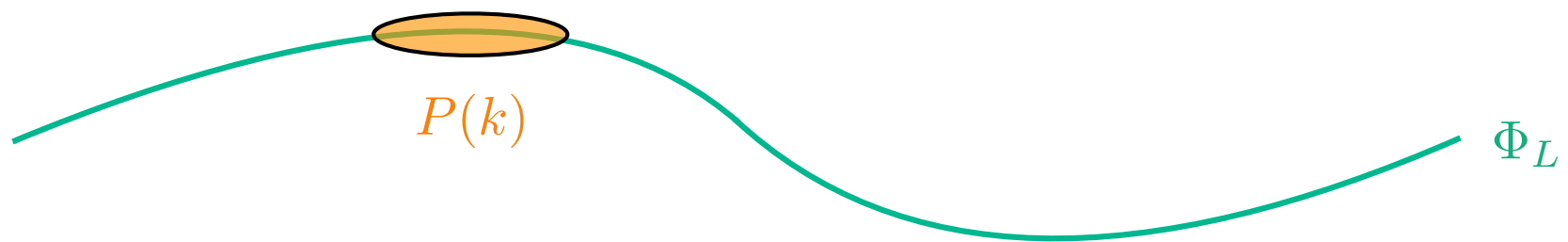
Simple covariance, easy to model, non-linear scales, complementary information.

Modulated power spectrum

Modulation of power spectrum

De Putter, 2018

$$P(k)|_{\Phi_L} = P(k)|_{\Phi_L=0} + \Phi_L(q)P_{\delta\Delta_{NG}}(k) + \delta_L(q)P_{\delta\Delta}(k) + \dots$$

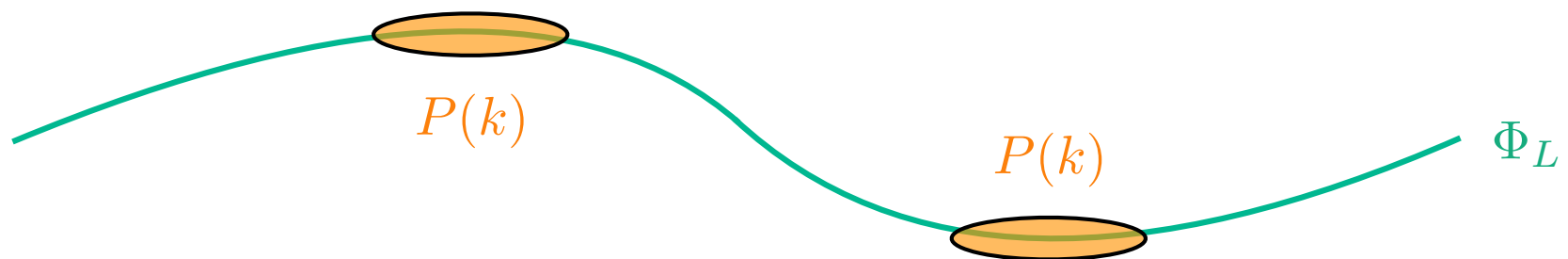


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$1/q^2$ modulation

Modulated power spectrum

Modulation of power spectrum

De Putter, 2018

$$P(k)|_{\Phi_L} = P(k)|_{\Phi_L=0} + \underbrace{\Phi_L(q)P_{\delta\Delta_{NG}}(k)}_{\substack{\downarrow \\ 1/q^2 \text{ modulation}}} + \delta_L(q)P_{\delta\Delta}(k) + \dots$$

Tested in Quijote simulations.

Giri, Muehmer, Smith, 2023

$$0.006 < q < 0.047$$

$$0.5 < k < 3.0$$

Modulated power spectrum

Modulation of power spectrum

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Tested in Quijote simulations.

Giri, MÜchmeyer, Smith, 2023

$$0.006 < q < 0.047$$

$$0.5 < k < 3.0$$

Powerful probe using matter field.

Due to shot noise, comparable with scale-dependent bias using halos (in Quijote).

Squeezed bispectrum

Leading orders in q

$$B(q, k_1, k_2) = P(q)P(k) \left[\frac{1}{q^2} R_{NG}(k_1) + \frac{\vec{k}_1 \cdot \vec{q}}{q^2} + R_1(k_1) + R_\theta(k_1) (\hat{k}_1 \cdot \hat{q})^2 \right] + (1 \leftrightarrow 2)$$

Squeezed bispectrum

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Factor k dependence by averaging over it

Esposito, Hui, Scoccimarro, 2019

Goldstein, Esposito, et. al., 2022

$$B(q) \sim \sum_k B(q, k_1, k_2) \sim \frac{1}{q^2} a_{NG} + a_0 + a_1 q + \dots$$

Squeezed bispectrum

Leading orders in q

$$B(q, k_1, k_2) = P(q)P(k) \left[\frac{1}{q^2} R_{NG}(k_1) + \frac{\vec{k}_1 \cdot \vec{q}}{q^2} + R_1(k_1) + R_\theta(k_1)(\hat{k}_1 \cdot \hat{q})^2 \right] + (1 \leftrightarrow 2)$$

Factor k dependence by averaging over it

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$$B(q) \sim \sum_k B(q, k_1, k_2) \sim \frac{1}{q^2} a_{NG} + a_0 + a_1 q + \dots$$

Don't find substantial improvement over scale-dependent bias.

$$0.006 < q < 0.047$$

$$0.5 < k < 3.0$$

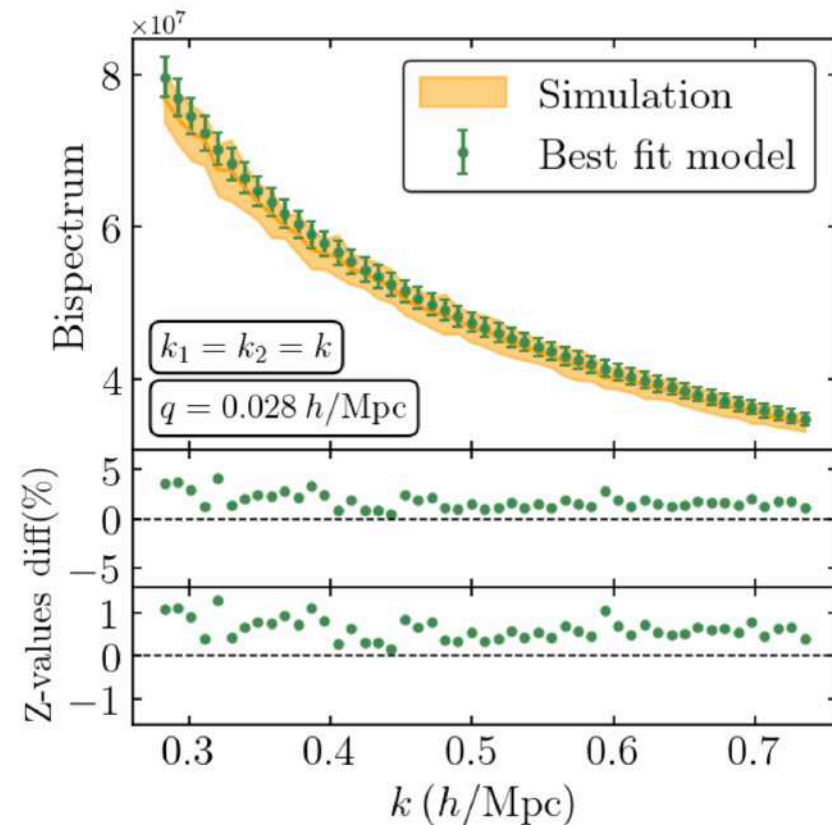
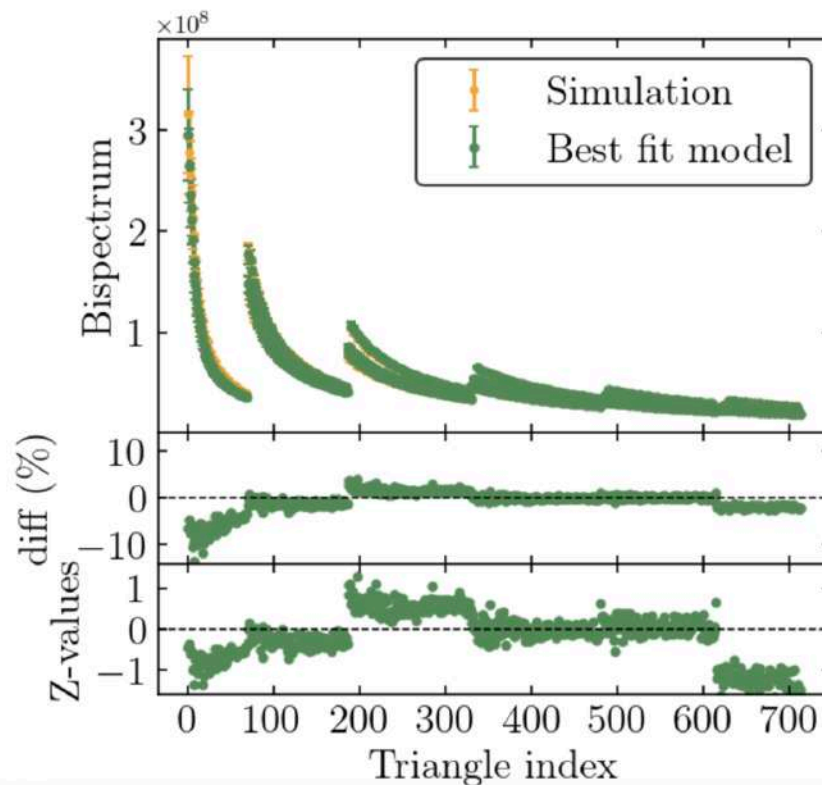
Very easy to model, simple covariance, GR effects.

Squeezed bispectrum

$$B(q, k_1, k_2) = P(q)P(k) \left[\frac{\vec{k}_1 \cdot \vec{q}}{q^2} + R_1(k_1) + R_\theta(k_1)(\hat{k}_1 \cdot \hat{q})^2 \right] + (1 \leftrightarrow 2)$$

$$R_1(k) = \frac{10}{7} + S_1^1 k + S_2^1 k^2 + S_3^1 k^3$$

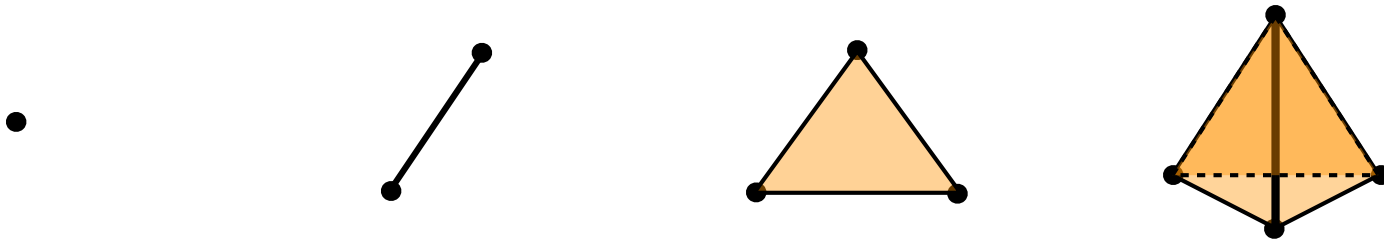
$$R_\theta(k) = \frac{4}{7} + S_1^\theta k + S_2^\theta k^2 + S_3^\theta k^3$$



Persistent Homology

Persistent Homology

Use topological features to constrain NG



In 3D three homology groups

H_0 “Clusters”

H_1 “Loops”

H_2 “Voids”

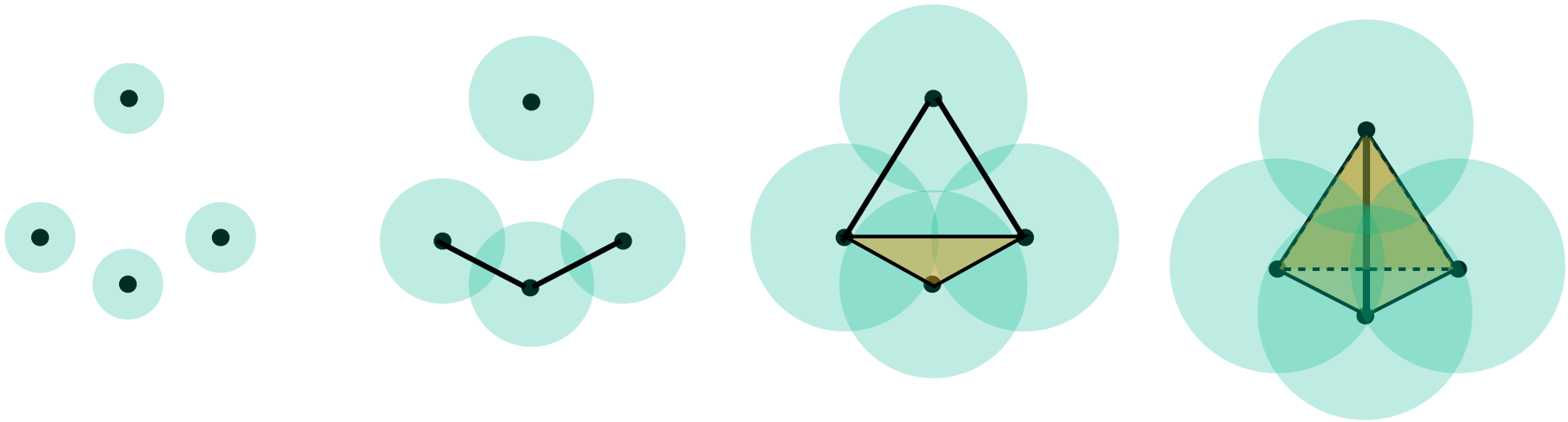
Build simplicial complexes out of your point cloud. Thus have a systematic way of counting structures.

Cole, Shiu, 2017

Biegetti, Cole, Shiu, 2020

Persistent Homology

How to build the complex? Connect points if they are less than a certain distance apart.



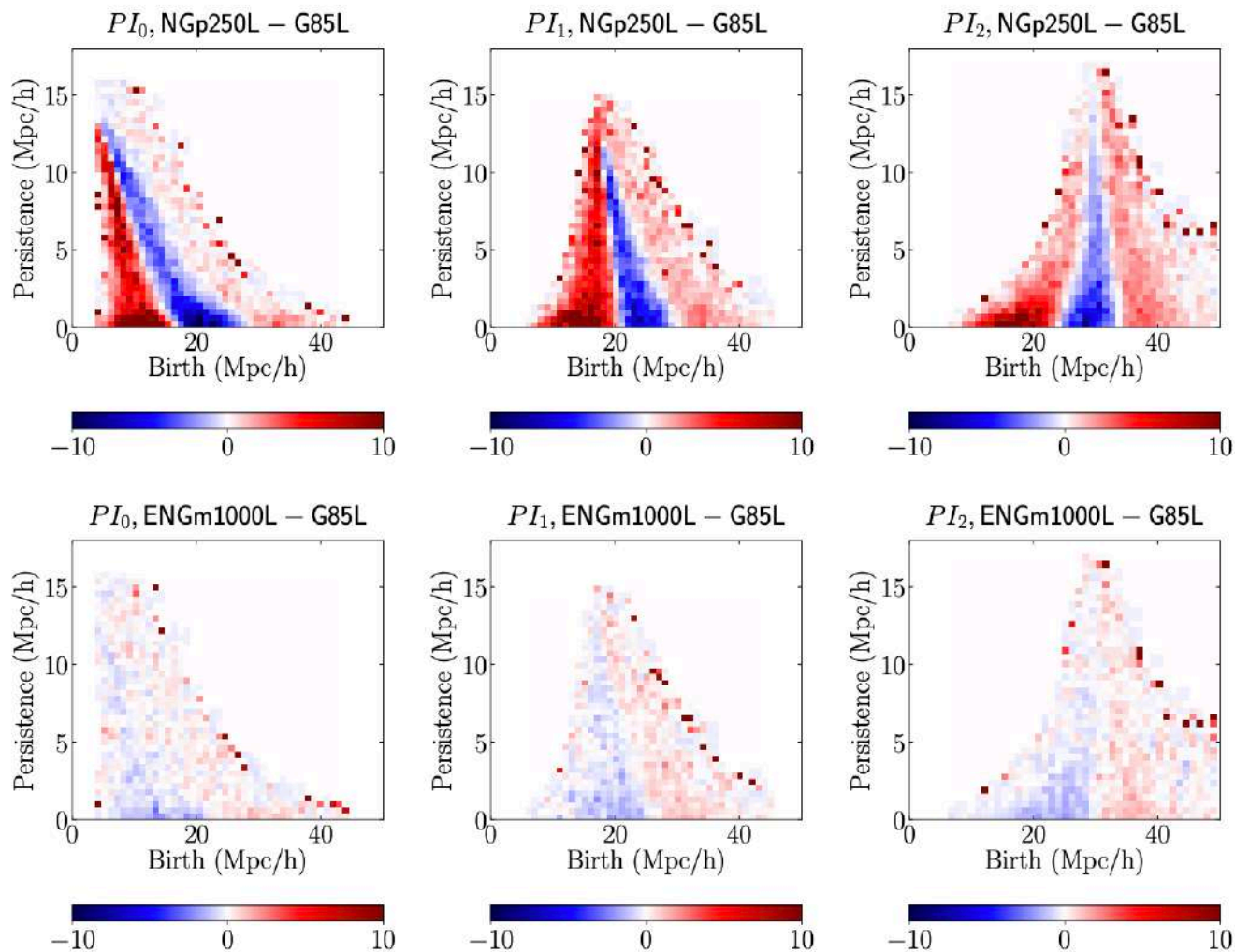
Very sensitive to outliers. A single halo in a void can destroy it.

Choose a better “filtration”, with a “scale” parameter.

Thus, for each feature we have a birth and a death.

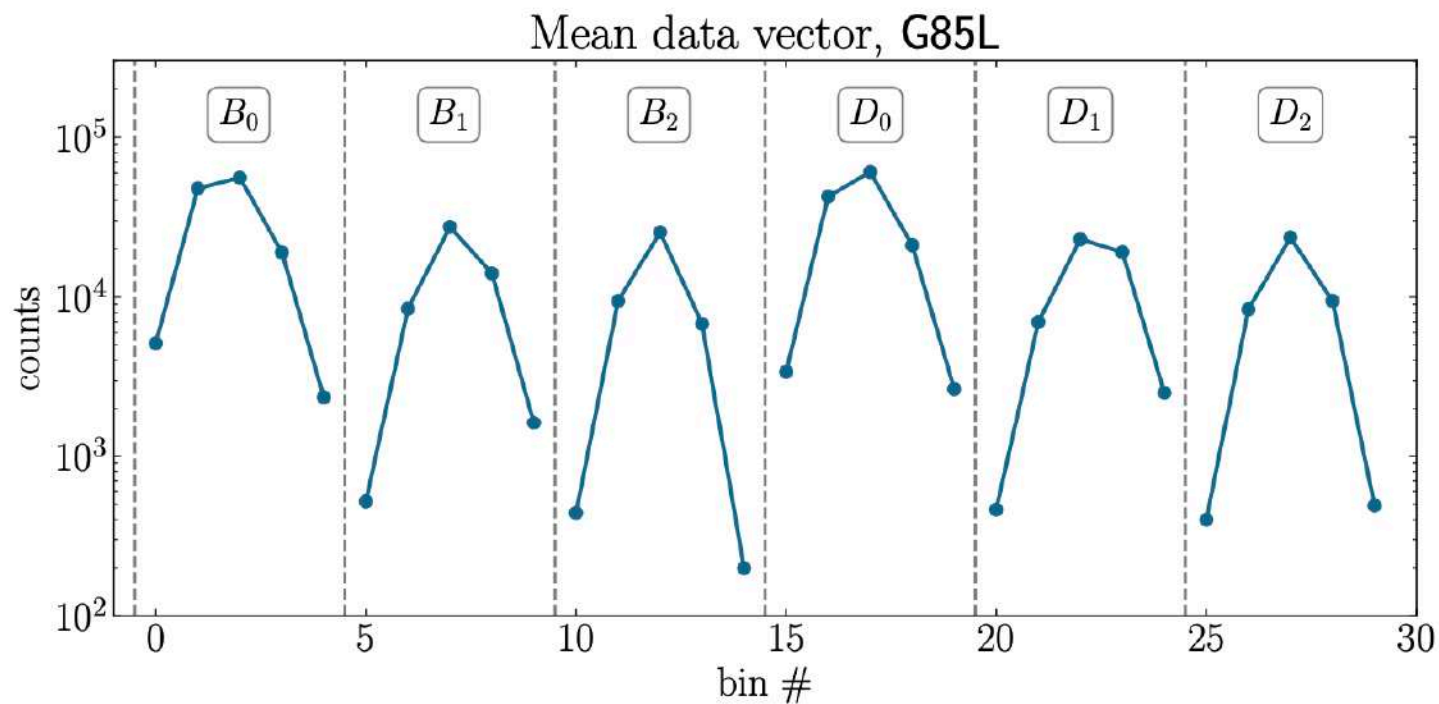
Persistent Homology

Plot weighted histograms of birth vs persistence (death - birth)



Persistent Homology

As data vectors choose a few bins of births and deaths



Persistent Homology

Use simulations to compute derivatives for a Fisher analysis

In redshift space, with different types of errors.

Spectroscopic

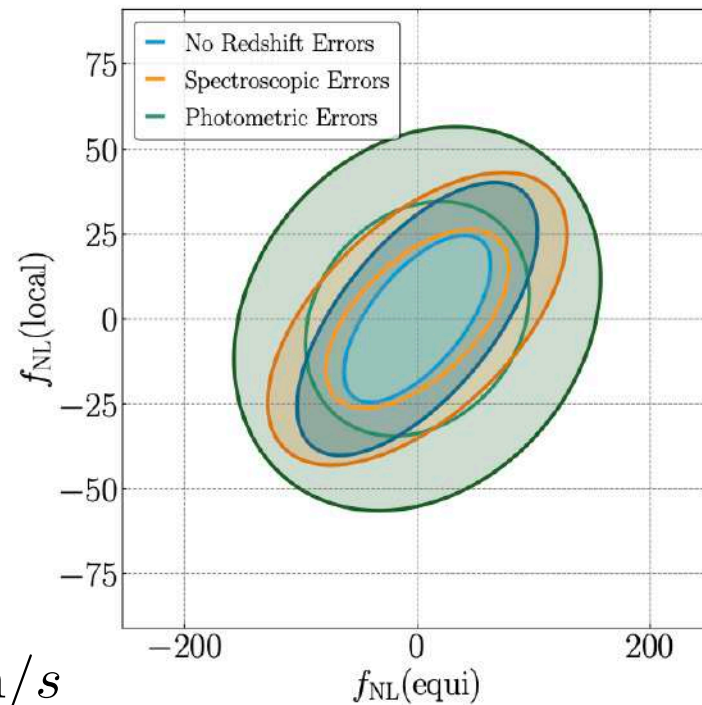
$$\vec{s} = \vec{x} + \frac{\vec{v} \cdot \hat{n}}{\mathcal{H}} \hat{n} + \frac{\vec{v}_{spec}}{\mathcal{H}} \hat{n}$$

$$\sigma_{spec} \sim \sigma_v / 5 \sim 30 \text{ km/s}$$

Photometric

$$\vec{s} = \vec{x} + \frac{\vec{v}_{photo}}{\mathcal{H}} \hat{n}$$

$$\frac{\Delta z}{(1+z)} \sim 0.01 \rightarrow \sigma_{photo} \sim 3000 \text{ km/s}$$



$$V \sim 1 \text{ (Gpc/h)}^3$$

Conclusions

- Hard to get $\Delta f_{\text{NL}}^{\text{equi}} \sim 1$ with standard techniques.
- Scale dependent bias still most promising probe for local PNG. With bispectrum to break degeneracies.
- Squeezed limit techniques could be refined to provide “many scale dependent biases”.
- Persistence homology techniques can be used to greatly improve limits on equilateral PNG. But still need a lot of work.

THE
END