

The Big Bang in the Lab: Simulating the Early Universe with Bose-Einstein Condensates

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[in prep]

Related work :

JHEP 07 (2018) 014

JHEP 10 (2019) 174

Cosmology from Home 2023

Overview

A. Inflationary Cosmology Review

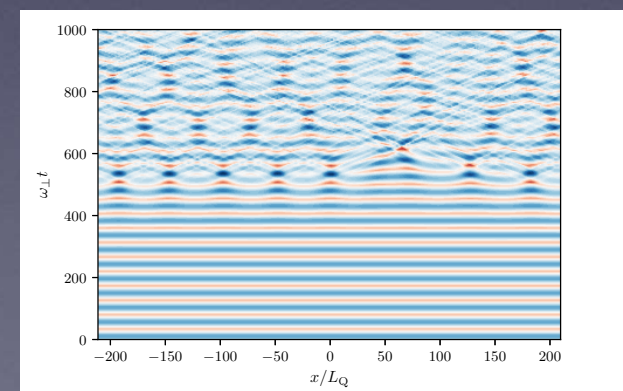
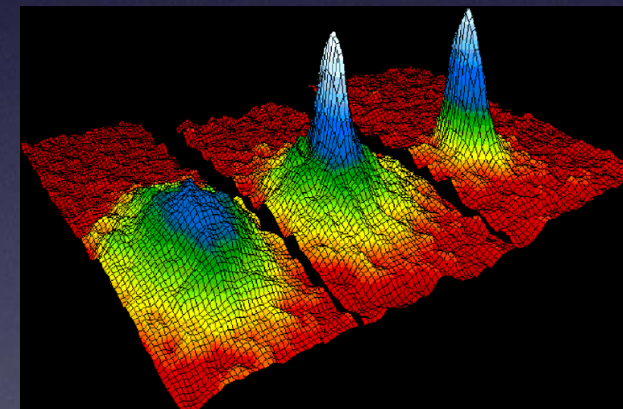
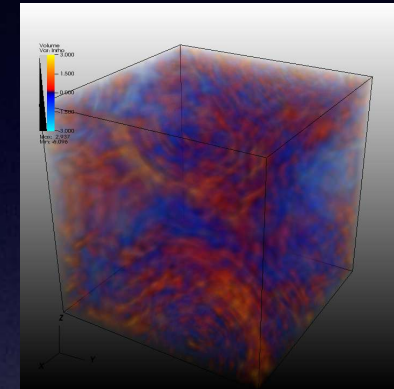
- ▶ Epochs of nonlinearity
- ▶ Preheating
- ▶ Phase transitions (vacuum decay)

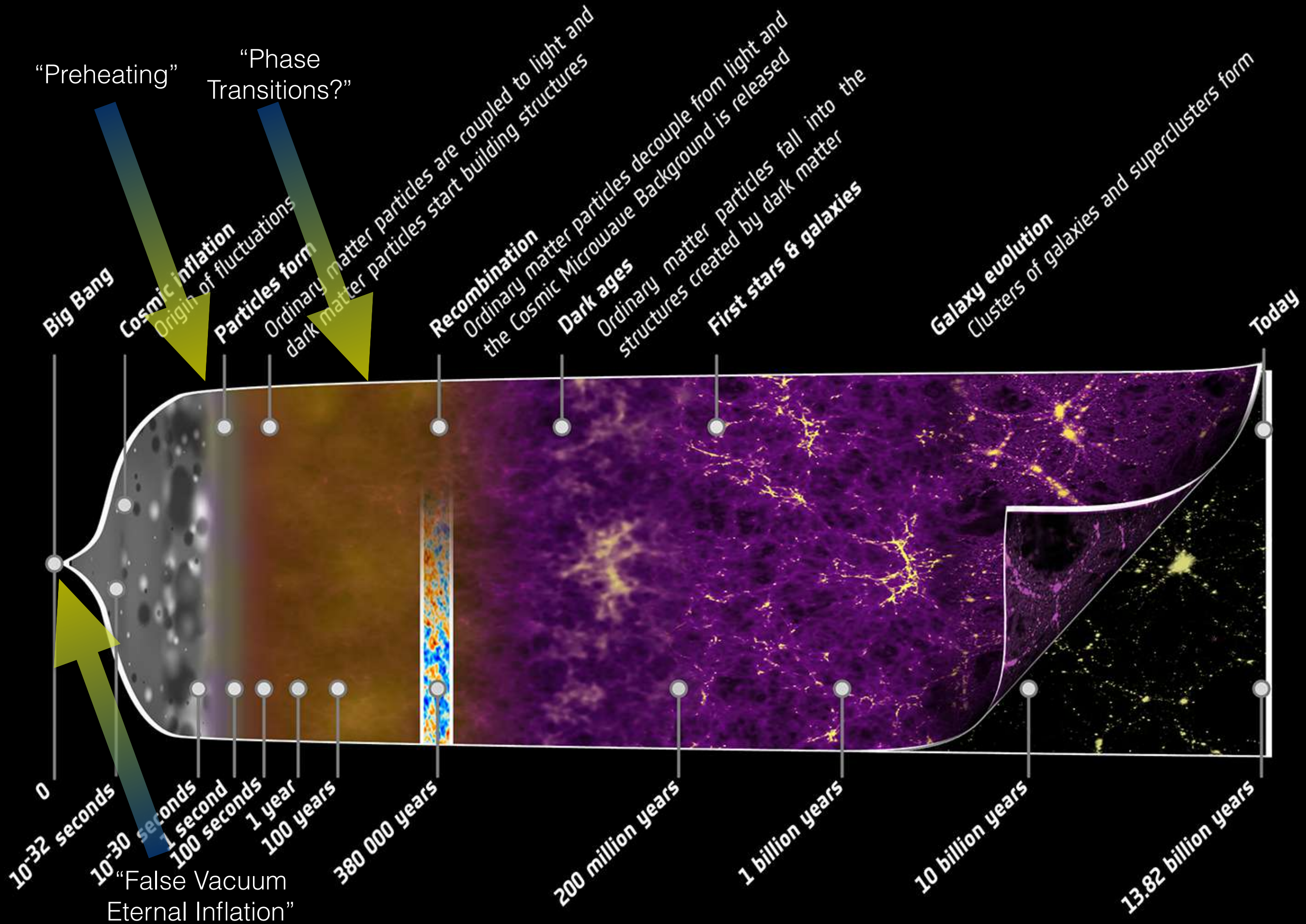
B. Analog BEC Dynamics

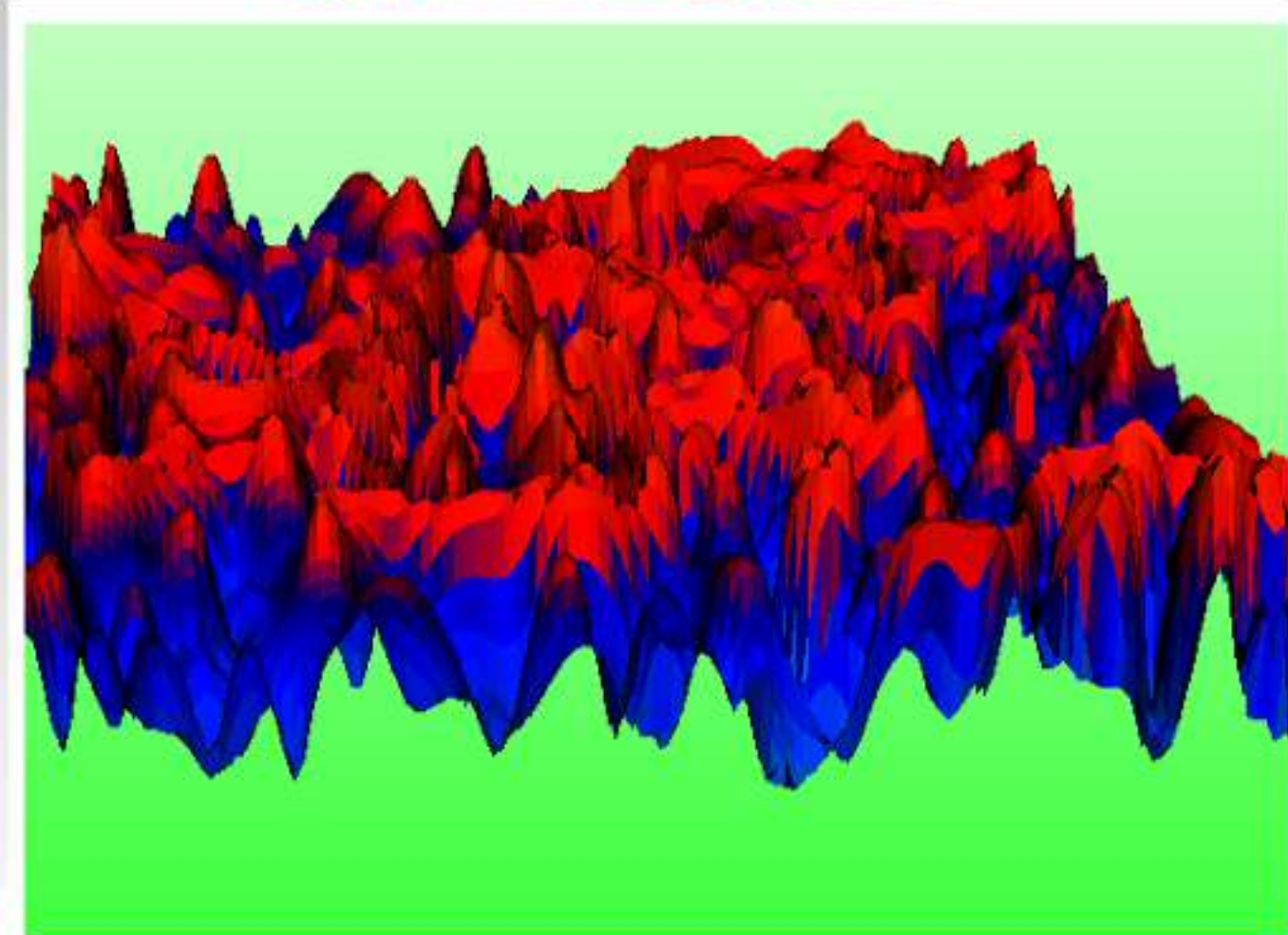
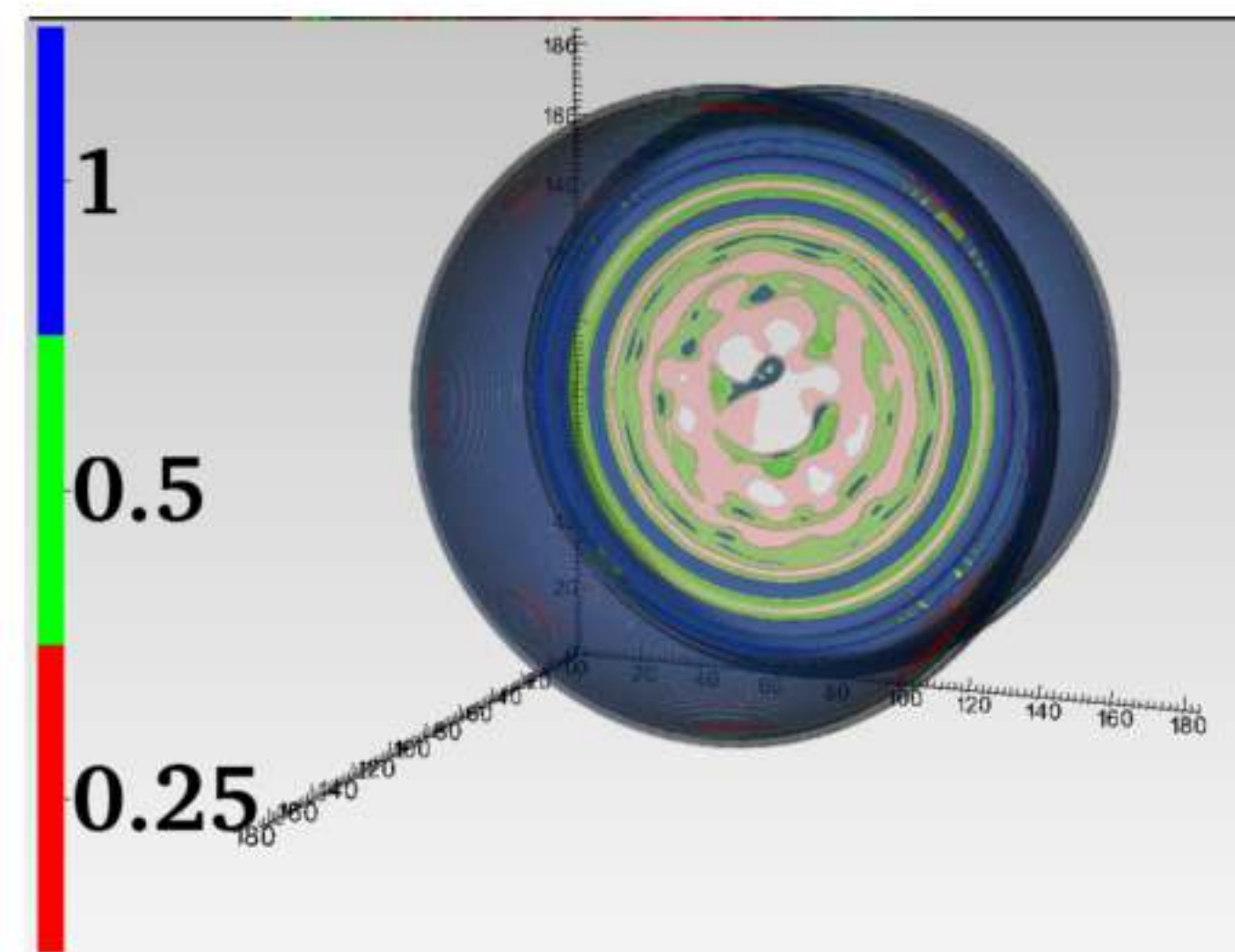
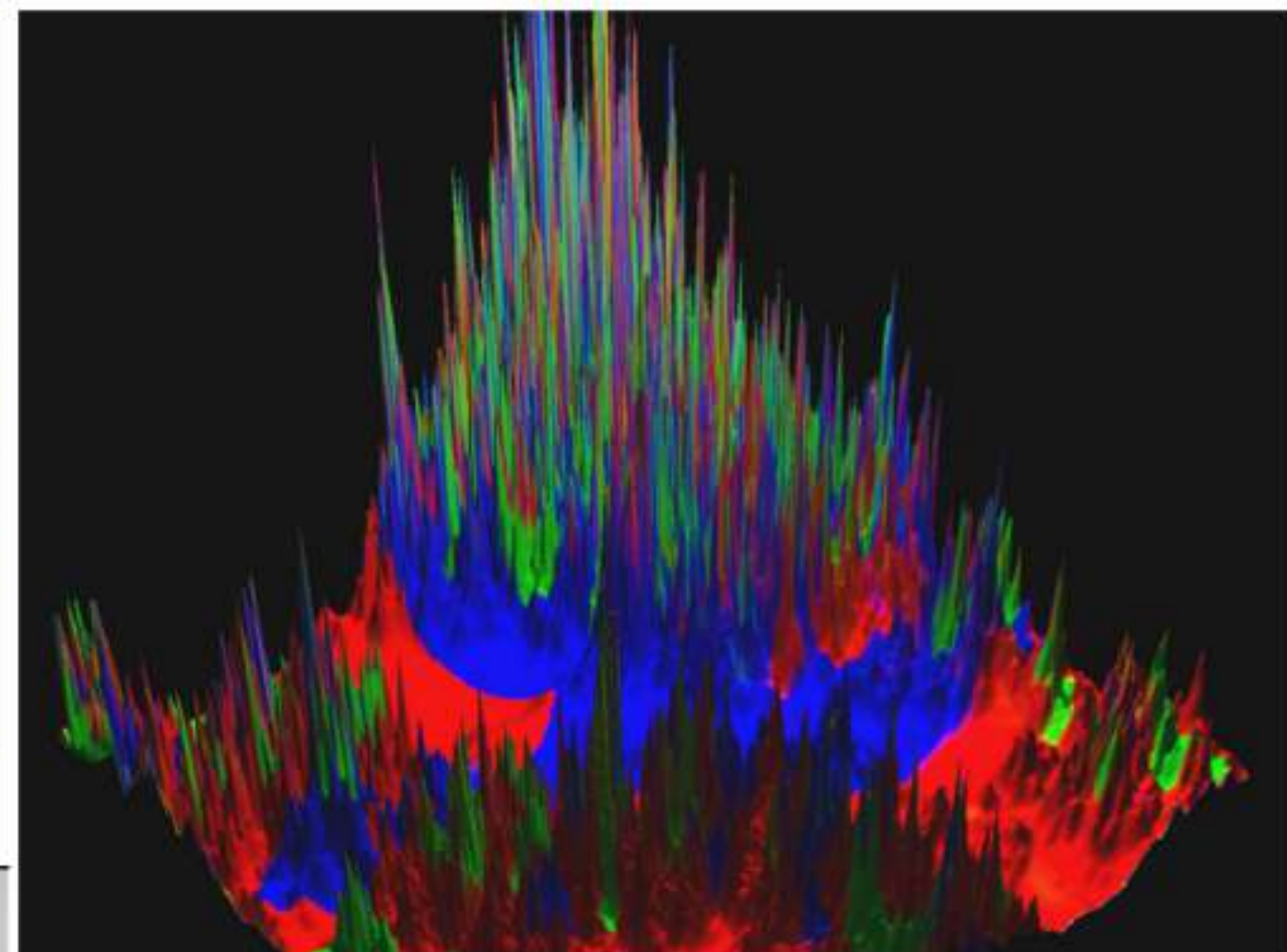
- ▶ Mapping between cosmology and cold atoms

C. Preheating (end of inflation) in cold atom systems

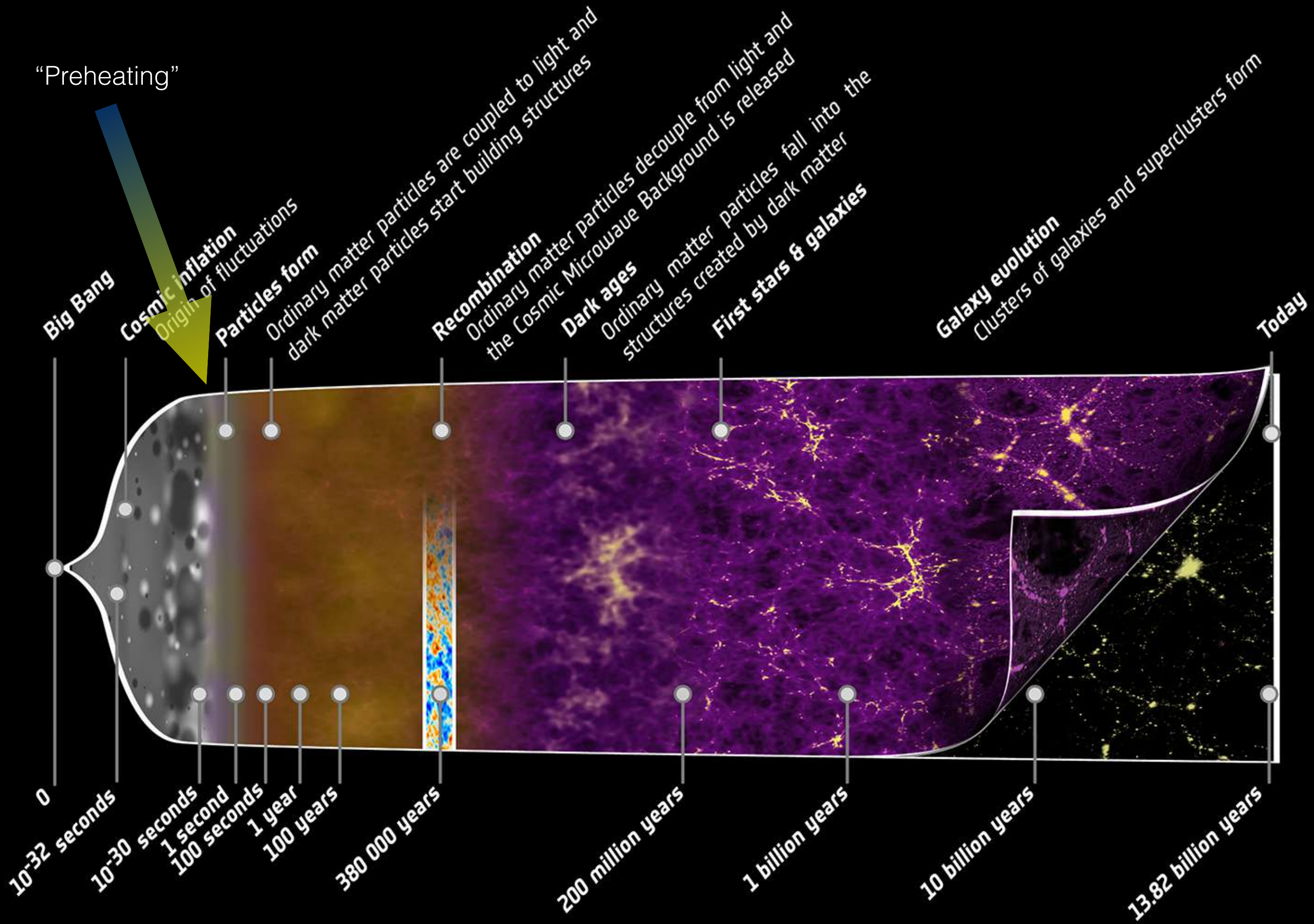
D. A little vacuum decay mixed in



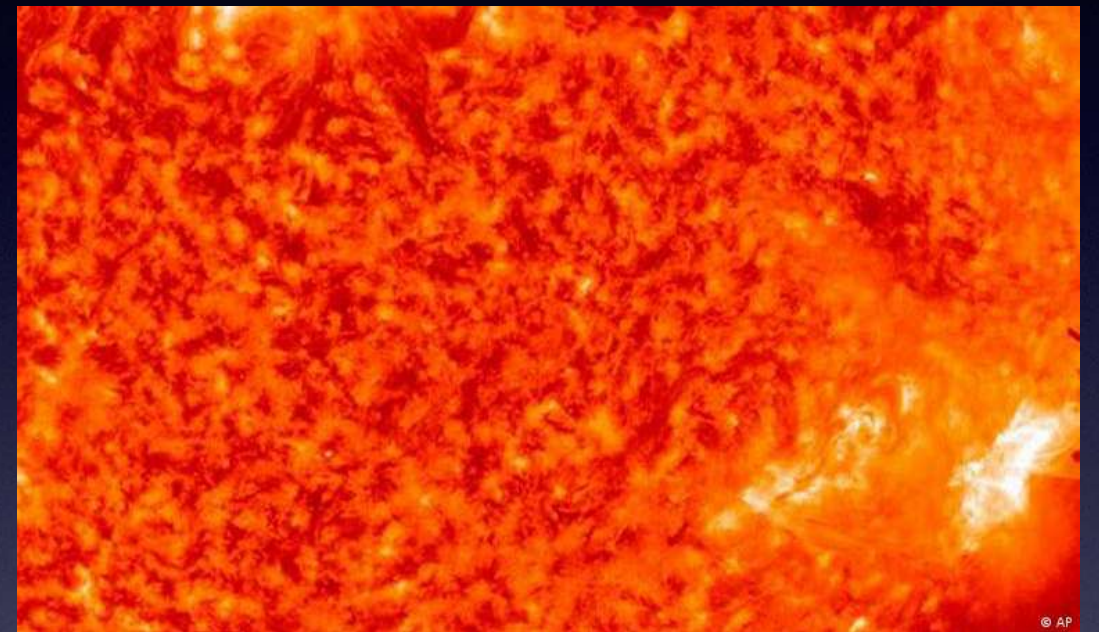
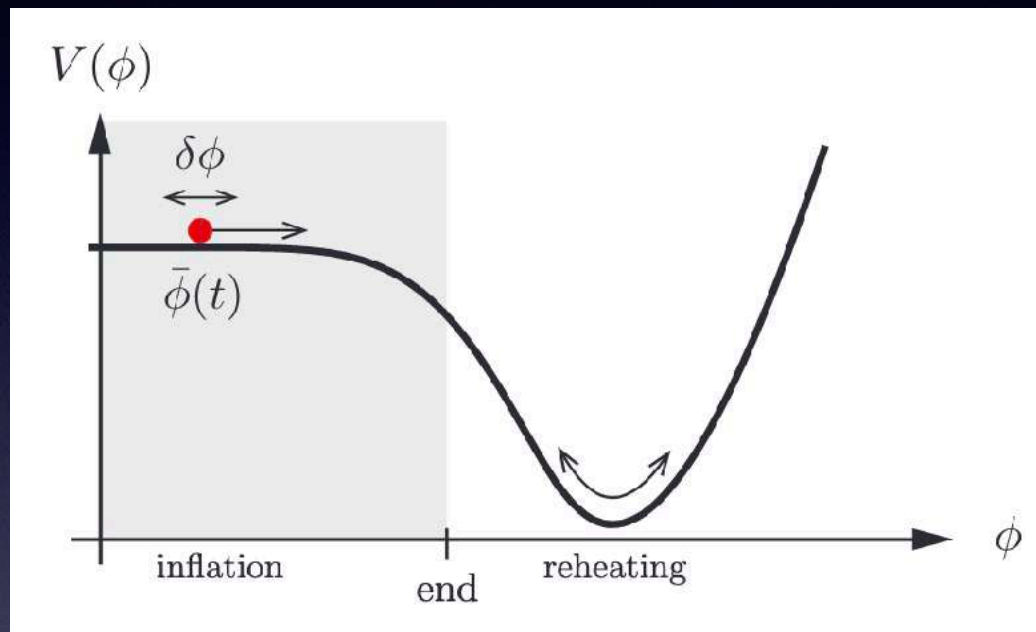




“Preheating”



Starting the Big Bang



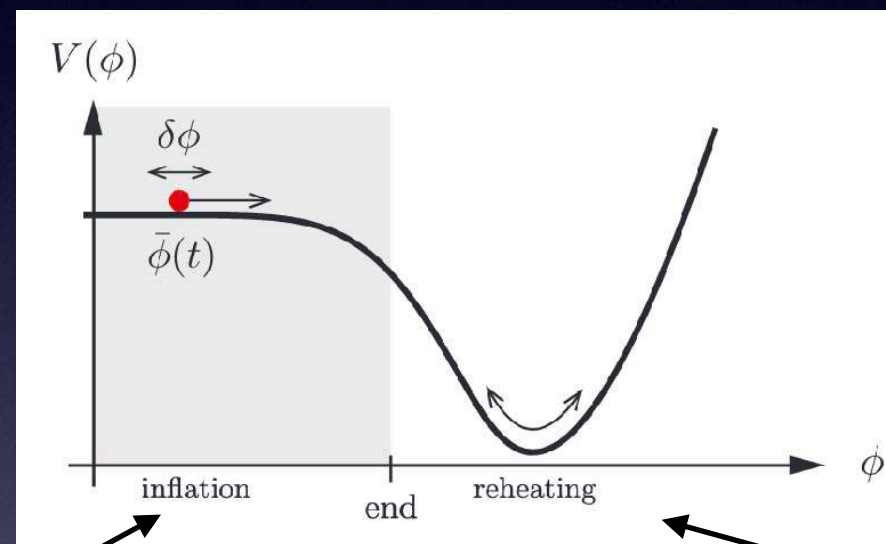
- ▶ Cold $T \approx 0$
- ▶ $\frac{S}{V} \approx 0$
- ▶ Few active d.o.f.

- ▶ Hot ($T \gtrsim T_{\text{BBN}} \sim 10\text{MeV}$)
- ▶ $\frac{S}{V} \propto g_{\text{eff}}(T)T^3$
- ▶ Many active d.o.f. (i.e. particles)

Huge production of entropy (information processing)

Theorist's View of the Early U

$$\mathcal{L}_\phi = -\frac{1}{2}G_{IJ}\partial_\mu\phi_I\partial^\mu\phi_J - V(\phi)$$



During Inflation

- ▶ Subhorizon homogeneity
- ▶ (Small) superhorizon perturbations
- ▶ $\phi(x, t) \rightarrow \bar{\phi}(t)$

End of Inflation

- ▶ $[\delta\hat{\phi}, \delta\hat{\dot{\phi}}] \implies \langle |\delta\phi_k|^2 \rangle, \langle |\delta\dot{\phi}_k|^2 \rangle > 0$
- ▶ $\phi(x, t) \rightarrow \bar{\phi}(t) + \delta\hat{\phi}(x, t)$
- ▶ Variety of instabilities

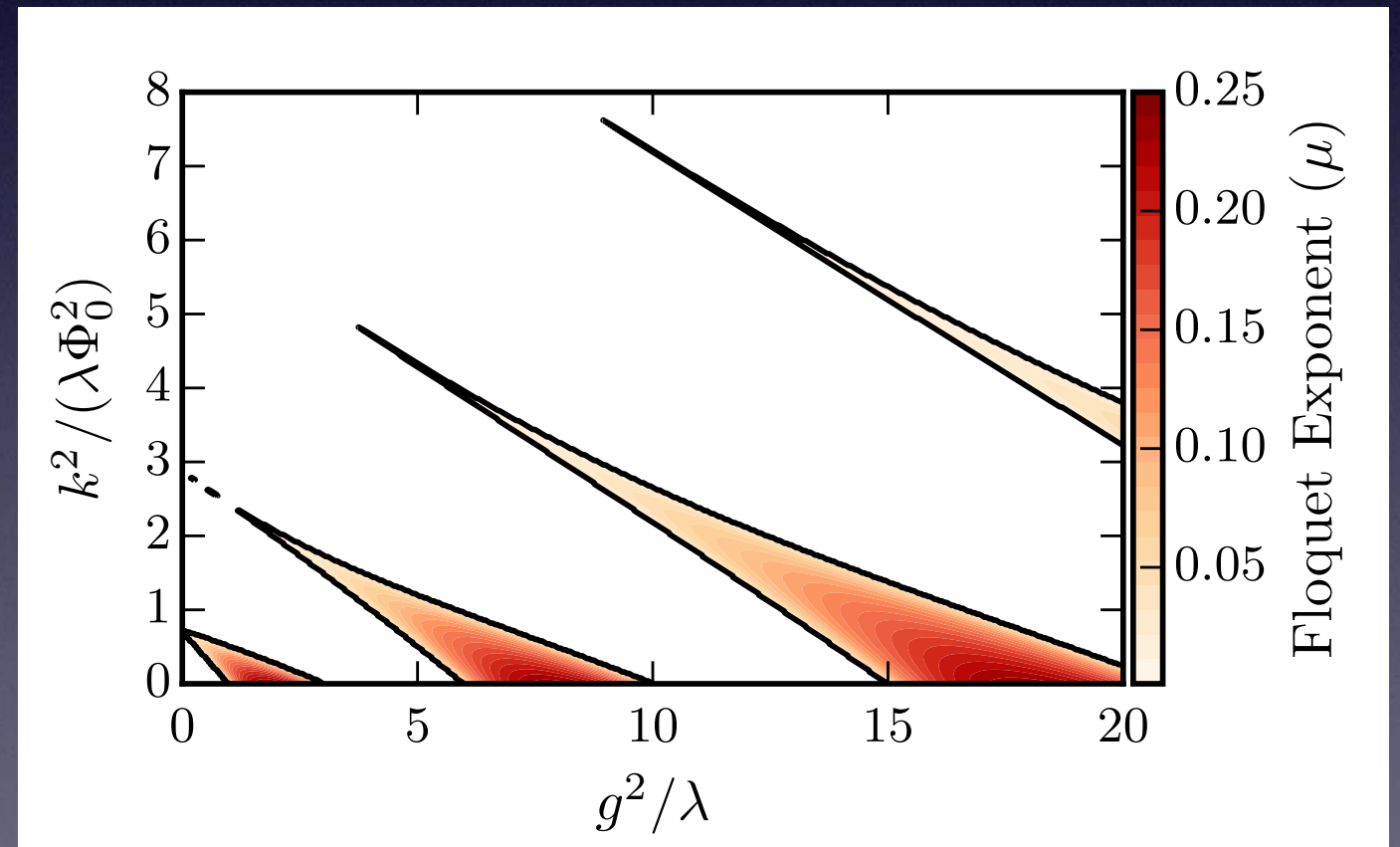
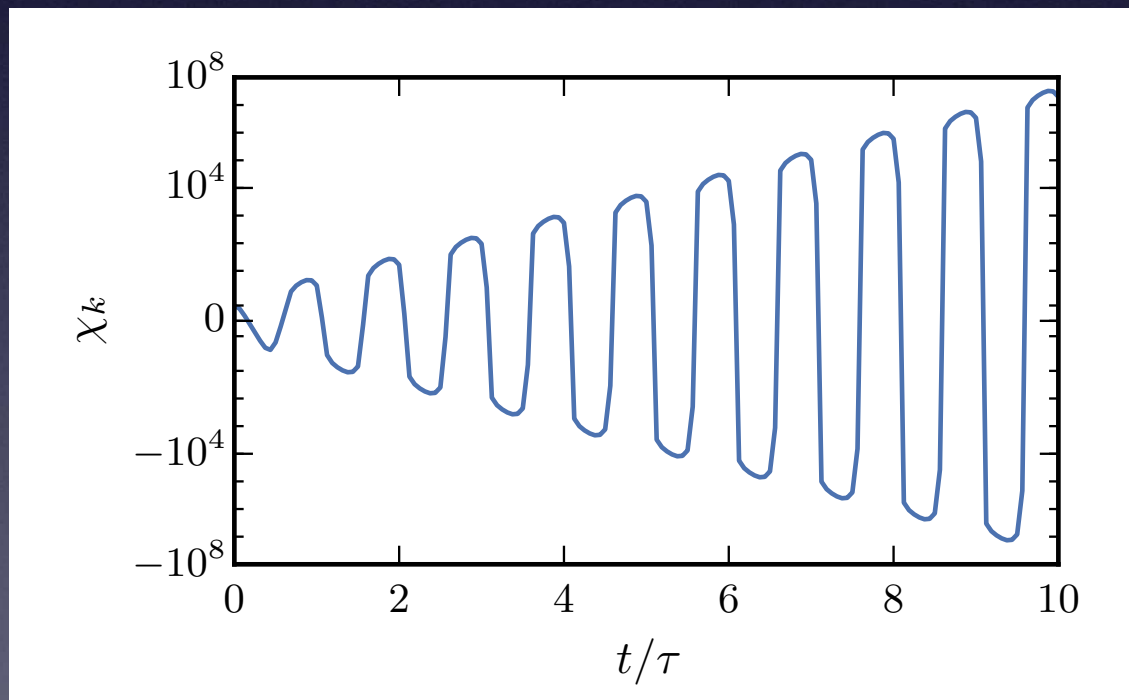
Preheating: Linear Theory

[e.g. Traschen, Bradendberger /
Kofman, Linde, Starobinski]

$$\phi = \bar{\phi}(t) + \delta\phi$$

$$\delta\ddot{\phi} + (k^2 + V''(\bar{\phi})) \delta\phi = 0$$

$$V(\phi) = \frac{\lambda}{4}\phi^4$$



$\bar{\phi}$ acts as external driver

Full treatment includes backreaction and rescattering

Lattice Simulations

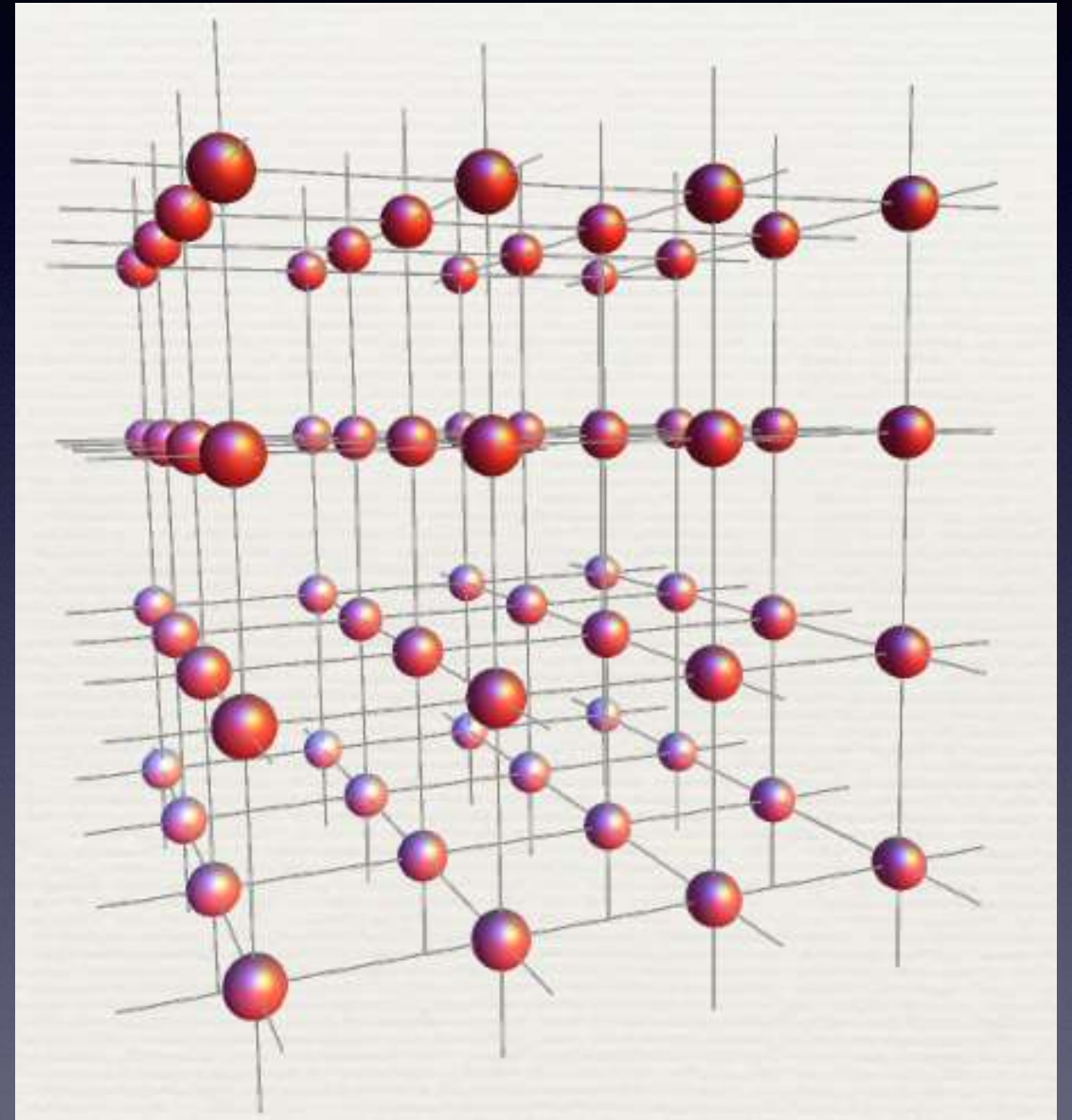
[Braden]

- Solve field equation

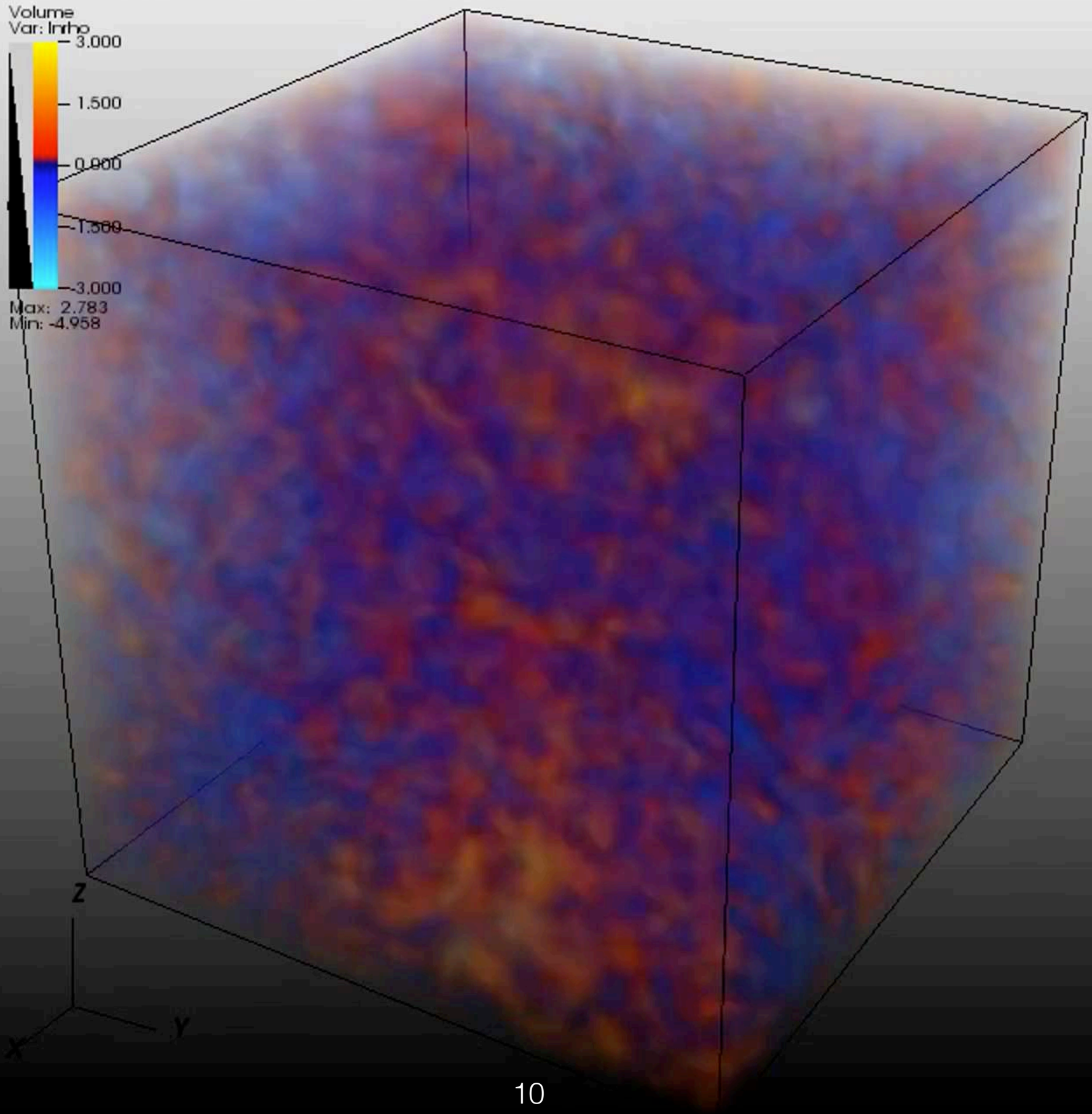
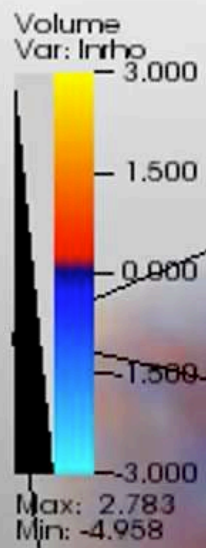
$$\ddot{\phi} + 3H\dot{\phi} + a^{-2}\nabla^2\phi + V'(\phi) = 0$$

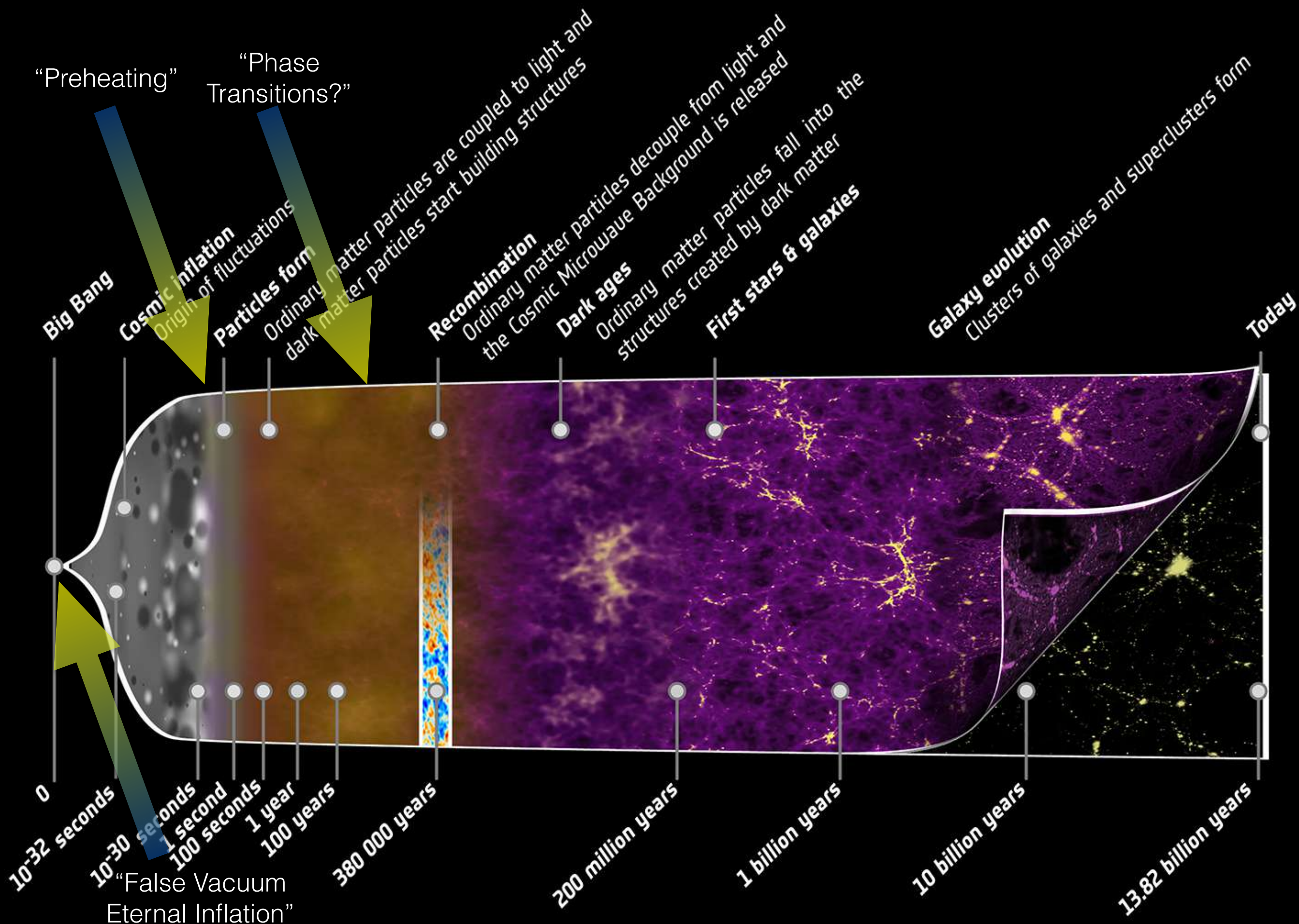
$$H^2 = \frac{\rho}{3M_P^2}$$

- Finite-difference or pseudospectral
- 10th order Gauss-Legendre (general) or 8th order Yoshida (nonlinear sigma model)
- Quantum fluctuations \longrightarrow random field realization

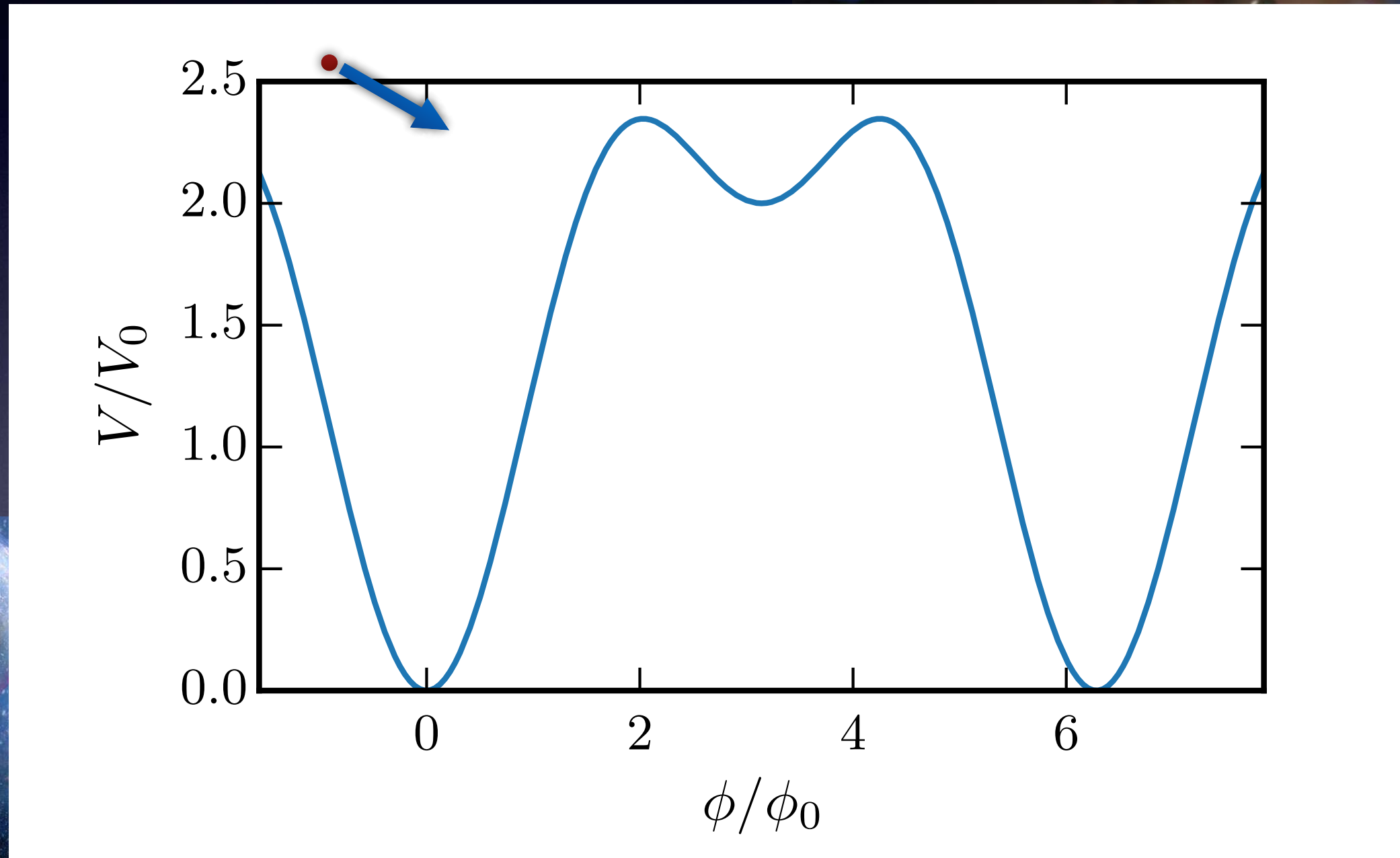


$\mathcal{O}(10^{-15})$ convergence

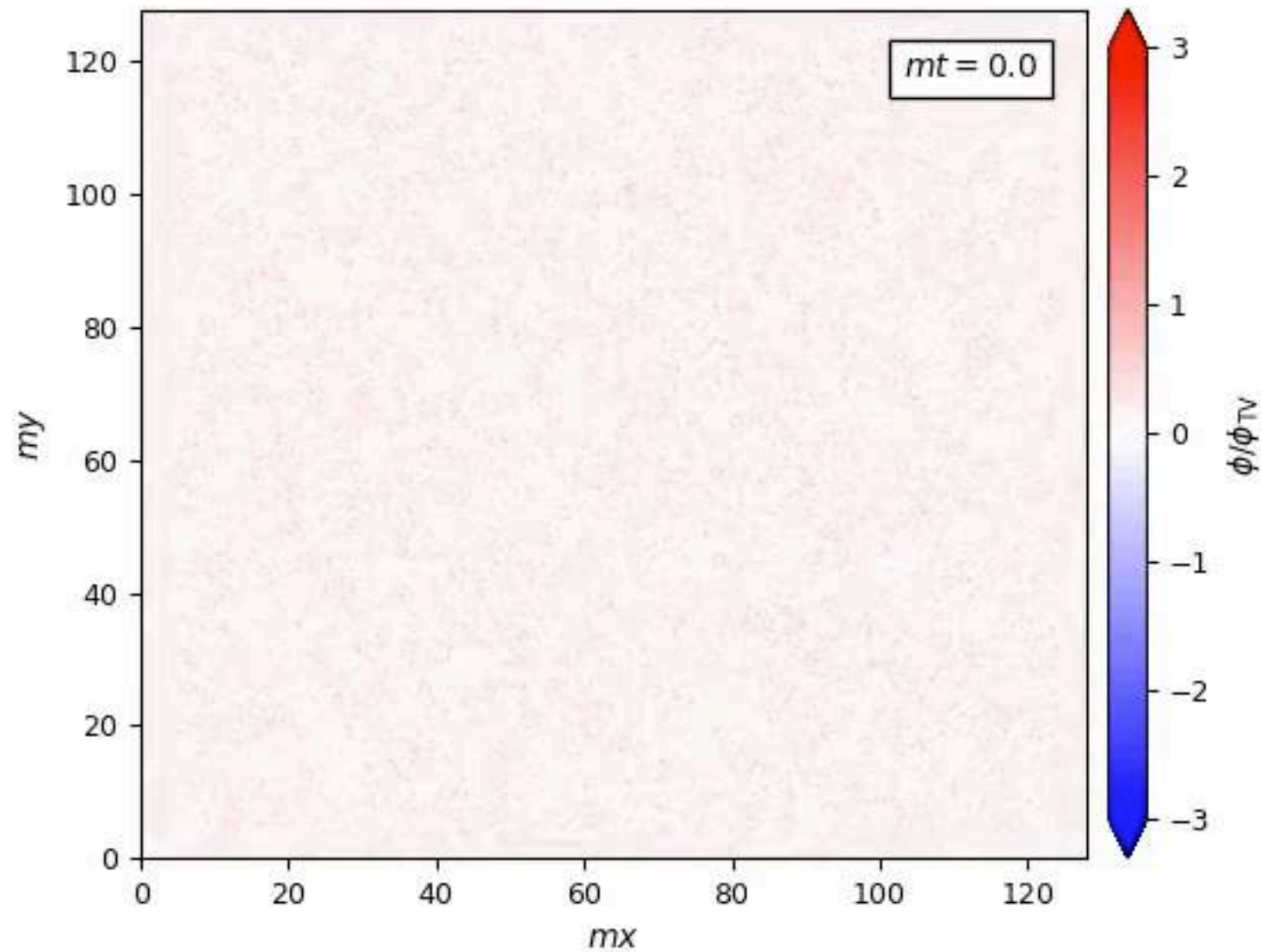




First Order Phase Transitions

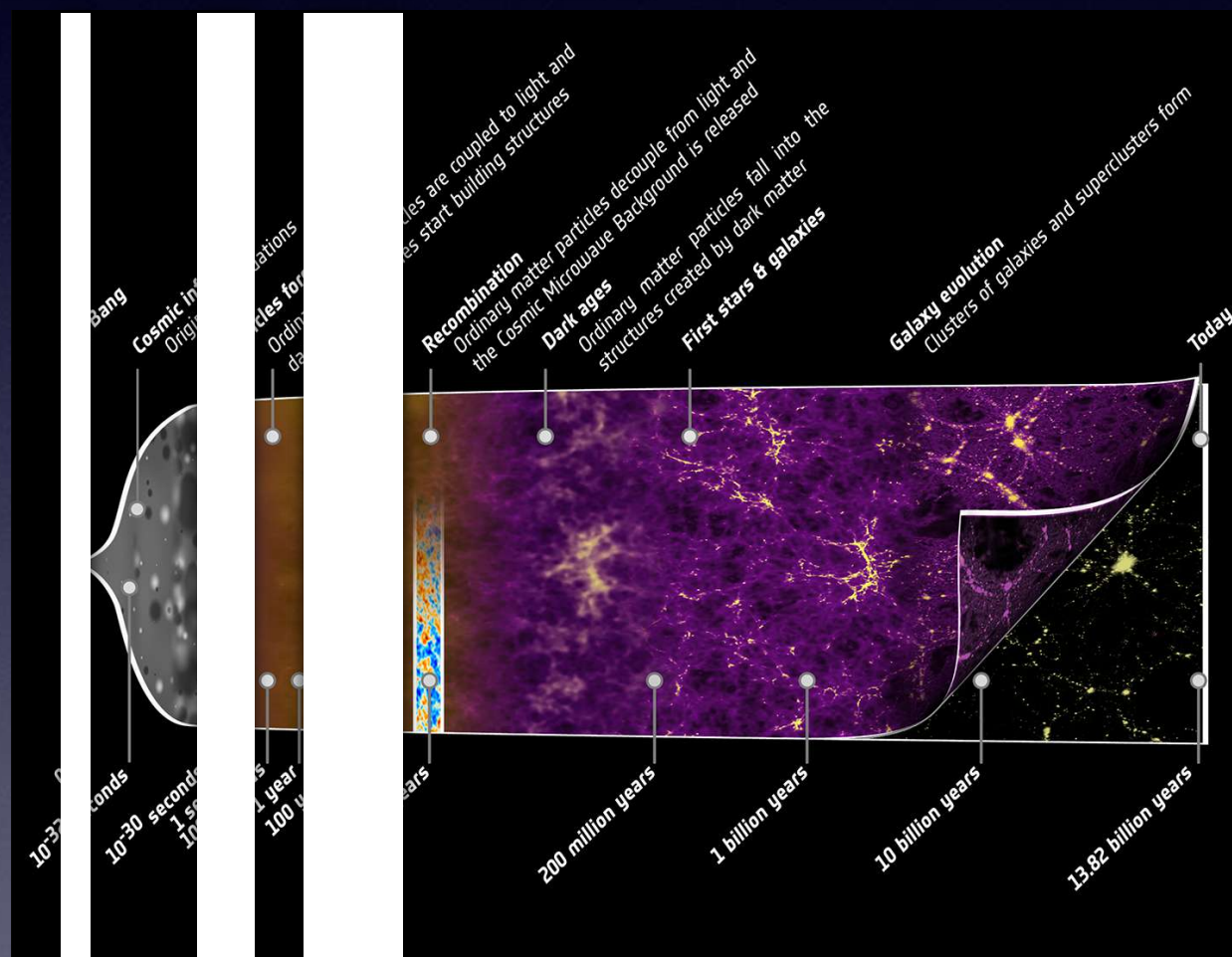


FVD in Action



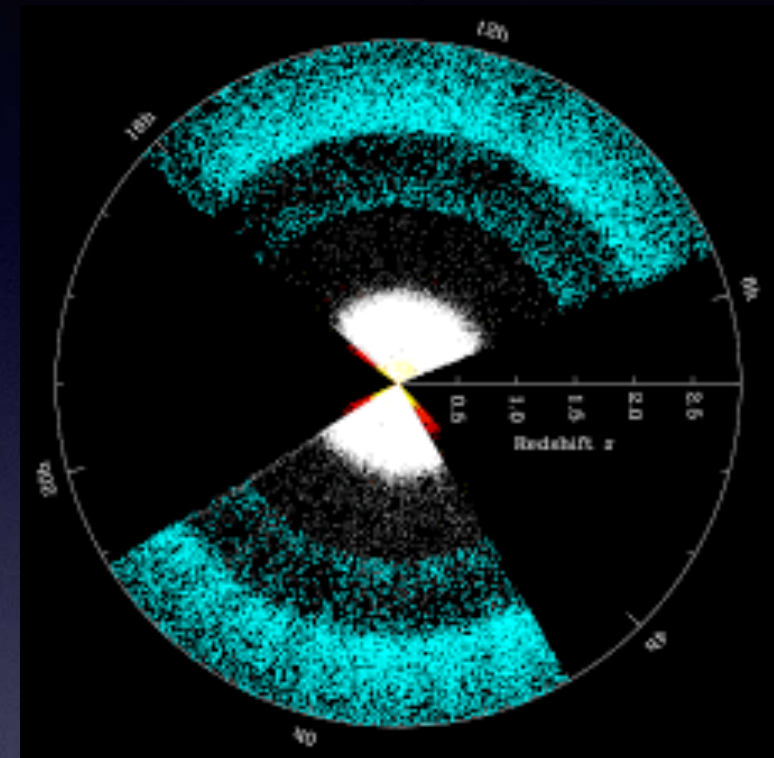
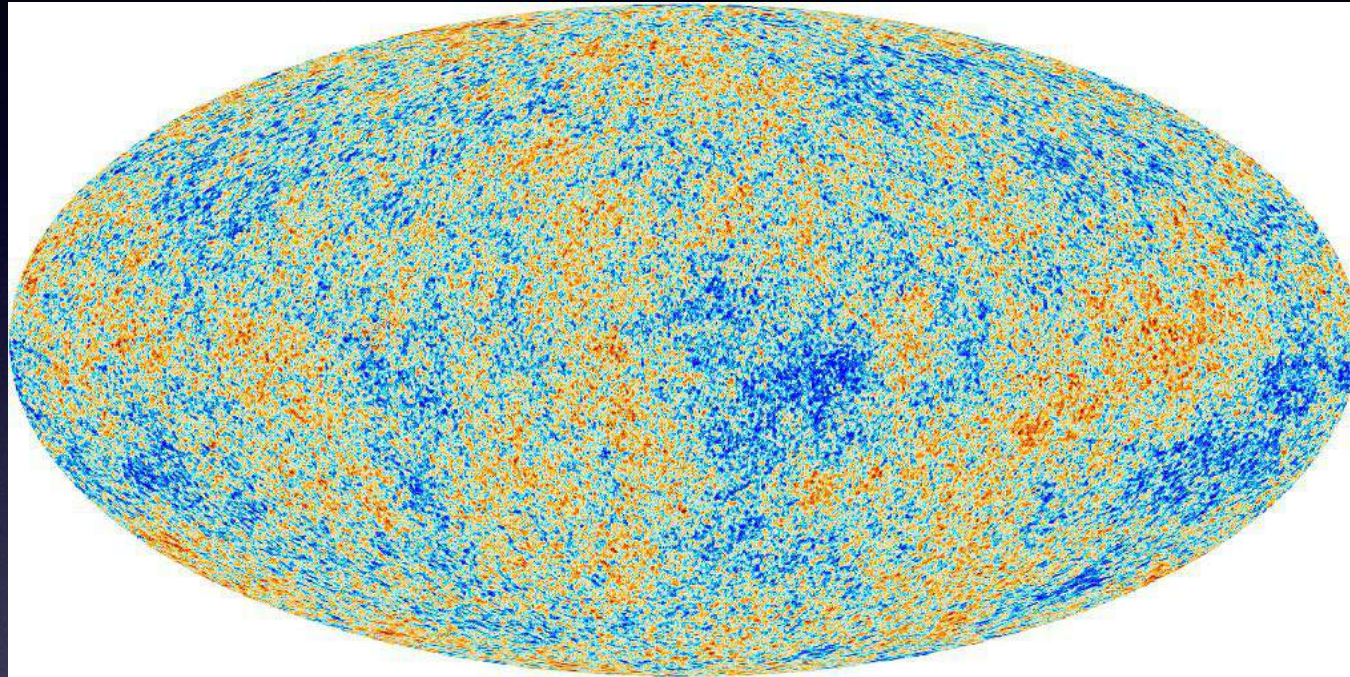
Why Should I Care

- Nonequilibrium physics is fundamental to nature
 - ▶ Intersection of nonlinear QFT and gravity
- Modern cosmology is incomplete otherwise



- Observational Considerations

Inflation and Cosmology



Standard Inflation: Models a few parameters

$$P_s(k) = A_s k^{n_s - 1}$$

$$r = 16\epsilon$$

$$f_{\text{NL}}$$

**Perturbative
NonGaussianity**

What else can we use the data for?

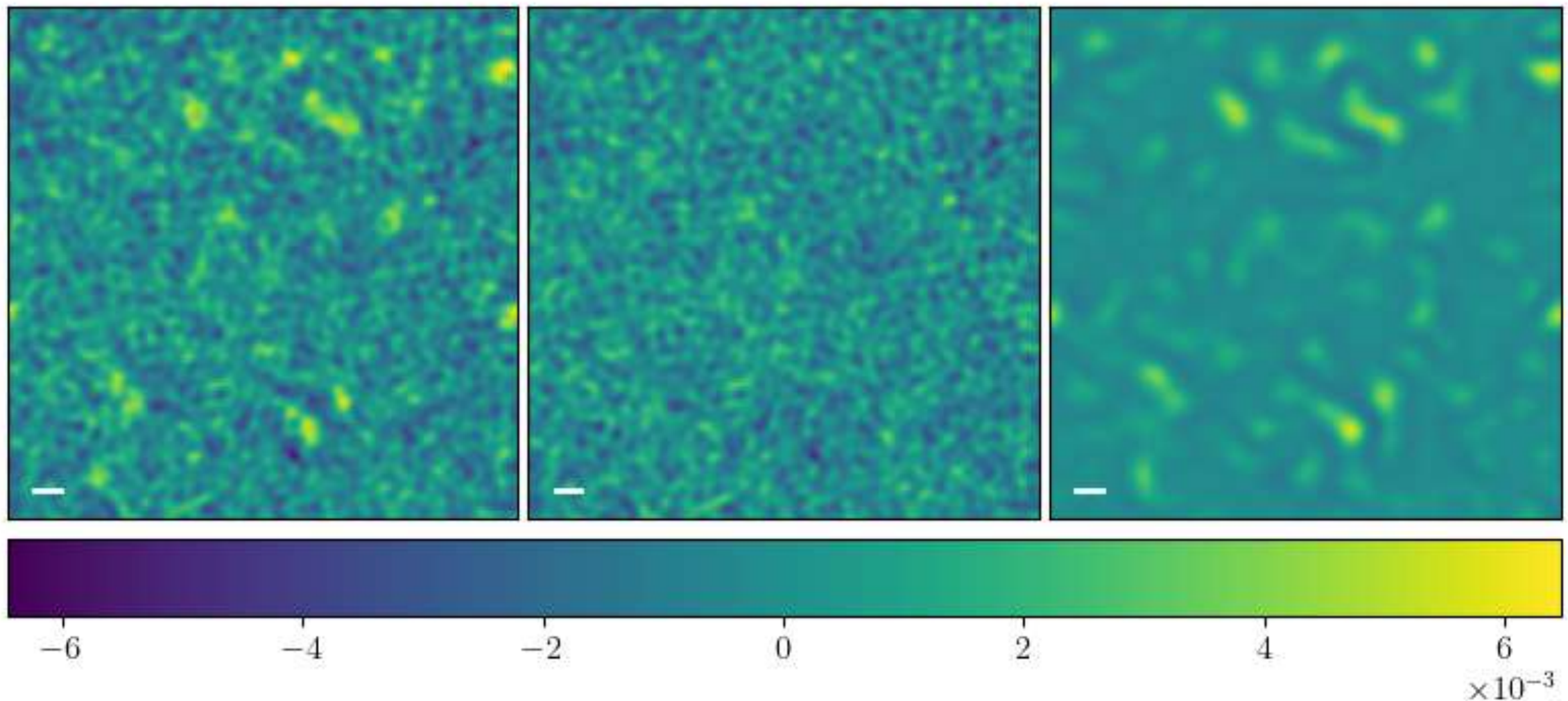
Early U Nonlinearity Imprints Novel Density Fluctuations

$$\zeta = \zeta_G + F_{\text{NL}}(\chi)$$

Total

Gaussian

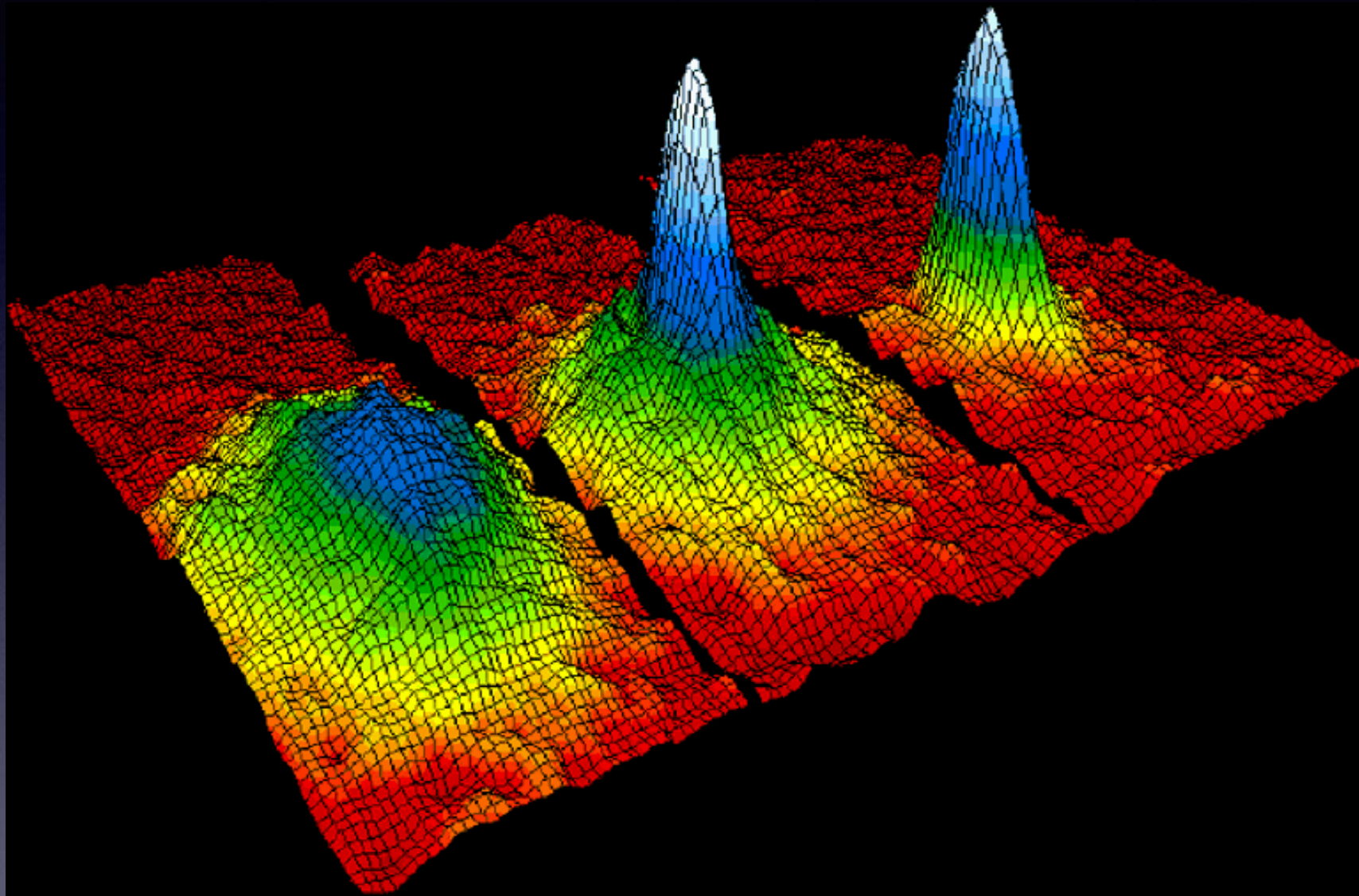
NonGaussian



Nonlinear, Nonperturbative, Nonequilibrium Quantum Field Theory (coupled to gravity)

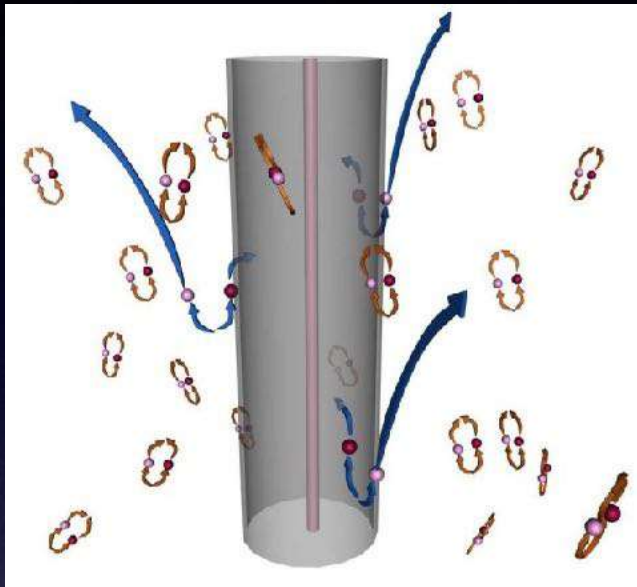
Our understanding of (p)reheating
and false vacuum decay rests
on reasonable but experimentally untested
approximations to non equilibrium QFT

Analogue Systems?

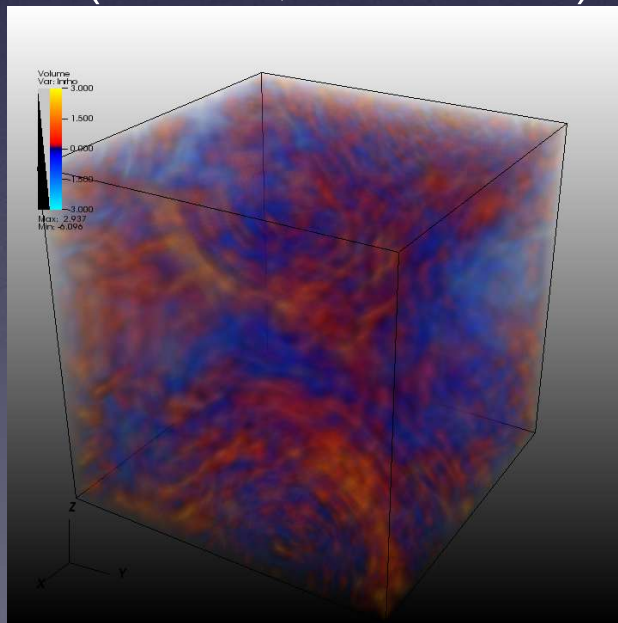


Bose-Einstein Condensates

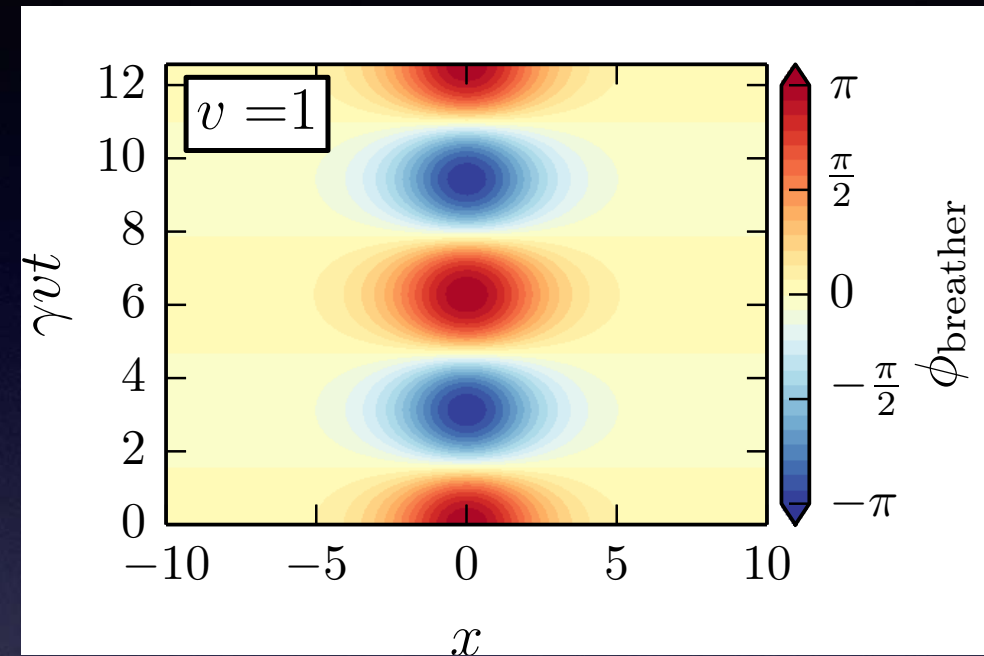
Analogue Systems



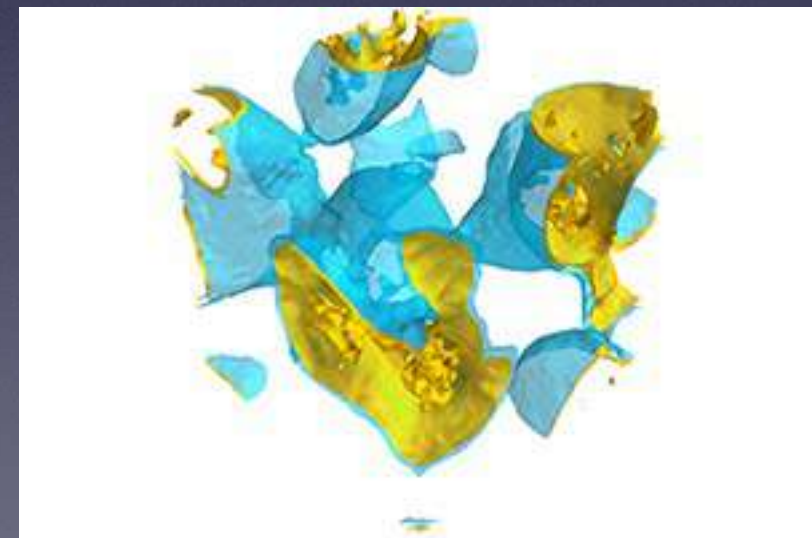
Hawking Radiation
(Linear, Quantum)



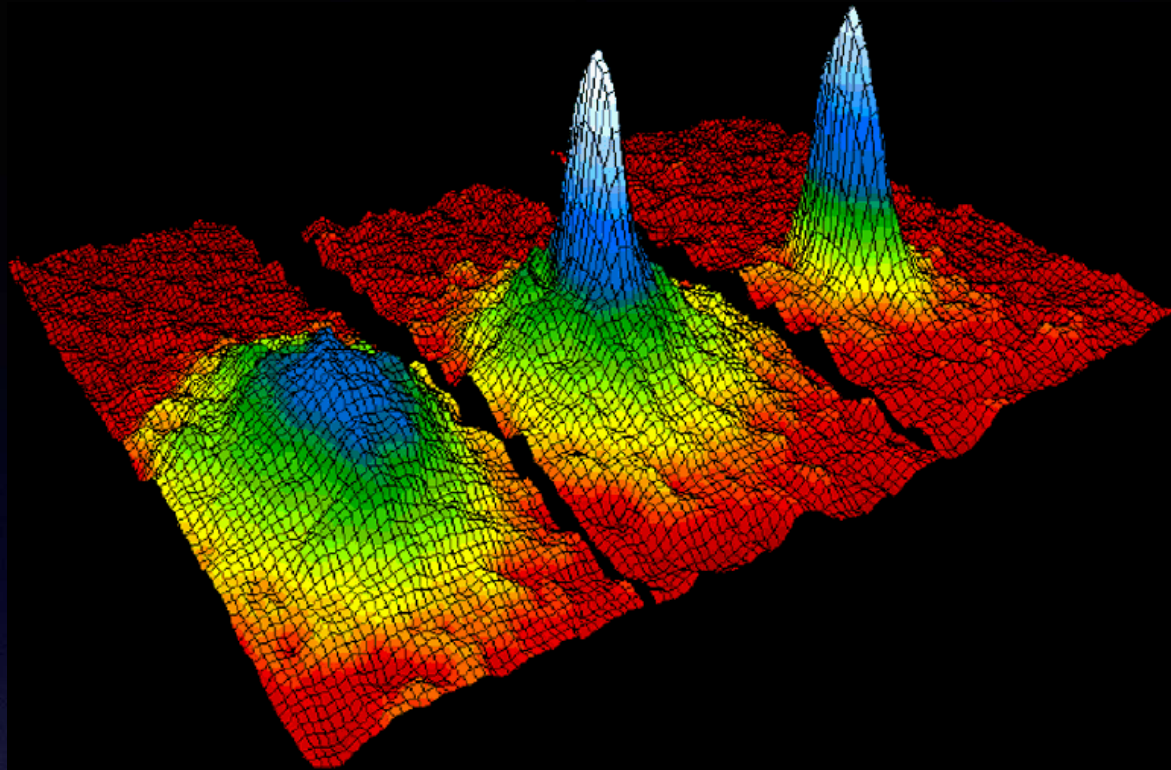
Preheating
(Linear, Quantum) ->
(Nonlinear, Classical)



Solitons
(Nonlinear, Classical)



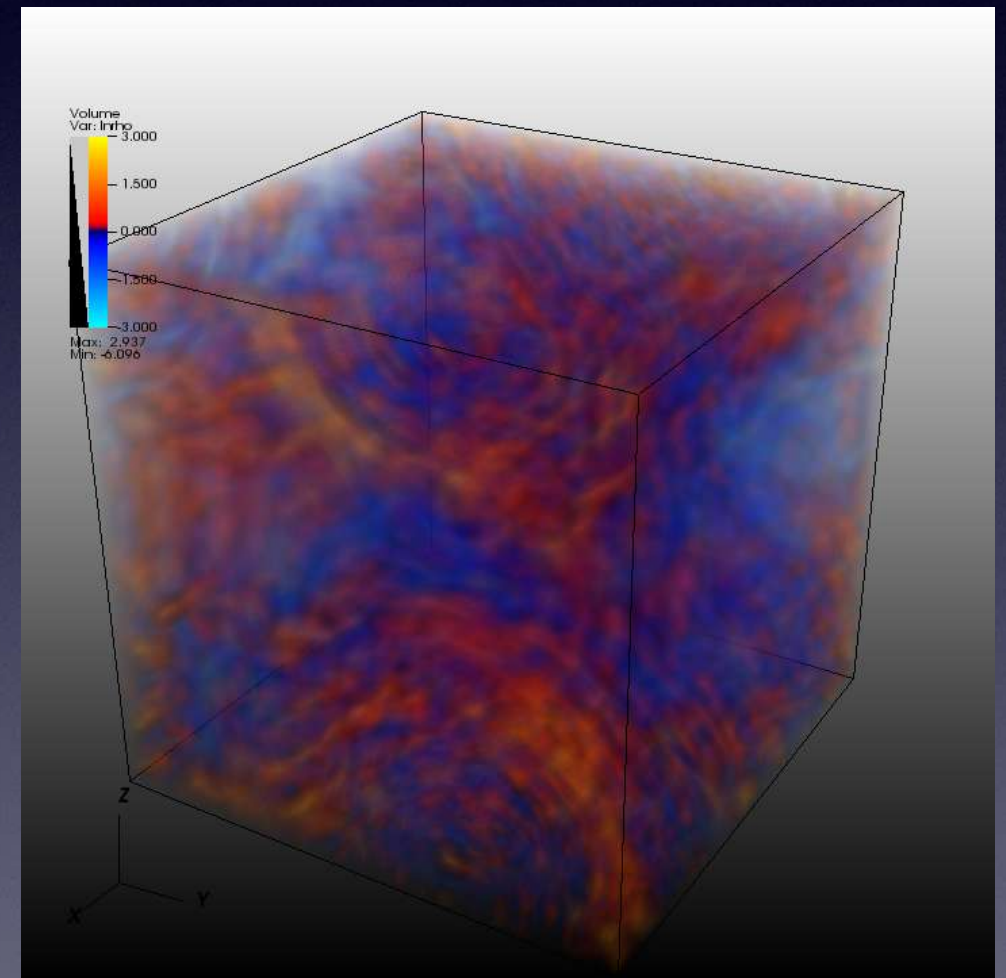
Bubble Nucleation
(Nonlinear, Quantum)



$$\mathcal{L}_{\text{eff}} \sim G(\phi) \frac{\dot{\phi}^2}{2} - c_s^2 \frac{(\nabla \phi)^2}{2} + \nu \Lambda \cos \phi + \dots$$

Finite number effects
 Trapping Potential
 Finite Temperature

$$i\hbar\dot{\psi}_i = \left(-\delta_{ij} \frac{\hbar^2}{2m_i} \nabla^2 + V(\mathbf{x}) + g_{ij} |\psi_j|^2 \right) \psi_i - \nu_{ij} \psi_j$$

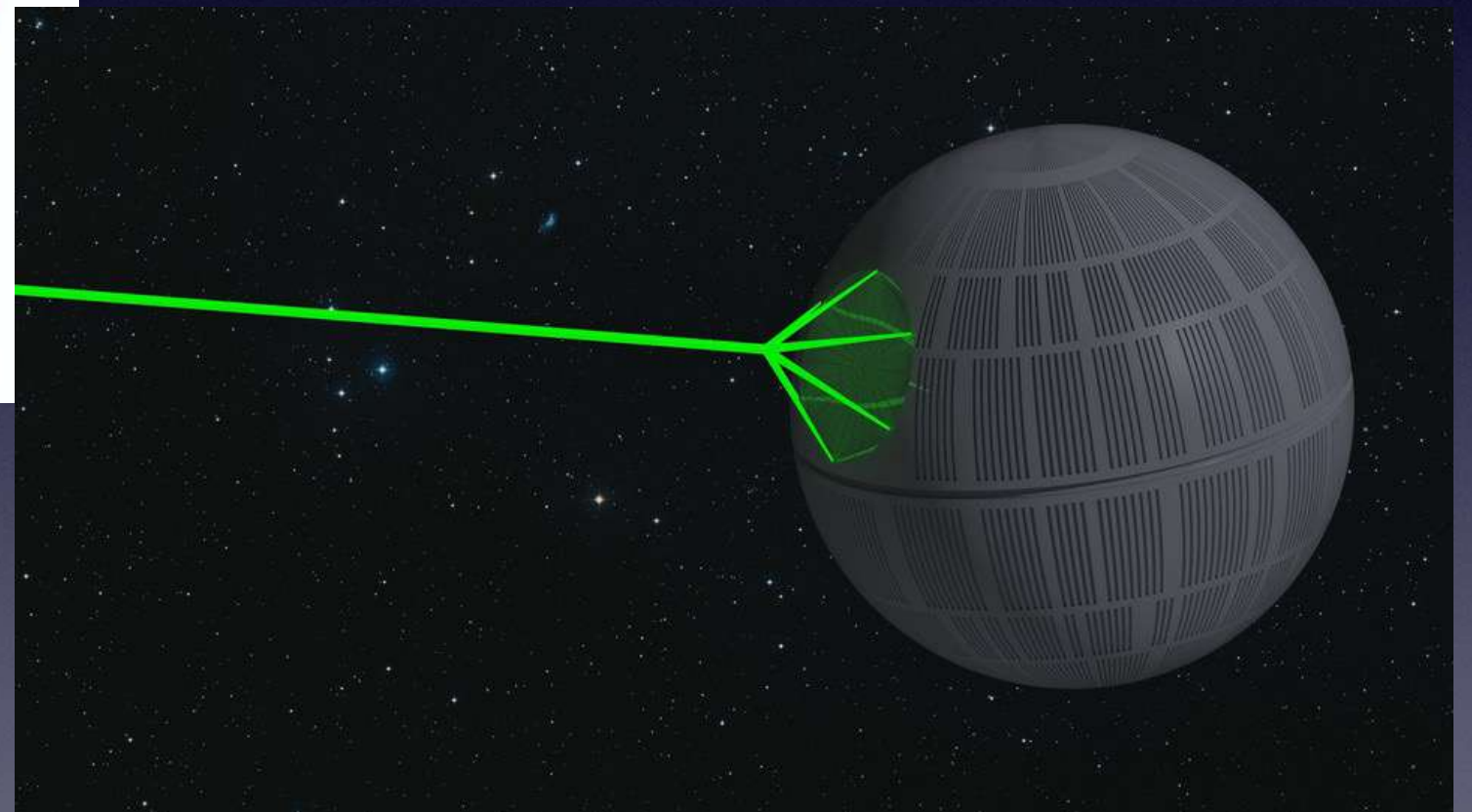
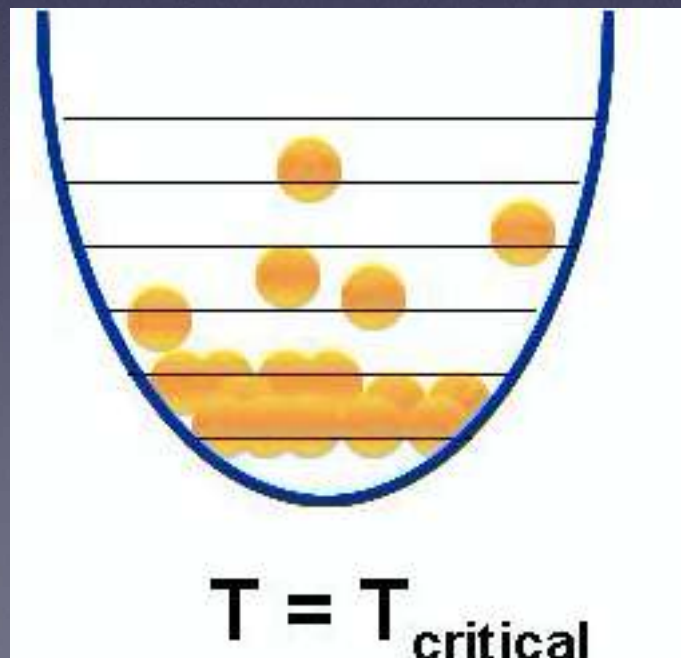
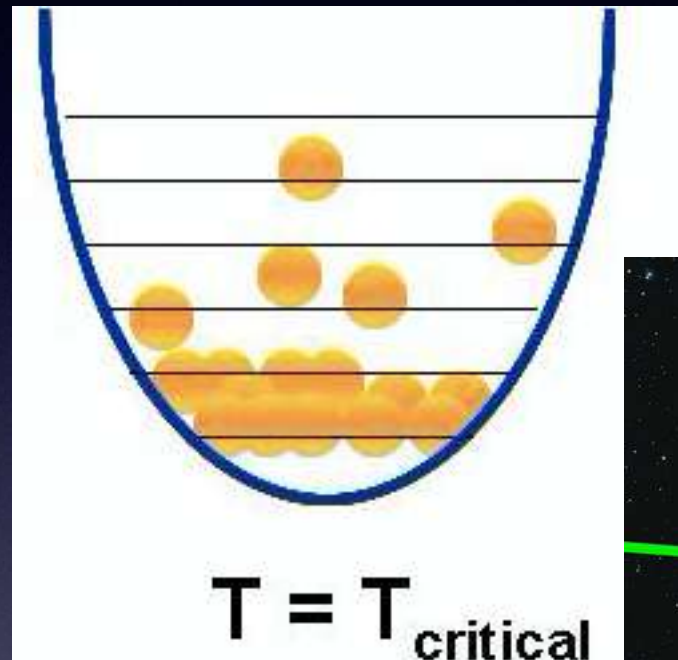


Building an Analogue System

[Fialko, Opanchuk, Sidorov, Drummond, and Brand]

[**JB**, Johnson, Peiris, Weinfurtner]

[Billam, Gregory, Michel, and Moss]



Important Scale : Healing Length

Crossover between wave and particle dispersion relationship

BECs and Relativity

$$\psi_i = \sqrt{\rho_i} e^{-i\phi_i}$$

Can. Momentum
(Particle density)

Can. Position
(Complex phase)

Assumptions

$$\rho_i(x, t) = n_i + \delta\rho_i(x, t)$$

Useful limit $\tilde{\nu} \equiv \frac{\nu}{g\bar{n}} \ll 1$

1) Homogeneous

2) Small

First Hint: Relativistic Dispersion for $k < k_{\text{heal}}$

$$\hbar^2\omega^2 = m^2 + c^2k^2 \left(1 + \frac{k^2}{k_{\text{heal}}^2} \right) \quad k_{\text{heal}}^2 = \frac{4mg\bar{n}}{\hbar^2} = 4 \frac{m}{g\bar{n}} \frac{g^2\bar{n}^2}{\hbar^2}$$

Small Density Fluctuations

A convenient variable is $\varphi = \phi_2 - \phi_1$

Integrate out fluctuations in number $Z_{\text{eff}} \propto \int d\phi e^{i\mathcal{L}_{\text{eff}}}$

Relative phase φ is described by sine-Gordon model

$$\mathcal{L}_{\text{eff}} = \frac{v}{2} (\partial_t \varphi)^2 + v\Lambda \cos \varphi + \dots$$

$$c_s^2 \approx \frac{g'}{m}$$

$$m_\varphi \approx \sqrt{\tilde{v}} \frac{2g\bar{n}}{\hbar}$$

$$L_\varphi = \frac{c_s}{m_\varphi} \sim \frac{L_{\text{heal}}}{\sqrt{\tilde{v}}}$$

A Small Detour: False Vacuum Decay

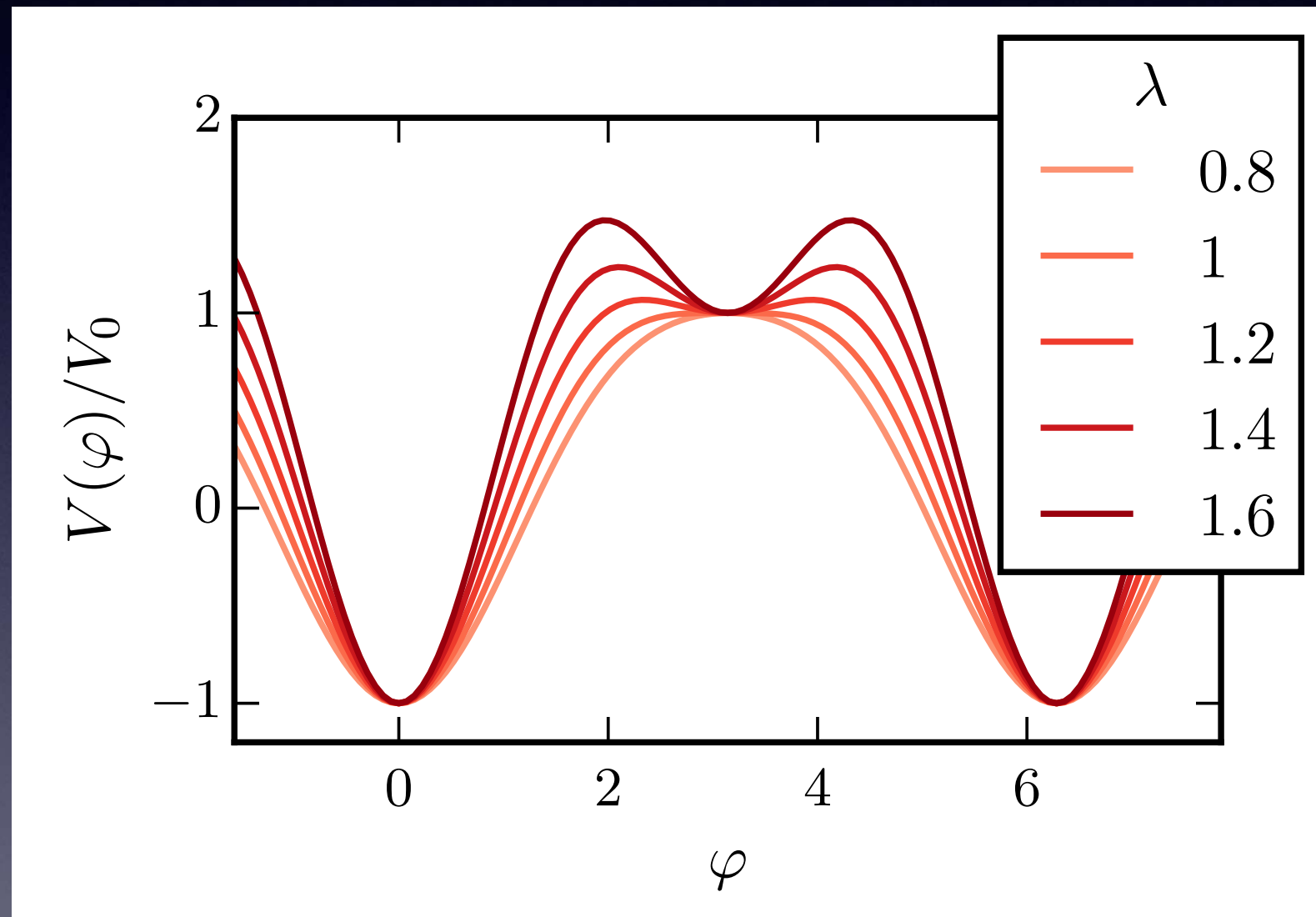
Modulate Transition Rate

$$\nu = \nu_0 + \delta\hbar\omega \cos(\omega t)$$



Time Averaged Potential

$$\lambda = \delta \left(\frac{2g\bar{n}}{\nu_0} \right)^{1/2}$$



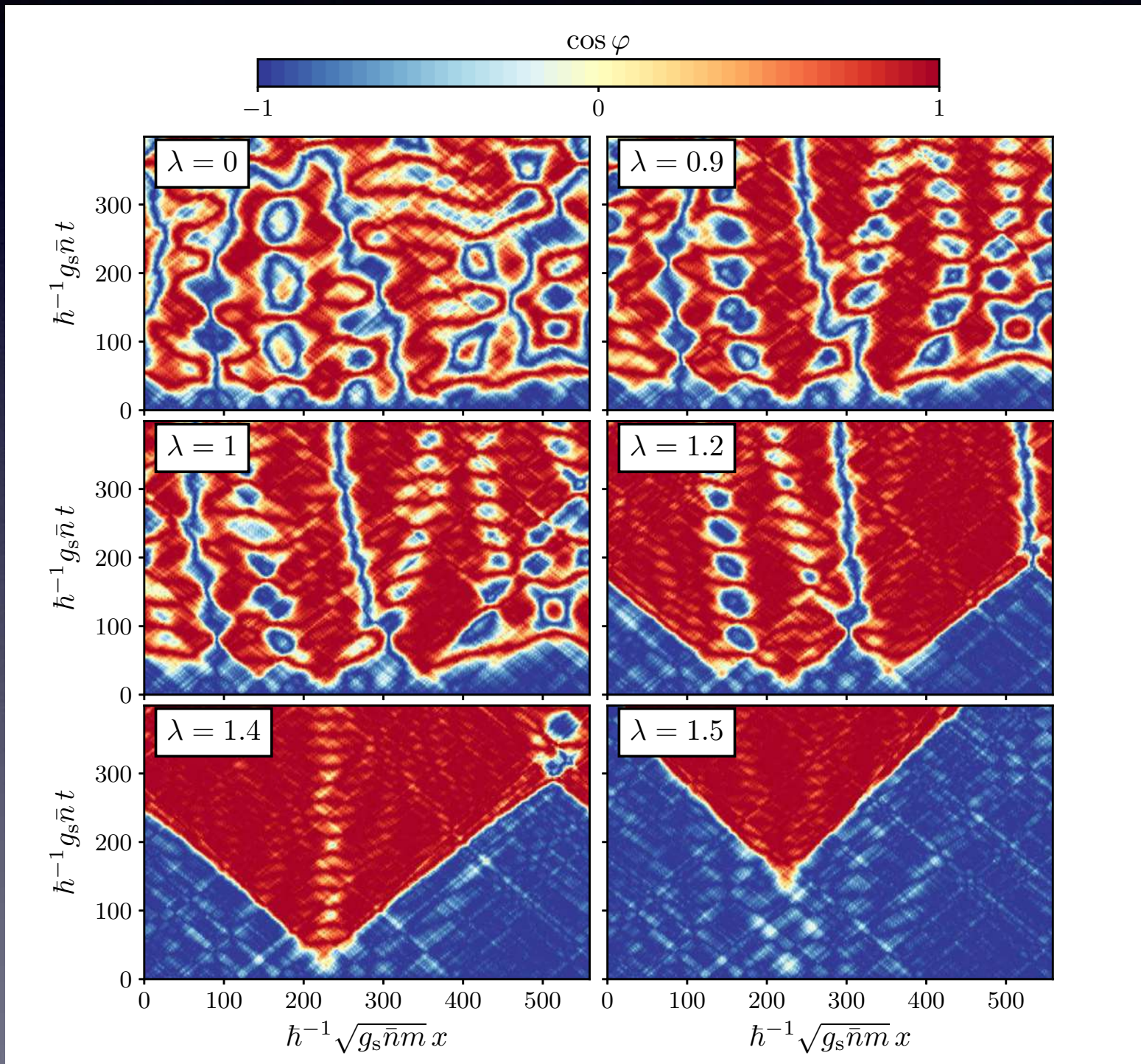
$$V(\phi) = V_0 \left(-\cos \left(\frac{\phi}{\phi_0} \right) + \frac{\lambda^2}{2} \sin^2 \left(\frac{\phi}{\phi_0} \right) + 1 \right)$$

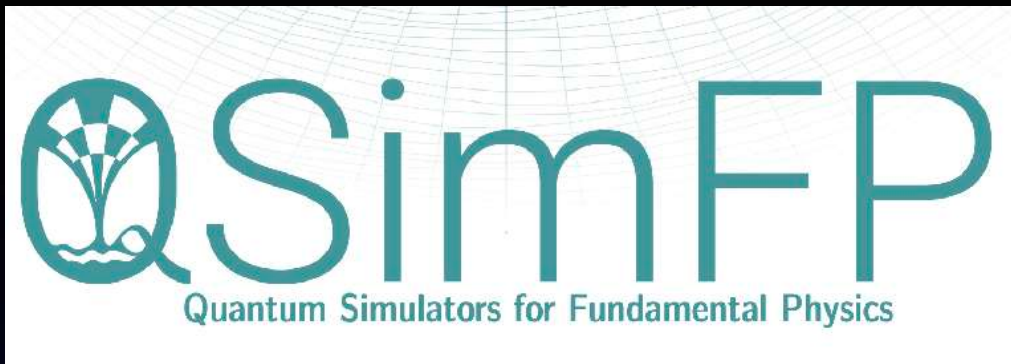
BEC Evolution

2nd-Order
Phase Transition

Rapid 1st-Order
Phase Transition

Slower 1st-Order
Phase Transition





Building an
analog FVD
experiment
@
Cambridge

More info at
www.qsimfp.org



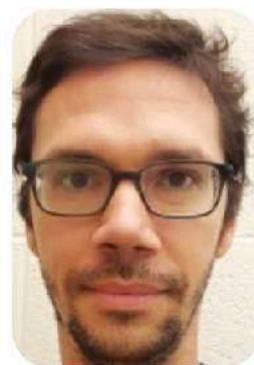
Vanessa Augustus



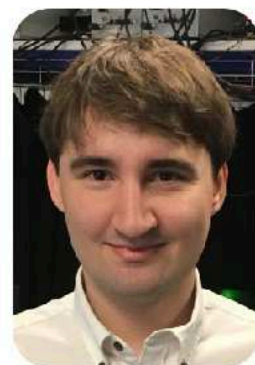
Carlo Barenghi



Thomas Billam



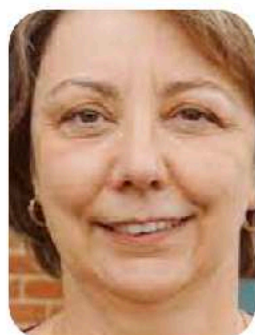
Jonathan Braden



Christoph Eigen



Sebastian Erne



Ruth Gregory



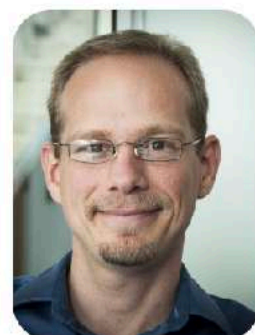
Zoran Hadzibabic



Gregoire Ithier



Alexander Jenkins



Matthew Johnson



Anthony Kent



Friedrich Koenig



Jorma Louko



Ian Moss



John Owers-Bradley



Hiranya Peiris



Andrew Pontzen



Radivoje Prizia



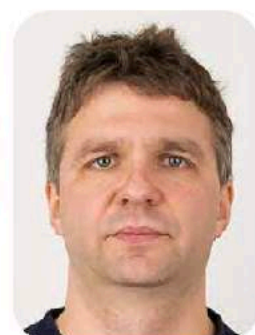
Xavier Rojas



Joerg Schmiedmayer



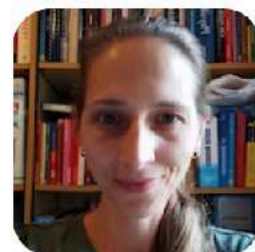
Theo Torres



Viktor Tsepelin



William G. Unruh



Silke Weinfurtner



Patrik Švančara

Back to our Regularly Scheduled Programming: Preheating

Small Density Fluctuations

A convenient variable is $\varphi = \phi_2 - \phi_1$

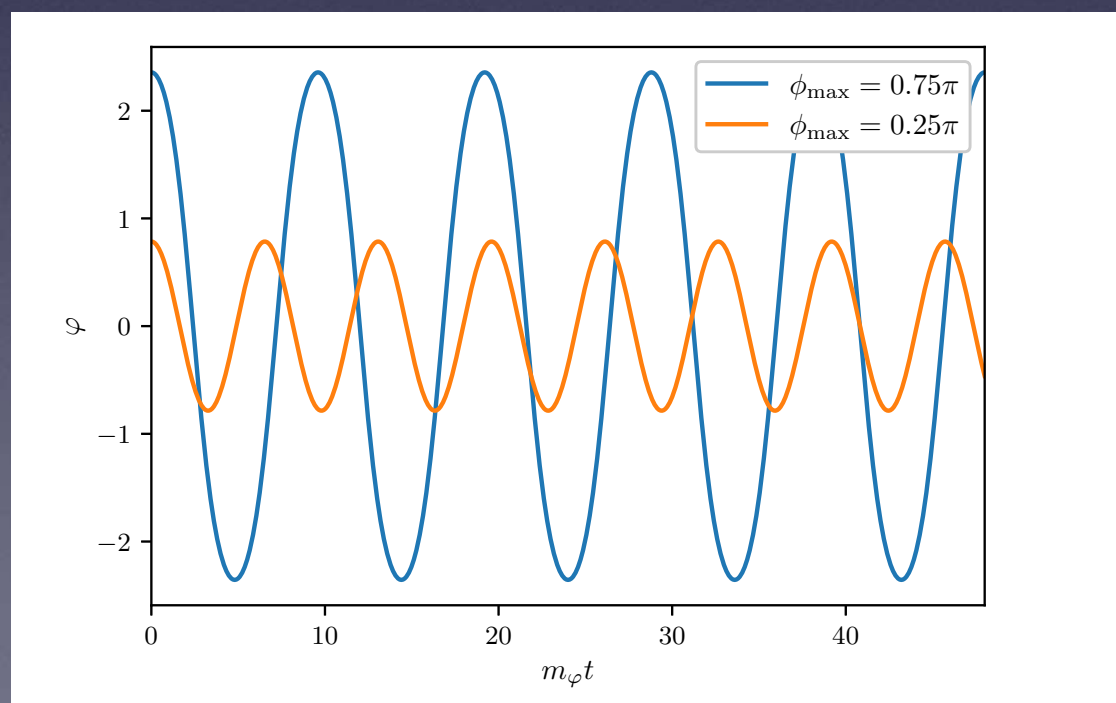
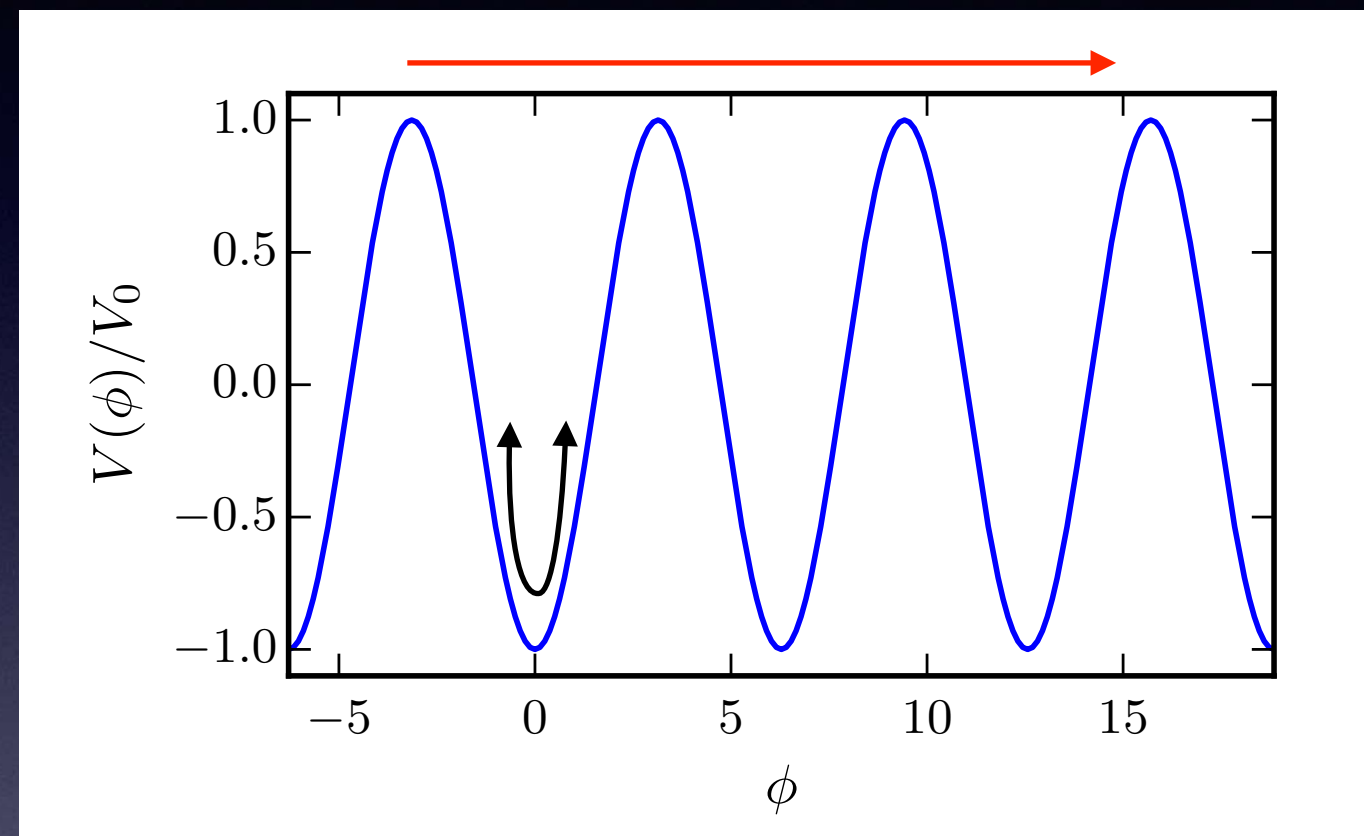
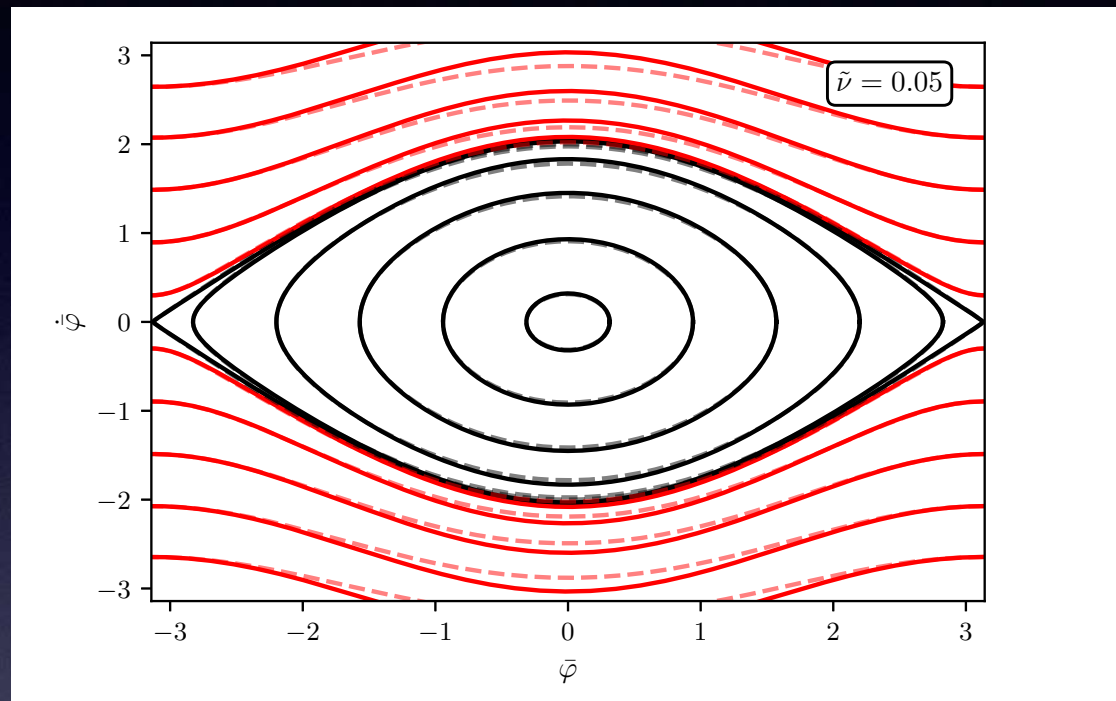
Integrate out fluctuations in number density $Z_{\text{eff}} \propto \int d\phi e^{i\mathcal{L}_{\text{eff}}}$

Relative phase governed by sine-Gordon model

$$\mathcal{L}_{\text{eff}} \sim \frac{\dot{\varphi}^2}{2} - c_s^2 \frac{(\nabla \varphi)^2}{2} + \nu \Lambda \cos \varphi + \dots$$

$$c_s^2 \approx \frac{g\bar{n}}{m} \quad m_\varphi \approx \sqrt{\tilde{\nu}} \frac{2g\bar{n}}{\hbar} \quad L_\varphi = \frac{c_s}{m_\varphi} \sim \frac{L_{\text{heal}}}{\sqrt{\tilde{\nu}}}$$

Homogeneous Background Evolution



Two periodic motion types
Black - single minimum
Red - scan minima

Perturbation Equations

$$\vartheta = \theta_1 + \theta_2 \quad \psi_i(\phi_1, t) \phi_2 \bar{\psi}_i(t) \Pi_\vartheta = \frac{\rho_1 + \rho_2}{2\hbar\sqrt{\tilde{\nu}}} \quad \Pi_\varphi = \frac{\rho_2 - \rho_1}{2\hbar\sqrt{\tilde{\nu}}}$$

$$\begin{aligned} \frac{d\delta\Pi_\vartheta}{d\tilde{t}} &= \kappa^2 \left(\delta\vartheta + \sqrt{\tilde{\nu}}\bar{\Pi}\delta\varphi \right) \\ \frac{d\delta\vartheta}{d\tilde{t}} &= -\frac{\tilde{\nu}}{1-\tilde{\nu}\bar{\Pi}^2}\kappa^2 \left(\delta\Pi_\vartheta - \sqrt{\tilde{\nu}}\bar{\Pi}\delta\Pi_\varphi \right) - \delta\Pi_\vartheta \\ &\quad - \sqrt{\frac{\tilde{\nu}}{1-\tilde{\nu}\bar{\Pi}^2}} \sin\bar{\varphi}\delta\varphi + \left(\frac{\tilde{\nu}}{1-\tilde{\nu}\bar{\Pi}^2} \right)^{3/2} \bar{\Pi} \cos\bar{\varphi} \left(\delta\Pi_\varphi - \sqrt{\tilde{\nu}}\bar{\Pi}\delta\Pi_\vartheta \right) \\ \frac{d\Pi_\varphi}{d\tilde{t}} &= \kappa^2 \left(\delta\varphi + \sqrt{\tilde{\nu}}\bar{\Pi}\delta\vartheta \right) \\ &\quad + \cos\bar{\varphi}\sqrt{1-\tilde{\nu}\bar{\Pi}^2}\delta\varphi + \frac{\sqrt{\tilde{\nu}}\sin\bar{\varphi}}{\sqrt{1-\tilde{\nu}\bar{\Pi}^2}} \left(\delta\Pi_\vartheta - \sqrt{\tilde{\nu}}\bar{\Pi}\delta\Pi_\varphi \right) \\ \frac{d\delta\varphi}{d\tilde{t}} &= -\frac{\tilde{\nu}}{1-\tilde{\nu}\bar{\Pi}^2}\kappa^2 \left(\delta\Pi_\varphi - \sqrt{\tilde{\nu}}\bar{\Pi}\delta\Pi_\vartheta \right) - \delta\Pi_\varphi \\ &\quad + \frac{\tilde{\nu}\bar{\Pi}\sin\bar{\varphi}}{\sqrt{1-\tilde{\nu}\bar{\Pi}^2}}\delta\varphi - \frac{\tilde{\nu}\cos\bar{\varphi}}{(1-\tilde{\nu}\bar{\Pi}^2)^{3/2}} \left(\delta\Pi_\varphi - \sqrt{\tilde{\nu}}\bar{\Pi}\delta\Pi_\vartheta \right). \end{aligned}$$

$$\frac{d\mathbf{y}}{dt} = L(t)\mathbf{y}$$

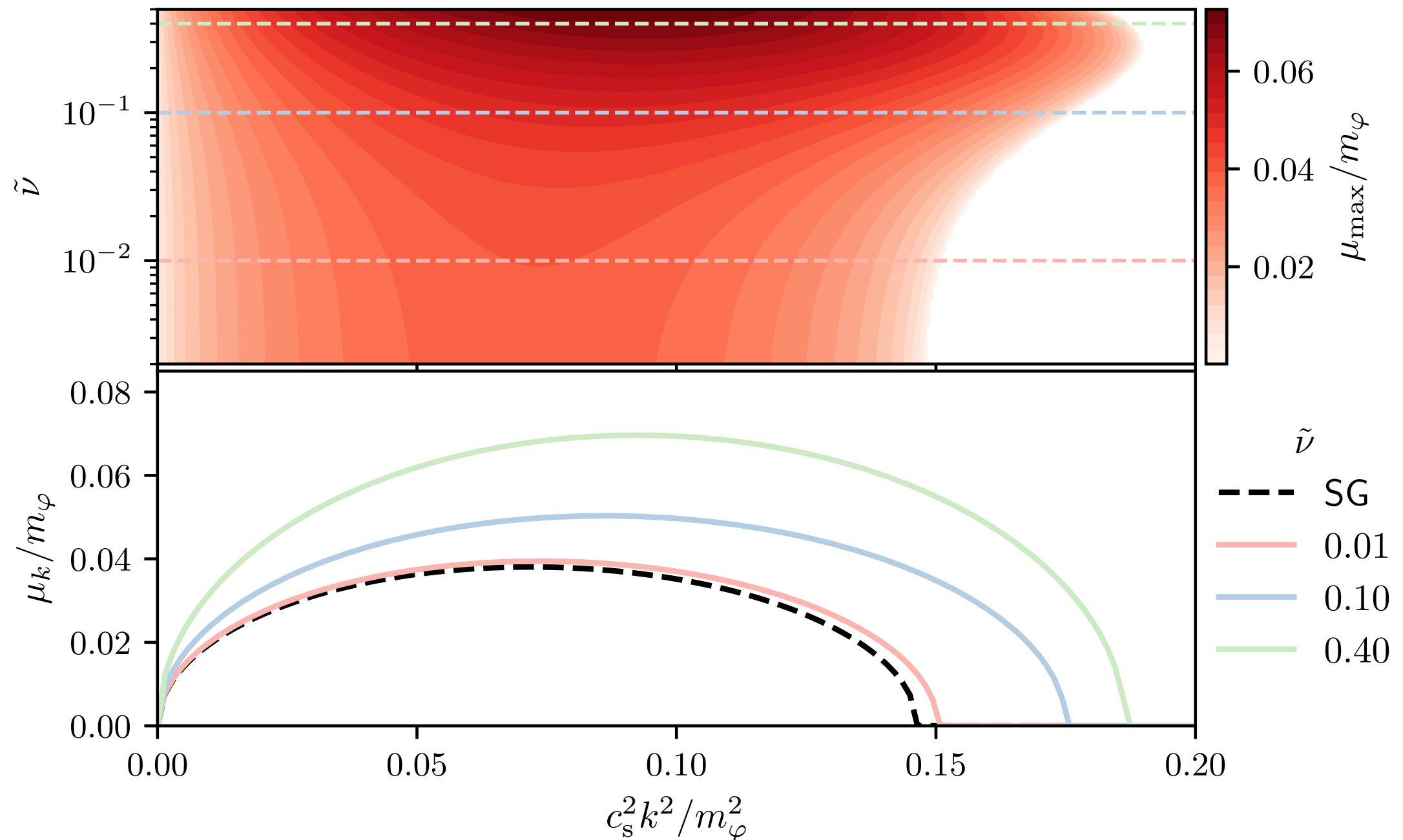
$$L(t+T) = L(t)$$

$\text{Re}(\mu) > 0$:
exponential growth

Periodic $\bar{\varphi}$ and $\bar{\Pi}_\varphi \implies$ solutions $\delta\mathbf{y} = \mathbf{P}(t)e^{\mu t}$

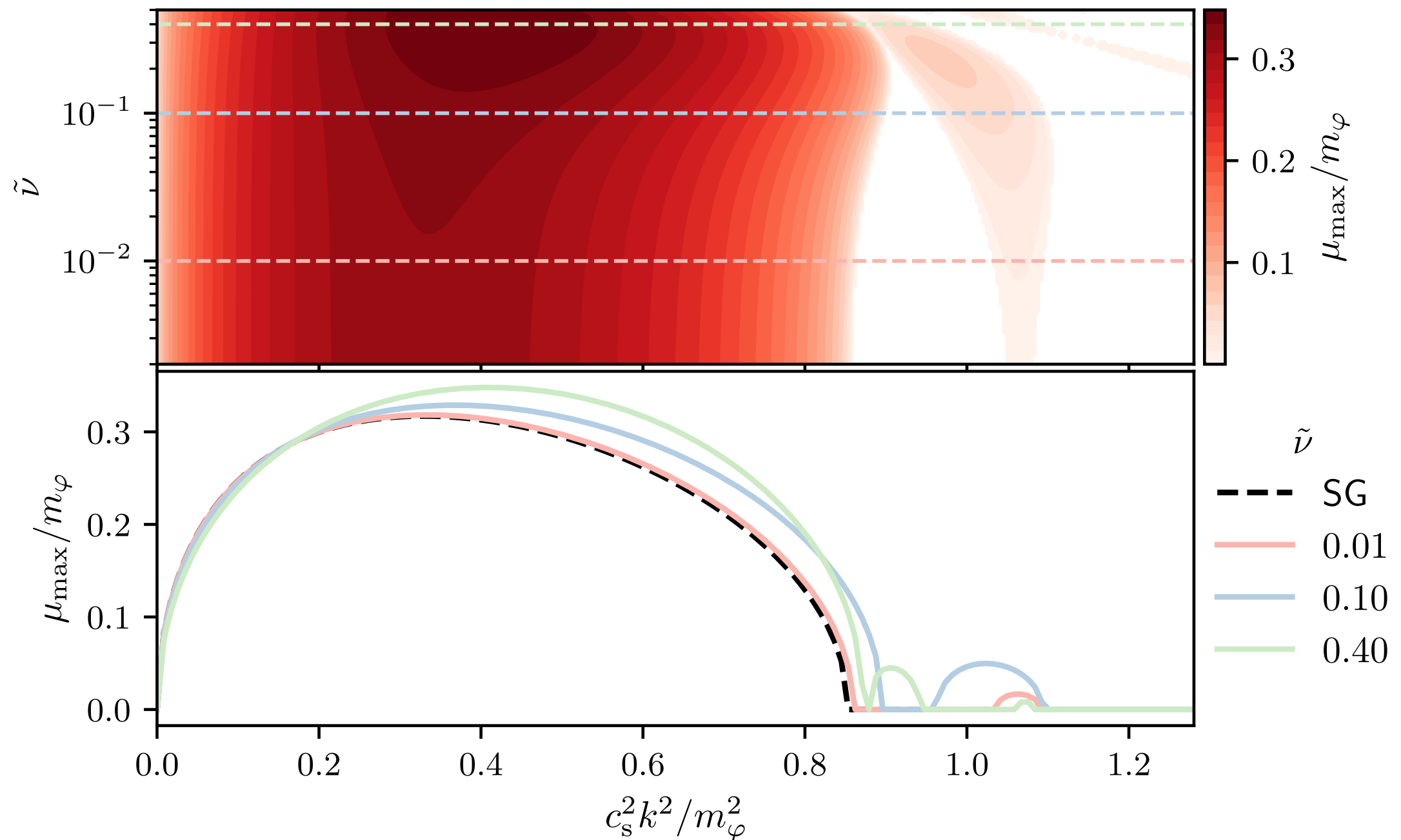
Linear Instability

$(\phi_{\max} = 0.25\pi)$



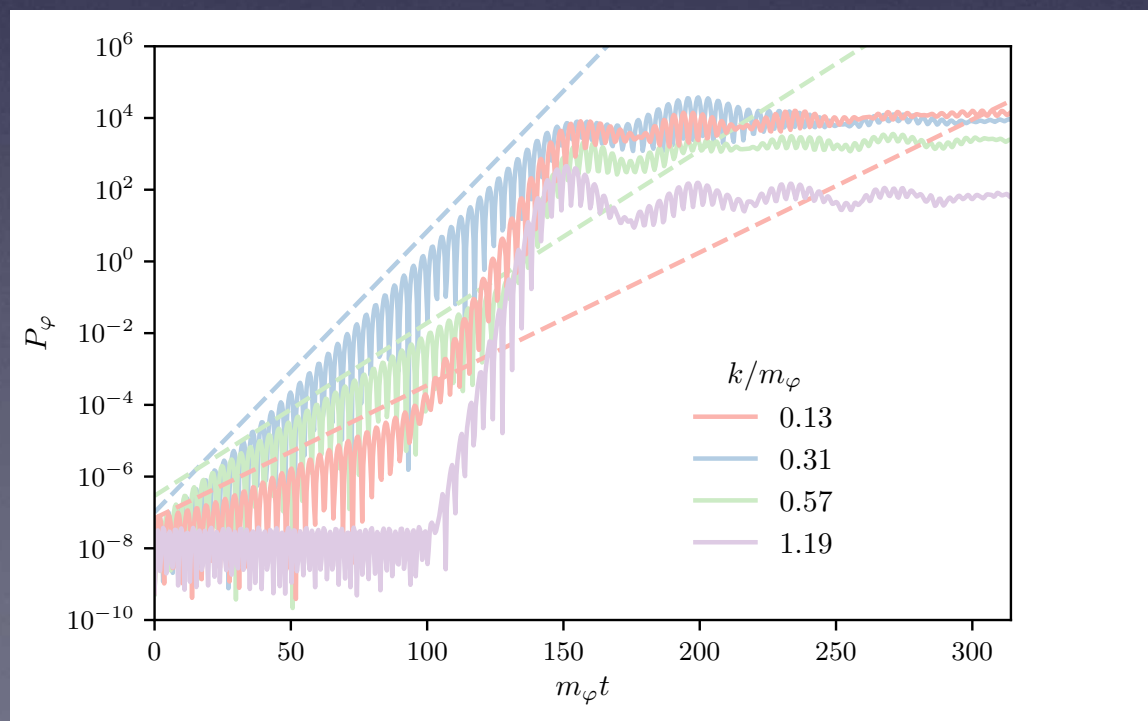
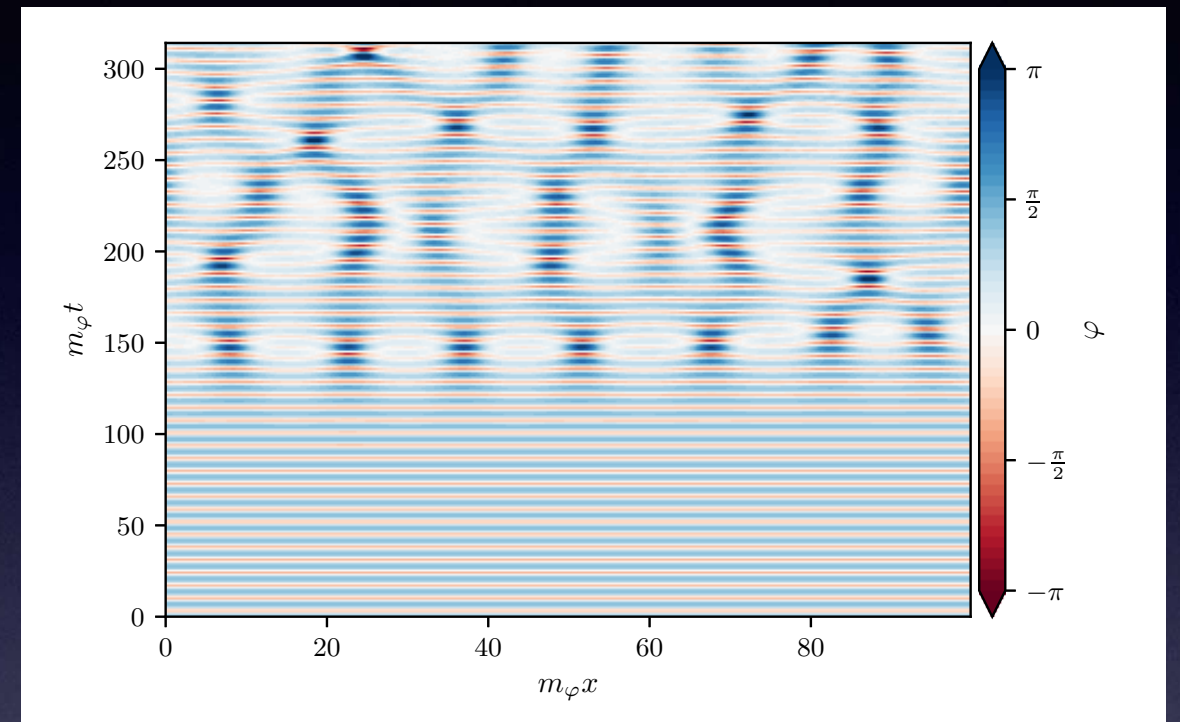
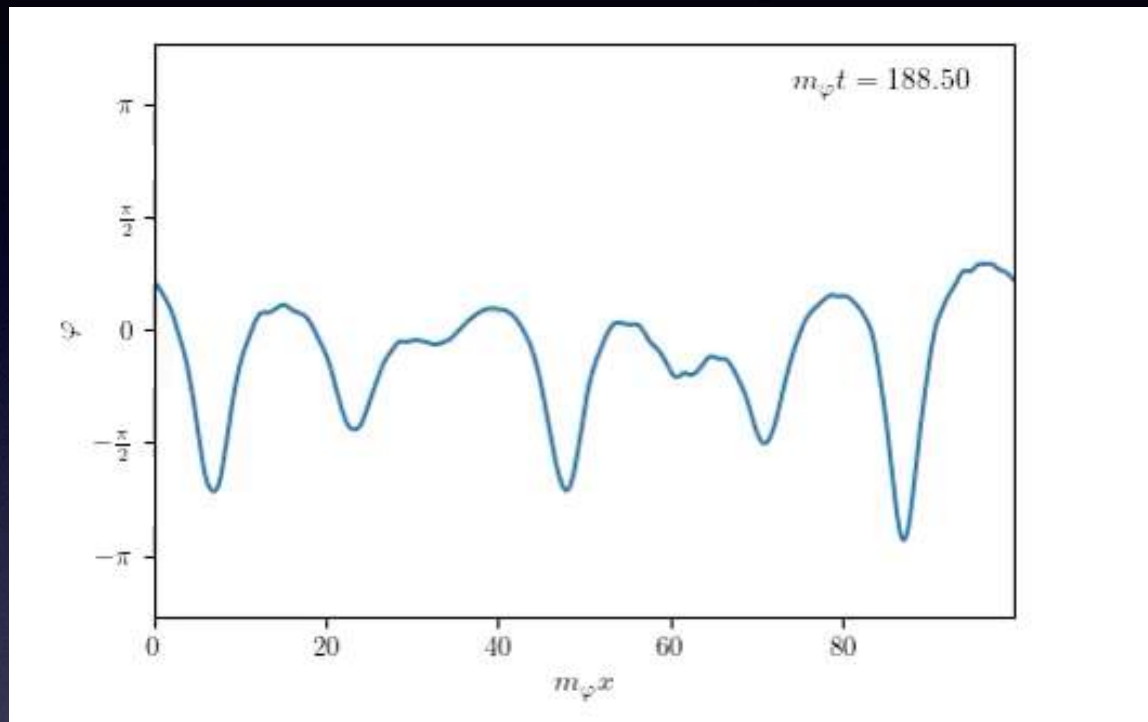
Linear Instability

$(\phi_{\max} = 0.75\pi)$



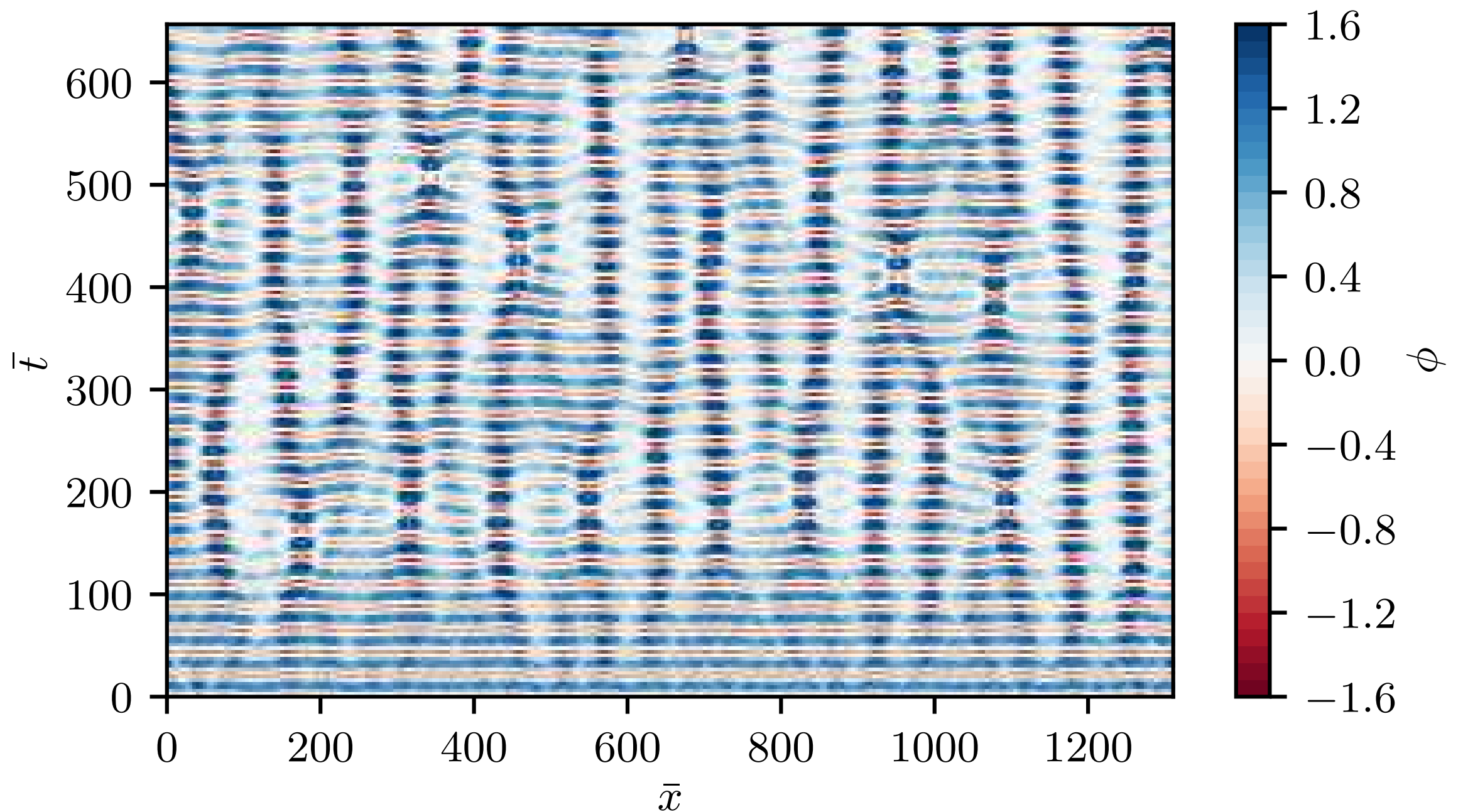
What About Nonlinearity

1D Analog Preheating

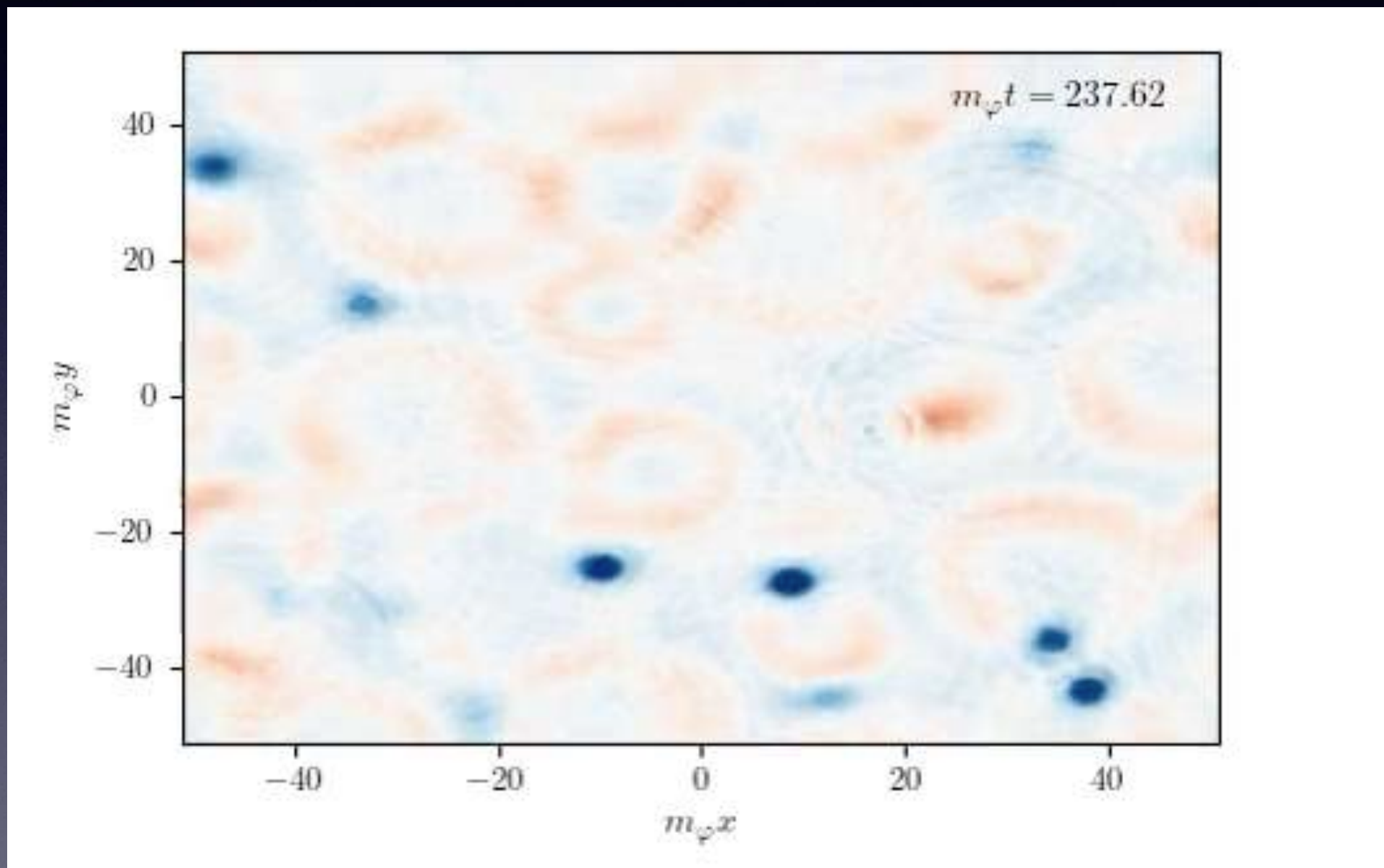


Early : Linear (Floquet) theory
Intermediate : Rescattering
Late : Solitary waves

Larger initial fluctuations (i.e. fewer particles)



2D Analog Preheating



Field Momentum $\dot{\phi}_{39}$ with mean removed

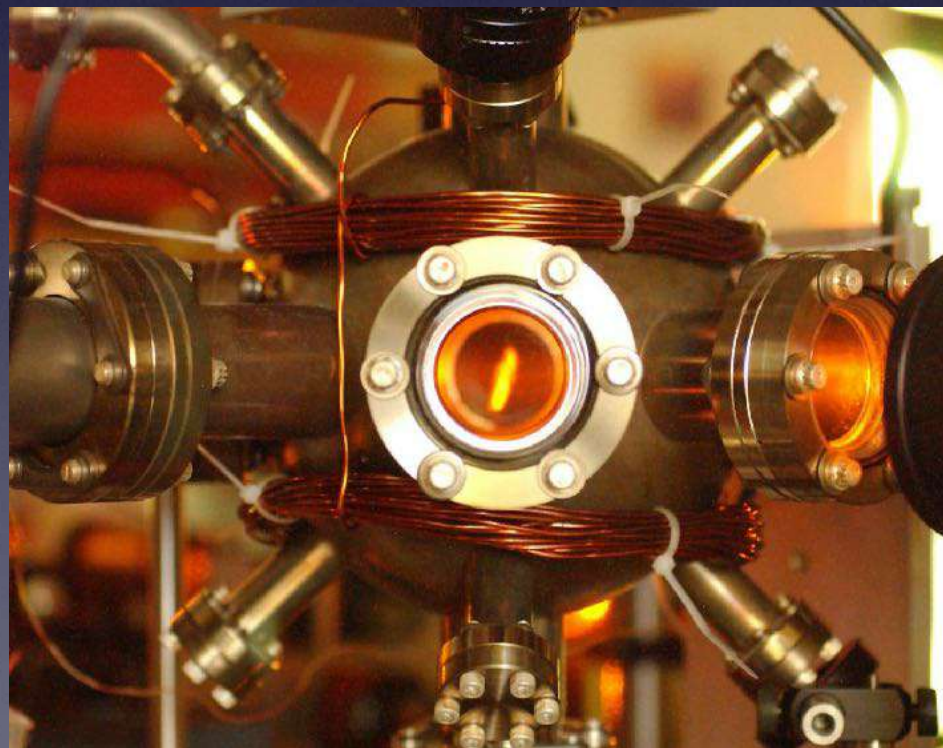


BUT ...

... Real BECs are Trapped

$$i\hbar\dot{\psi}_i = \left(-\delta_{ij} \frac{\hbar^2}{2m_i} \nabla^2 + \cancel{V(x)} + g_{ij} |\psi_j|^2 \right) \psi_i - \nu_{ij} \psi_j$$

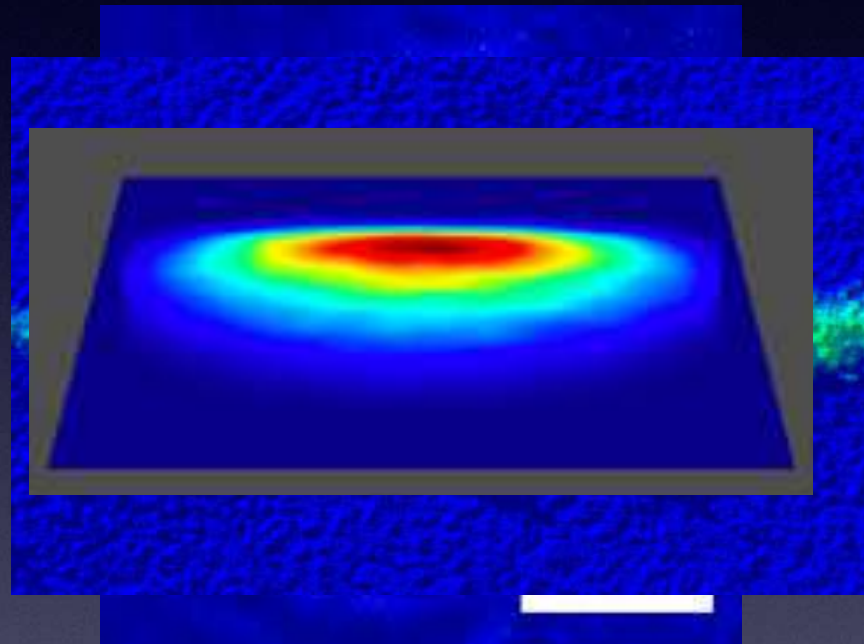
$V(x)=0$ homogeneous: $\bar{\psi}$ \rightarrow $V(x) \neq 0$ inhomogeneous $\bar{\psi}(x)$



Does parametric resonance still work?

Dimensional Reduction

Idea : Integrate out trapped directions



Pancake Trap

Cigar Trap

Doughnut Trap

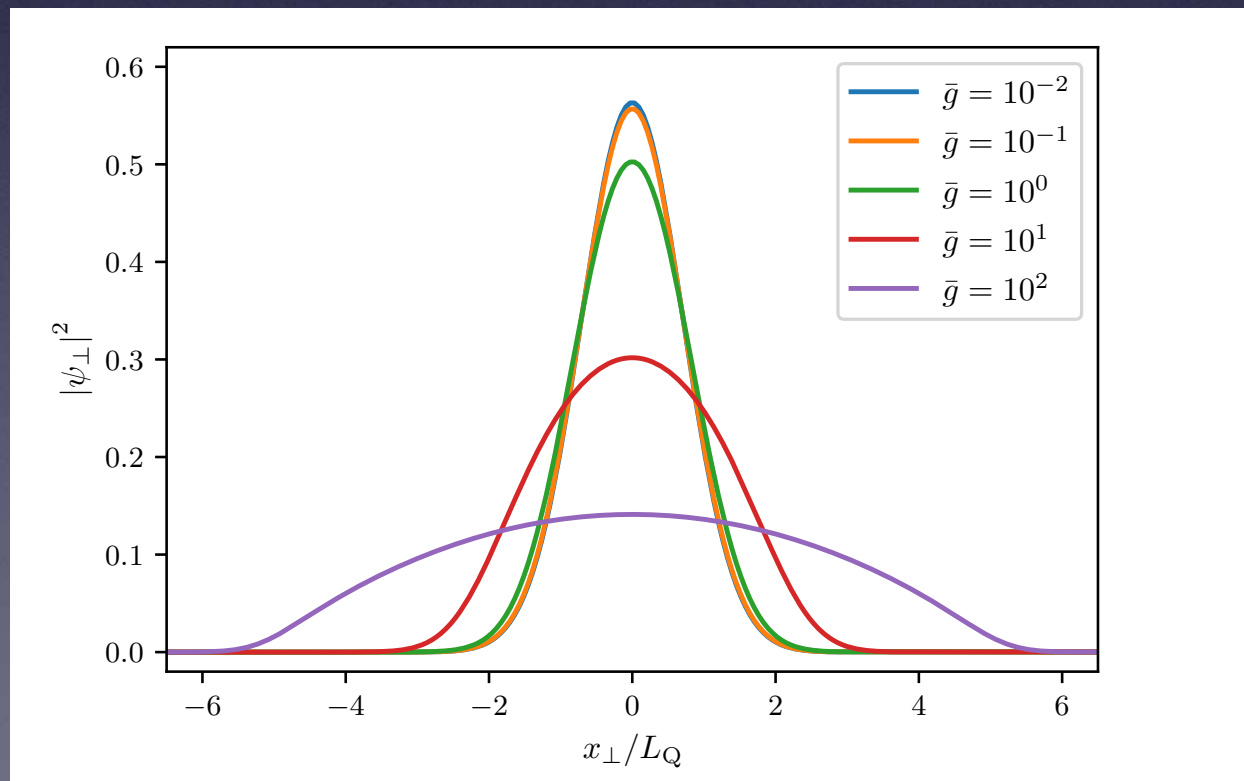
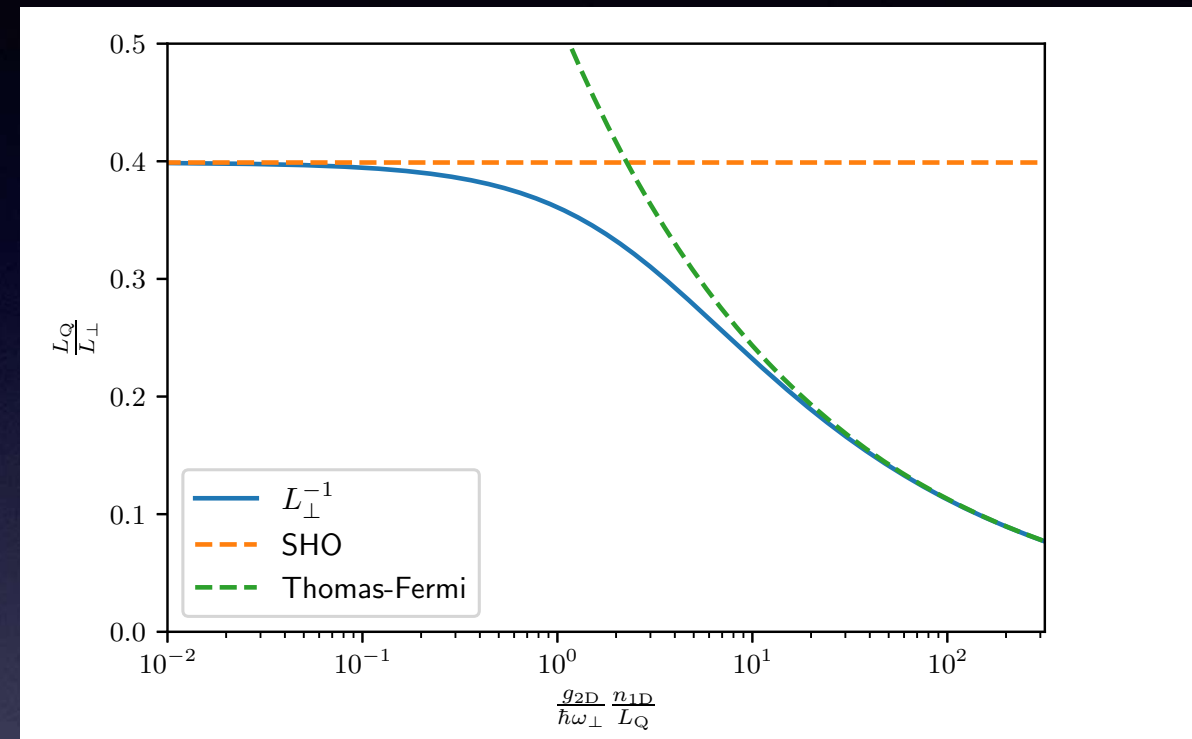
$$i\hbar\dot{\psi}_i = \left(-\delta_{ij} \frac{\hbar^2}{2m_i} \nabla^2 + V(\mathbf{x}) + g_{ij} |\psi_j|^2 \right) \psi_i - \nu_{ij} \psi_j$$

$$\frac{g^{D-1}}{g^D} = \frac{1}{L_\perp} = \frac{\int dx_\perp n^2}{\int dx_\perp n} = \frac{\int dx_\perp |\psi_\perp|^4}{\int dx_\perp |\psi_\perp|^2}$$

Harmonic Trap

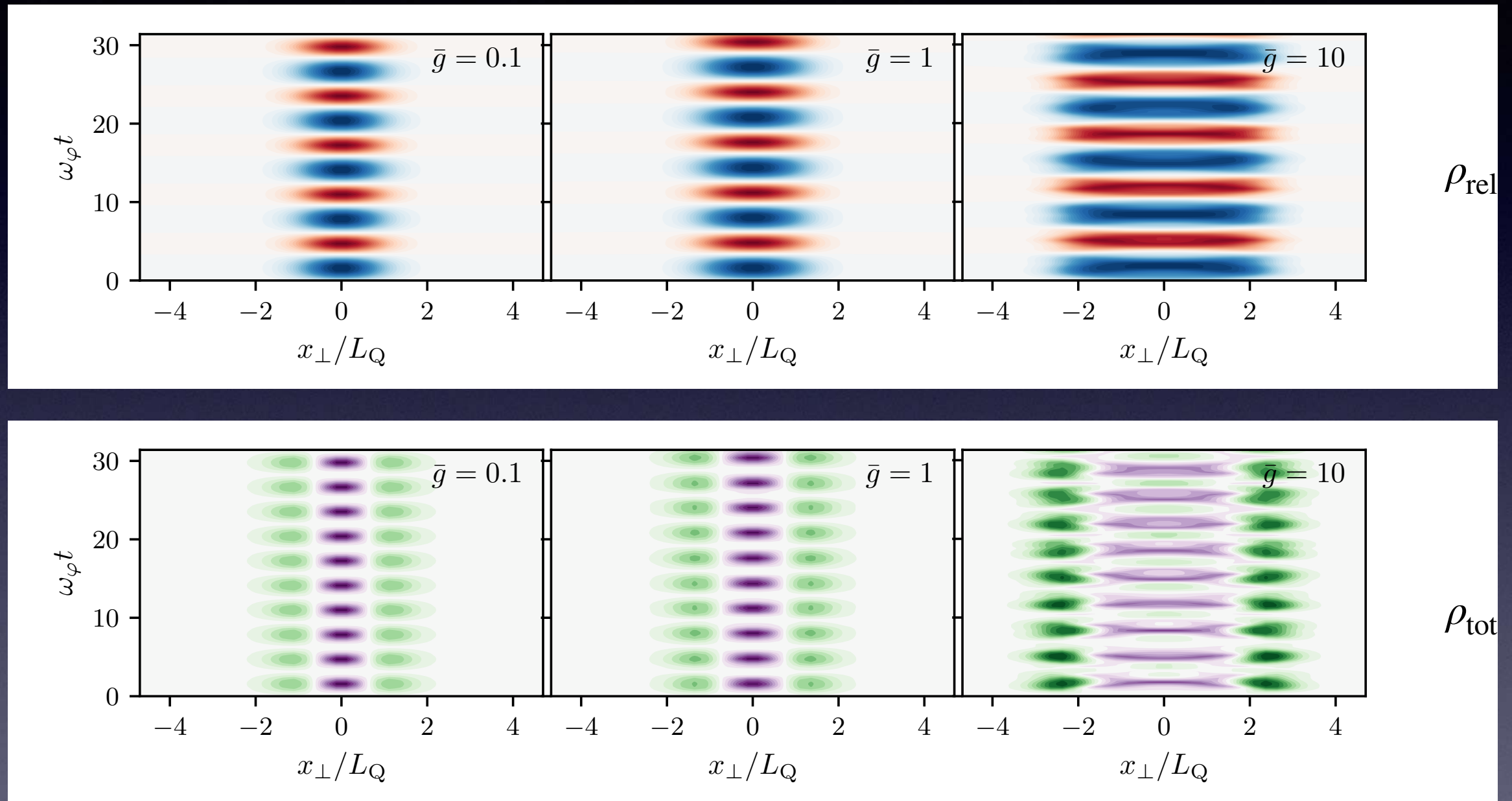
$$V(x_{\parallel}, x_{\perp}) = \frac{1}{2} m \omega_{\perp}^2 x_{\perp}^2 \quad L_Q^2 = \frac{\hbar}{m \omega_{\perp}}$$

$$\mu \psi_{\perp} = \left[-\frac{\hbar^2}{2m} \nabla_{\perp}^2 + V_{\text{trap}}(x_{\perp}) + g^D |\psi_{\perp}|^2 \right] \psi_{\perp}$$



$$|\bar{g}| \sim \frac{L_Q^2}{L_{\text{heal}}^2} \sim \frac{1}{\tilde{v}} \frac{m_{\phi}^2 L_Q^2}{c_s^2}$$

Background Evolution

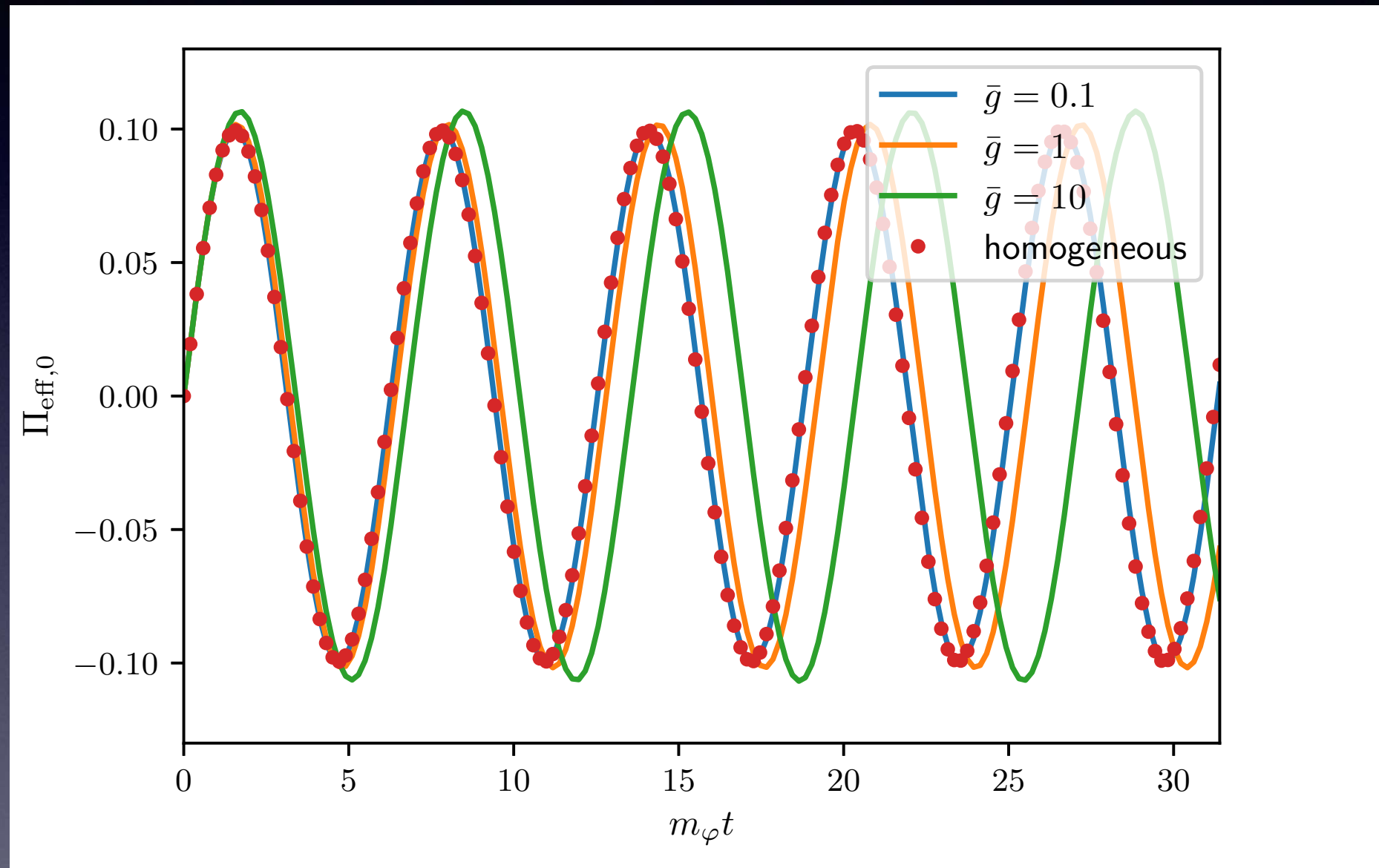


SHO limit

Thomas-Fermi Limit

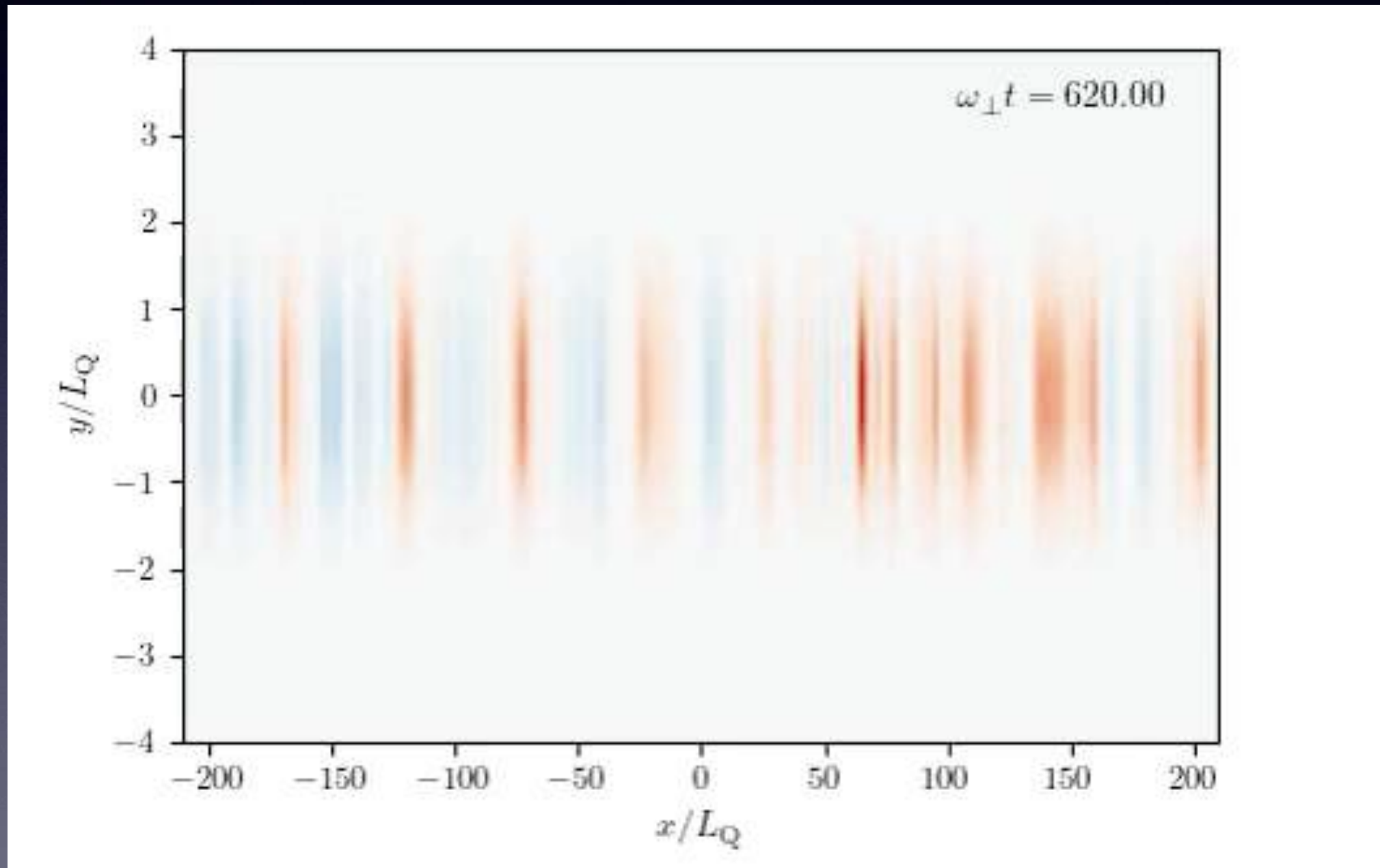
Decrease strength of transverse trap

Evolution Matches Dimensionally Reduced Simulations

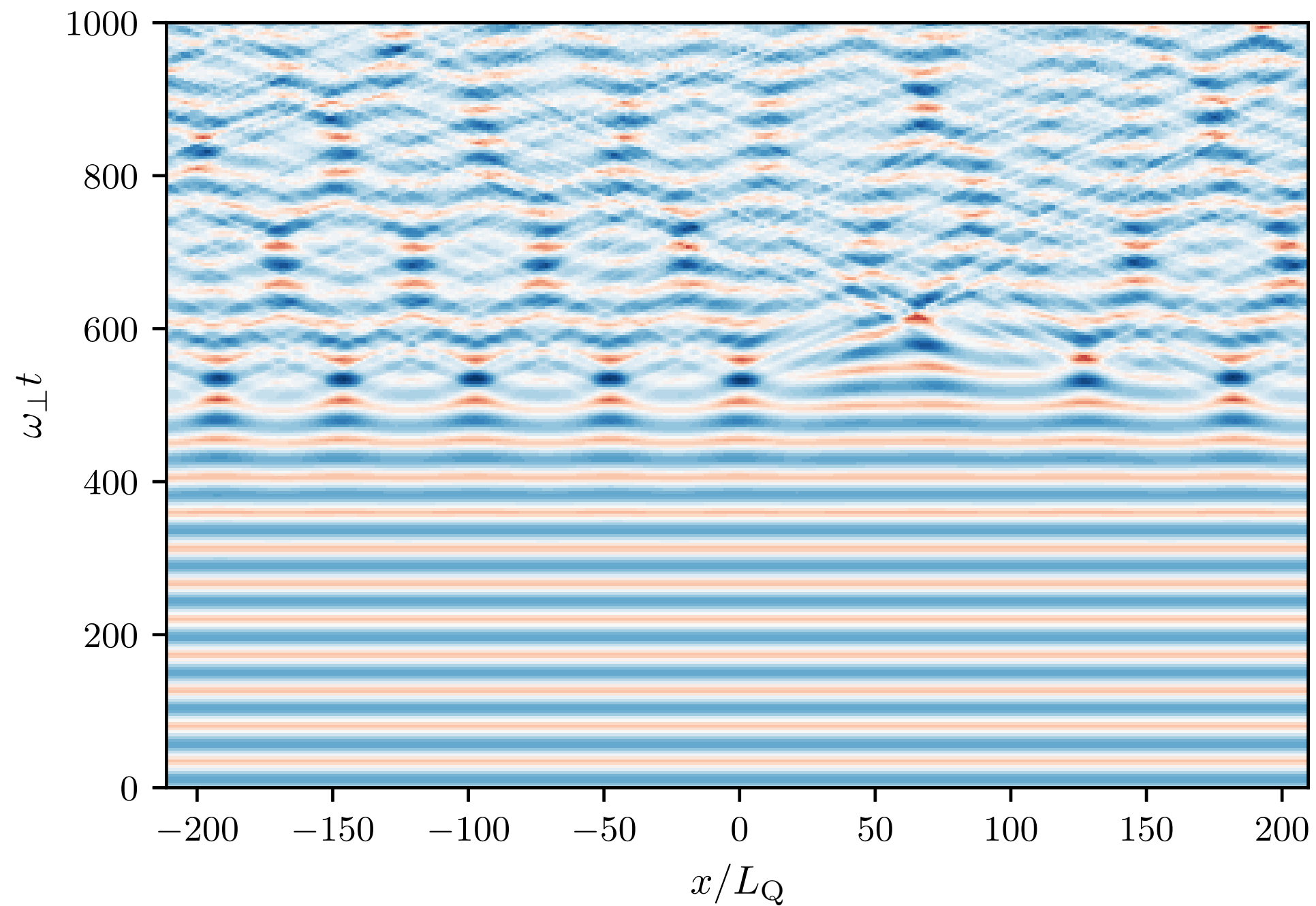


Is trapping maintained with parametric resonance?

Preheating in a Trap



Effective 1D Field Evolution



Summary

- Coupled dilute gas Bose-Einstein condensates can behave as relativistic fields
- Setups exist to mimic end-of-inflation (self-induced parametric resonance \rightarrow nonlinearity)
- Linear theory : BEC and cosmological calculation match, deviations controlled by tunable parameter
- Solitary waves in nonlinear regime
- Persists in more realistic case of trapped BEC
- Similar setup can also be used for vacuum decay